The MARS Code System User's Guide
Version 13(95)

Nikolai V. Mokhov

Fermi National Accelerator Laboratory
P.O. Box 500, Batavia, Illinois 60510

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Nikolai V. Mokhov
Fermi National Accelerator Laboratory
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Abstract

This paper is a user's guide to the current version of the MARS Monte Carlo code. MARS performs fast inclusive simulations of three-dimensional hadronic and electromagnetic cascades, muon and low energy neutron transport in shielding and in accelerator and detector components in the energy range from a fraction of an electronvolt up to 30 TeV. The code has undergone substantial improvements since the last documented version MARS 10 and all these as well as other specific features of the MARS code system are explained in detail. Descriptions of general input and output with commentary and recommendations are given. Examples are given for running the program with distributed sources, complex compounds, arbitrary geometries, and magnetic fields. Use of the code in a multistage mode, coupled with event generators (DTUJET), with the STRUCT program for tracking particles in accelerator lattices with beam loss recording, and with physics analysis and graphics packages is demonstrated with typical input and output examples.
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1 Introduction

The MARS code system is a set of Monte Carlo programs for inclusive simulation of three-dimensional hadronic and electromagnetic cascades in matter, of muon and low-energy neutron production and transport in radiation shielding, accelerator and detector components at energies up to 30 TeV. It allows fast cascade simulation with modest memory requirements in complex geometries with composite materials, in presence of arbitrary magnetic fields, with a variety of variance reduction techniques, other optimization and scoring options.

The MARS code has undergone substantial improvements since the last documented version MARS10 [1]. The purpose of this document is to be a user’s guide for a new MARS, version 13(95), so the paper reproduces the needed sections of the previous manual with a detailed description of all the main improvements and additional options. The main changes compared to MARS10 and to MARS12 released in 1992 [2], concern the hadron production model, muon and low-energy neutron production and transport, photo-hadron and photo-muon production, $e^+$, $e^-$, $\gamma$ and hadron production by muons, synchrotron radiation, muon decays, precise particle tracking in a magnetic field, extended geometry module, radioactivation calculation, computing performance, complete double precision mode, new input and output with significantly extended scoring and visualization capabilities. If needed, the code couples to the DTUJET event generator [3] and to the STRUCT code for multi-turn particle tracking in an accelerator lattice [4]. By 1995 there are about 50 MARS users worldwide.
2 MARS System

2.1 Base

Feynman's ideas concerning an inclusive approach to multiparticle reactions [5] and a biasing techniques served as a basis for the original program MARS [6] as well as for the program CASIM [7]. To construct a cascade tree only a fixed number of particles from each vertex is chosen (depending on the problem considered) and each carries a statistical weight which is equal, in the simplest case, to the partial mean multiplicity of the particular event. Energy and momentum are conserved on the average over a number of collisions. The practical reasons for the inclusive scheme as described in [8, 9, 10, 11] are:

- CPU time per incident particle grows only logarithmically with incident energy, compared to the linear rise in the exclusive mode, which allows the easier simulation of multi-TeV cascades;

- in many applications one considers effects due to the simultaneous interactions of a huge number of particles, so to describe the cascade it is sufficient to obtain the first moment of the distribution function using the inclusive cross-sections, in the same manner as with Boltzman's equation;

- experimental data on inclusive spectra are more readily available than on exclusive ones;

- the use of statistical weights allows the production of a given particle type to be enhanced within the phase-space region of interest, especially for rarely produced particles.

A disadvantage of this approach is the impossibility of directly studying fluctuations from cascade to cascade. Other codes have in the meantime adopted a similar approach to simulate very high energy electromagnetic[12] and hadronic[13] showers.

The mathematical foundation and the physical model of the MARS system as well as benchmarks for numerous applications are described in detail in [2, 9, 11, 14, 15]. The program has been developed over many years. Besides the original version [6] the milestones were: MARS3 [16], MARS4 [17], MARS6 [18], MARS8 [19], MARS9 [20], MARS10 [1, 21], MARS12 [2, 22], MARS93 [23], the hydrodynamical MARS/MESA/SPHINX package [24] and a parallel version of MARS12 [25].
2.2 Major Features

The main features of the MARS code system, version 13(95), in physics, geometry and analysis/computing categories are listed here.

1. Physics Model:

- all interactions of hadrons, leptons and photons during their passage through matter are taken into account for the energy range from 30 TeV down to about 0.2 MeV (0.00215 eV for neutrons);

- each $hA$ vertex is constructed in a problem-oriented way from a simplest case with one or two hadrons in the final state (e.g., leading particle bias), through the analog (exclusive) type vertex, to the sophisticated cases with enhancements in phase-space and/or particle type (e.g., $\bar{p}$ production [15, 26]);

- simulation of hadron-nucleus interactions at $E \geq 5$ GeV is based on a set of the semi-theoretical formulas for a proton target, coupled with the additive quark model of hadron-nucleus interactions for fast secondaries and phenomenological model for slow particles [11, 14, 27, 28, 29, 30, 31]; production of diffractive particles and nuclear de-excitation processes are modeled separately; at $E \leq 5$ GeV, hadron-nucleus inelastic collisions are simulated using modified formulas [32]; special treatment was added in the new version to provide better energy balance;

- optional coupling with DTUJET [3] event generator for primary high energy $pp$- or $\bar{p}p$-collisions [33, 34];

- hadron-nucleus inelastic cross-sections are calculated for a set of nuclei in the framework of the optical model and tabulated in the energy range from 10 MeV to 30 TeV for subsequent interpolation [11, 35];

- interactions and transport of neutrons in the 0.00215 eV to 14.5 MeV energy range (see Table 7 of Sect. 3.4) is performed using 28-group library BNAB [36] in the $P_3$-approach [22] with a special treatment for very extended systems like multi-section labyrinths [37, 38]; optional is 49-group library for neutrons below 18 MeV (see Table 8), married with 15-group system for secondary photons in the 10 keV to 11 MeV energy range (Table 9) [23, 36];
• special attention is paid to processes with a small momentum transfer: multiple Coulomb scattering using Molière’s theory with allowance for nuclear size effects, elastic scattering, diffraction, $\delta$-rays and direct $e^+e^-$ production by hadrons [9, 20, 35, 39, 40];

• fast precise algorithm for simulation of ionization energy loss [35, 41];

• edge scattering with step optimization depending on direction and location of a particle near the surface [22, 42];

• muon production with forced decays of mesons and short-lived resonances [9, 11, 43, 44];

• very efficient algorithms for muon interactions (ionization, bremsstrahlung, direct $e^+e^-$ pair, and deep inelastic) and transport [35] well advanced compared to previous versions [1, 11, 43, 44];

• optional forced $\mu\rightarrow e\nu\bar{\nu}$ decays and synchrotron radiation generation [45];

• photoneutron production in a giant resonant energy region [22, 46];

• leading particle biased simulation via modified AEGIS code [47] of electromagnetic showers, initiated by $\pi^0$ decays, high energy $\delta$-rays, prompt $e^+e^-$ from hadrons, muons and photoneutrons;

• phenomenological algorithm for radionuclide production with a point kernel techniques for gamma dose rate [22].

2. Geometry, Transport and Materials:

• iterative step-wise particle tracking with precise localization of boundaries, which is especially refined near matter-vacuum edges [11, 48, 49]; advanced algorithm for tracking in a magnetic field;

• multi-medium multi-component geometrical module with optional superfine structure distributed, if desired, over thousands of meters, in presence of arbitrary magnetic fields;

• optional extended geometry description [50] with allowance for easier handling of an arbitrary combination of boxes, cylinders, spheres and cones with some visualization options;

• 22 materials are built into the code and others can be easily defined by the user; each material can be a mixture of up to 6 chemical elements; up to 20 different materials can simultaneously be involved in a specific calculation;
• fusion with the STRUCT code [4] for multi-turn particle tracking in the accelerator lattice represented by an arbitrary combination of magnetic elements and transfer matrices allowing unified approach to beam loss and radiation effects studies at modern accelerators [40, 48, 51, 52, 53, 54];

• geometry and phase-space tagging options, which allow very efficient way to study a source term;

• variance reduction options: phase-space and particle type biasing, exponential conversion of path length, mathematical expectation, splitting and Russian roulette for particle trajectories and statistical weights [11].

3. Computing and Analysis:

• scoring of three-dimensional distributions of star density, total and partial particle fluences, total and partial energy deposition densities, temperature rise, dose equivalent, residual dose rate; corresponding statistical errors; energy spectra in the predefined regions; tagged distributions and some integral values;

• substantially extended (compared to the previous versions) histogramming and graphics interfaces [55, 56];

• intermediate dumps of the key distributions with a frequency defined by the user;

• machine independent universal random number generator [57];

• the complete double precision version 13.1(95) and the version 13.2(95) with all crucial computations in double precision mode are currently supported on SunOS, SGI-IRIX, IBM-AIX and HP-UX; with certain limitations the version 13.2(95) is available for VMS and MS-DOS.

Initial particle kinetic energy is from threshold to 30 TeV. Threshold kinetic energies are 0.00215 eV for neutrons, 0.2 MeV for electrons and photons (0.01 Mev optional), and 2 MeV for all other particles. Transformed particle types and their codes (ij) are given in Table 1. All other particle types generated at the event processing stage are converted at the production point into the states listed in the table: \( \pi^0 \)-mesons are decayed into two photons; \( \rho^- \), \( \omega^- \), \( D^- \), and \( J/\Psi^- \) particles are decayed into muons etc. Heavy fragments (\( d \), \( t \), \( \alpha \) and others) deposit their energy locally.
Table 1: Transported particle types.

<table>
<thead>
<tr>
<th>jj=</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
<th>6</th>
<th>7</th>
<th>8</th>
<th>9</th>
<th>10</th>
<th>11</th>
<th>12</th>
</tr>
</thead>
<tbody>
<tr>
<td>p</td>
<td>n</td>
<td>π⁺</td>
<td>π⁻</td>
<td>K⁺</td>
<td>K⁻</td>
<td>μ⁺</td>
<td>μ⁻</td>
<td>γ</td>
<td>e⁻</td>
<td>e⁺</td>
<td>(\bar{p})</td>
<td></td>
</tr>
</tbody>
</table>

### 2.3 Structure

The MARS code system consists of a few hundred FORTRAN77 subroutines accompanied (for convenience) by the CERN FFREAD package to read format-free data cards, and by a few CERN service C-routines. The system is linked to the CERN library to use the HBOOK package [55]. The output files are analysed with the PAW system [56] and with other graphics packages (TOPDRAWER, GNUPLOT, XMGR and KALEIDAGRAPH). The source code file structure is listed in Table 2.

After compilation with a `-r8` option, version 13.1(95), or without this option for version 13.2(95), all object files are combined into the library `m13lib.a`. File `m13.f` is present in the working directory to allow problem-dependent modifications if needed. File `m13hist.f` can be kept in the same directory. The main input file `MARS.INP` must be present in the directory. A few other files should be present if a corresponding option is active: `GEOM.INP` for extended geometry description; link to `dat` directory with `.ndt` files with low-energy neutron cross-sections; files generated with `DTUJET` and `STRUCT` (if any). The main output file `MARS.OUT` along with files to interface with PAW and other graphics systems, and a few additional files for a multi-stage mode, will be created in the directory. All the I/O file names can easily be re-defined by the user in the main program.

An example of the `MAIN` program with `HBOOK` and random number generator initializations, with statements to open the main input and output files as well as files to communicate with graphics packages, is presented in the following section.
Table 2: Source code structure.

<table>
<thead>
<tr>
<th>File name</th>
<th>Content</th>
</tr>
</thead>
<tbody>
<tr>
<td>m13.f</td>
<td>main program and user routines;</td>
</tr>
<tr>
<td>m13bldt.f</td>
<td>all the BLOCK DATA modules;</td>
</tr>
<tr>
<td>m13cstuff.c</td>
<td>a few C service routines from the CERN library;</td>
</tr>
<tr>
<td>m13hist.f</td>
<td>histogram set up and entry subroutines;</td>
</tr>
<tr>
<td>m13io.f</td>
<td>I/O and initialisation subroutines;</td>
</tr>
<tr>
<td>m13mareg.f</td>
<td>event processing and geometry module;</td>
</tr>
<tr>
<td>m13ph.f</td>
<td>all the physics simulation subroutines;</td>
</tr>
<tr>
<td>m13tr.f</td>
<td>all the particle transport subroutines;</td>
</tr>
<tr>
<td>m13util.f</td>
<td>utility subroutines, including FREAD.</td>
</tr>
</tbody>
</table>
2.4 MAIN Example

IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)

LOGICAL IND
COMMON/LOGIND/IND(20)
COMMON/INPOUT/NREAD,NWR,NWPSI,NWMJL,NWEGH
PARAMETER (LUNHIST=40)
PARAMETER (NH=2500000)
COMMON/PAWC/H(NH)
CHARACTER*20 HISTFILE
CALL HLIMIT(NH)
HISTFILE='MARS.HIST'
CALL HROPEN (LUNHIST,'HBOOK',HISTFILE,'N',I024,ISTAT)
CALL RNDMST(12,34,56,78)
CALL RNDMTE(0)
NREAD=15
NWR=16
OPEN (UNIT=NREAD,FILE='MARS.INP',STATUS='OLD')
OPEN (UNIT=NWR, FILE='MARS.OUT',STATUS='UNKNOWN')
CALL BEGINN
IF(IND(18)) THEN
OPEN (UNIT=11,FILE='GEOM.INP',STATUS='OLD')
CALL GREAD
END IF
CALL MHISET
NWEGH=17
OPEN(UNIT=NWEGH,FILE='MUON.EGH',STATUS='UNKNOWN')
OPEN(UNIT=20, FILE='MUON.PLOT',STATUS='UNKNOWN')
CALL MARSON
NWPSI=13
OPEN(UNIT=NWPSI,FILE='PSINEU.GRA',STATUS='UNKNOWN')
NWMJL=12
OPEN(UNIT=NWMJL,FILE='EDMJL.GRA',STATUS='UNKNOWN')
CALL SERVN
IF(IND(13)) CALL MARACT
CALL HCDIR ('//HBOOK',' ')
CALL HROUT (0,ICYCLE, ' ')
CALL HREND ('HBOOK')
END
2.5 Usage

First, MARS.INP and other input files are created by the user. Then, if needed, he/she deals with user subroutines in a file m13.f re-defining any of the default parameters and adding a variety of the features to the considered problem (see Chapter 4). If low-energy neutron transport is to be considered with option \( \text{IND (14) = T} \), the code will look for a directory ndt_dir. The following soft link should be created: `ln -s ~/mars13/dat ndt_dir`. The command files for compiling and linking MARS under different operating systems are given in Table 3 for version 13.1(95) in global double precision mode. If one uses version 13.2(95) with only all crucial computations in double precision mode, then \(-r 8\) option must be omitted. An executable file called rma will be created. For MARS code availability, comments and other related contact mokhov@fnalv.fnal.gov.

Table 3: Command files for version 13.1(95). The option \(-r 8\) (\(-qdp\) on IBM-AIX and \(-R 8\) on HP-UX) must be omitted for version 13.2(95).

<table>
<thead>
<tr>
<th>System</th>
<th>Command file</th>
</tr>
</thead>
<tbody>
<tr>
<td>SunOS</td>
<td><code>f77 -o rma -r 8 m13lib.a m13.f $ cernlib</code></td>
</tr>
<tr>
<td>SGI-IRIX</td>
<td><code>f77 -o rma -r 8 -mips2 m13lib.a m13.f $ cernlib</code></td>
</tr>
<tr>
<td>IBM-AIX</td>
<td><code>f77 -o rma -qdp -qextname m13lib.a m13.f $ cernlib</code></td>
</tr>
<tr>
<td>HP-UX</td>
<td><code>f77 -o rma -R 8 -K +ppu m13lib.a m13.f $ cernlib</code></td>
</tr>
</tbody>
</table>

3 General Input

3.1 Framework

There are three levels of the input data definition in MARS: defaults, input cards and user subroutines. All input data have some default values (see next two sections), so if these are acceptable for particular run their presence in the input sequence is not necessary. The shortest input consists of two cards in MARS.INP: a title and the STOP card.

Default values are built into the code. Input cards are defined by the user in two files: MARS.INP (main input sequence) and GEOM.INP (presented and used for the extended geometry description, if \( \text{IND (18) = T} \)). The names of these two files are defined in the main program (file m13.f) and can be changed. The main program also allocates the dynamic memory for HBOOK and gives control to the whole system. Histogram set up and entry routines, defined in the file m13hist.f,
can be re-defined by the user according to his/her specific needs. All the user sub-
routines are collected in the m13.f file in the dummy form, i.e. a subroutine name
followed by the RETURN and END statements. Corresponding subroutines can be
modified by the user if necessary.

The user can re-define any of the default parameters with the card input se-
quence in the file MARS.INP. This consists of option and data cards which allow
easy definition of beam parameters, materials, energy thresholds, termination con-
ditions, hA-vertices parameters, geometry and scoring parameters. Multi-medium
3-D cylindrically symmetric (r-z-phi) "standard" geometries and sources can be de-
\n
The "extended" geometry option extends this friendly way to a combination of boxes, cylinders, spheres and cones, defined in the input file GEOM.INP. In the case of complex composite materials, of very complex geometries, of arbitrary sources and magnetic fields appropriate
user subroutines should be provided (see Chapter 4). Arbitrarily complex geometry is usually built into "standard" or "extended" geometry.

The MARS.INP file consists of a title card in a A80-format followed by some
data cards. The standard CERN FFREAD package is used to read these format-free
data cards in the routine BEGINN. The structure of all these cards is the same:

\[ \text{KEYW a1 a2 ... N1= b1 b2 ... N2= c1 c2 ...} \]

where KEYW is a keyword assigned to the group of FORTRAN variables a(i),
b(i) and c(i); N1 and N2 are the addresses of the arrays b(i) and c(i). The
variables may be of the following types: integer, real and logical (represented by
T or F). Items are separated by blanks. The order of the the data cards and their
number are arbitrary. The only requirement is: the STOP card must be included
and must be the last in the sequence. The card image input sequence is described
in detail in the following section and cumulatively in Table 6 of Section 3.4. Ex-
amples of the input sequence are given in Section 3.5.

The units in MARS are: energy in GeV, dimensions in cm, azimuthal angle \( \phi \)
in degrees, temperature rise \( \Delta T \) in degrees Centigrade. The reference system is
global Cartesian coordinate system (GCS). Typically the origin \( x = y = z = 0 \)
coincides with the source starting point. Any number of local coordinate systems
(LCS) can be placed in it via GEOM.INP and/or user routines REG1, REG2, ALIGN
and SAGIT. In the global system the z is longitudinal, and the positive direction is
from left to right. The z axis is usually the center of symmetry, and as a rule the
beam strikes along this axis in the positive direction. The positive direction of the
x axis of the global system is up and the y axis is toward the viewer, completing a
right-handed system.
3.2 Card Input Sequence (MARS.INP)

The first card is a variable MTEXT (FORMAT (A80))—title of the problem, prints out a heading for the output.

Then some number of unformatted data cards follows.

**INDX**  IND(20)
Logical variables which control the options. Default: 20*F.

**IND(1)=T**—the program calculates and prints distributions of $e^+e^-$ and photon fluxes and of energy deposition density $\epsilon$ and related values: dose equivalent, instantaneous temperature rise $\Delta T$ at given initial temperature $T_0=\text{TEMPO}$ and number of particles per beam $N_0=\text{ANTN}$ (see VARS), and contact dose due to induced radioactivity at $N_0$ beam intensity. Other values discussed in Section 5.1 are calculated independently of the **IND(1)** meaning.

**IND(1)=F**—does not calculate the energy deposition related values and reduces execution time drastically in the TeV energy region.

**IND(2)=T**—initiates "Z—sandwich standard" geometry as a basis (see ZSEC and RSEC).

**IND(2)=F**—initiates "R—sandwich standard" geometry as a basis (see ZSEC and RSEC).

**IND(3)=T**—initiates calling of user subroutines described in Chapter 4 and defined in a file m13.f.

**IND(3)=F**—no user subroutines calling.

**IND(4)=T**—indicates the presence of magnetic or electric fields in the system; field components must be defined in a user subroutine FIELD.

**IND(4)=F**—no magnetic or electric fields.

**IND(5)=T**—activates $e^+, e^-, \gamma$ and $h$ production by muons at **IND(10)=T** and knock-on electron and $e^+e^-$-pair production by hadrons; **IND(5)=F**—above production mechanisms are not activated.

**IND(6)=T**—provides the use of mathematical expectation method: scoring of probabilities for a fluence calculation rather than direct (analog) particle contributions; can be effective for "deep penetration" problem (thick shields etc.); use with care, test first.
IND (6) = F – "analog" scoring of transported particles.

IND (7) = T – indicates that an incident particle interacts with probability EFF (see VARS) with the point-like target placed in the system at coordinates \((x_0, y_0, z_0)\) and starts with probability \((1 - EFF)\) from this point interacting with the rest of the system. By convention the material index for this target is equal to IM=1, and can be easily re-defined in a BEG1 user subroutine.

IND (7) = F – no point-like target.

IND (8) = T – provides the maximum amount of output to be printed.

IND (8) = F – corresponds to standard output (see Section 5.1).

IND (9) = T – provides the use of special algorithm (as in [22, 42]) for construction of particle trajectories prior to the first two inelastic nuclear interactions (edge scattering problem) and of Landau fluctuations simulation.

IND (9) = F – does not use the above sophisticated algorithms (time saving).

IND (10) = T – activates forced muon production in long and short lived meson decays as well as produced and incident muon transport with all possible interaction processes included (see Section 2.2).

IND (10) = F – turns off muon option (much faster).

IND (11) = T – provides an azimuthal structure of scoring; adds the \(\phi\)-dimension to the standard \((r-z)\) geometries.

IND (11) = F – no azimuthal division option.

IND (12) = T – activates biased antiproton production at every \(\Lambda\)-vertex above threshold with a consequent transport.

IND (12) = F – no forced antiproton production.

IND (13) = T – provides specific activity and residual dose rate calculation.

IND (13) = F – no induced radioactivation calculation.

IND (14) = T – provides multi-group neutron transport in the 0.00215 eV to 14.5 MeV energy range using the default 28-group neutron cross-section library; includes photoneutron production at \(\text{IND (1)} = T\).

IND (14) = F – turns off low-energy transport; their contribution into energy deposition distribution is considered in a simplified way (much faster).


IND (15) = F – no splitting for neutrons.
IND (16) = T – variance reduction via low-energy neutron delta-scattering and weight window at IND (14) = T.
IND (16) = F – no delta-scattering and weight window for neutrons.
IND (17) = T – initiates using of the 49-group cross-section library for neutron transport below 18 MeV and photon production and transport in the 0.01 to 11 MeV energy range produced by these neutrons at IND (14) = T.
IND (18) = T – initiates using of the extended geometry algorithm with a reading of an input file GEOM.INP (see Section 3.3).
IND (18) = F – no extended geometry description and a file GEOM.INP presented.
IND (19) and IND (20) are not used in the current version.

NEVT NSTOP, NTIME
NSTOP – number of events (or incident particles) to be run. Default: 200.
NTIME – number of intermediate dumps of key results, defined in a user subroutine DUMP. Default: 0.

ENRG E0, EM, PSTAM, EMI, EMNU
E0 – incident particle kinetic energy. Default: 100.
EM – hadron energy cutoff for flux and spectra scoring. Default: 0.0145.
PSTAM – star production threshold momentum. Default: 0.3 GeV/c (corresponds to 47 MeV for nucleons and 191 MeV for pions).
EMI – muon energy cutoff for flux and spectra scoring. Default: 0.005.
EMNU – neutrino energy cutoff. Default: 0.5.

IPIB I0, IBEAM
I0 – incident particle type (see Table 1). Default: 1.
IBEAM – type of the incident beam, represents the beam profile:
0 – laterally infinitesimal beam;
1 – beam is distributed uniformly in the rectangular area with SIXX and SIYY half-sizes (along the corresponding axis);
2 – beam is Gaussian with R.M.S. equal to $\sigma_x = $SIXX and $\sigma_y = $SIYY;
3 – beam is Gaussian as defined for $\text{IBEAM}=2$ and additionally has a Gaussian angular spread with R.M.S. equal to $\sigma(\theta_x) = $SITX and $\sigma(\theta_y) = $SITY in radians. Default: 0.
BEAM  SIXX, SIYY, SITX, SITY
       SIXX, SIYY, SITX, SITY—defined above, active if IBEAM $\geq 1$.
       Default: 0 0 0 0.

INIT  XINI, YINI, ZINI, DXIN, DYIN, DZIN
       XINI, YINI, ZINI—are initial $x_0, y_0, z_0$ coordinates of the beam spot center.
       Default: 0 0 0.
       DXIN, DYIN, DZIN—are initial direction cosines of the beam.
       Default: 0 0 1.
       The user can provide his own arbitrary source in a subroutine BEG1 (See Section 4.2).

SMIN  STEPEM, STEPH
       STEPEM—is an accuracy (in cm) of boundary localization in the iterative transport algorithm,
       this is a second (in addition to NSTOP) control parameter for calculational accuracy.
       Recommendation: $\text{STEPEM} \approx 0.3 \times t_{\text{min}}$, where $t_{\text{min}}$ is a smallest
       dimension of the smallest cell in the considered system. The larger $\text{STEPEM}$, the faster
       calculations. Default: 0.3.
       STEPH—is a parameter which modestly controls calculational time per event
       ($\approx \log(\text{STEPH} / \text{STEPEM})$) and accuracy of construction of particle trajectories
       up to the first two inelastic nuclear interactions for the default case IND (9) = F. Recomendation:
       $\text{STEPH} = \min(\lambda, l)$ where $\lambda$ is the mean inelastic length for hadrons and $l$ is the length
       of the characteristic zone in the direction of predominant propagation of the particles.
       Default: 10.

VARS  EFF, DLEXP, TEMPO, AINT
       EFF—point-like target efficiency; active if IND (7) = T. Default: 0.
       DLEXP—is an exponential conversion factor, which allows for increasing
       (DLEXP $\geq 1$) or decreasing (DLEXP $\leq 1$) the effective inelastic mean free path;
       useful correspondingly for “deep penetration” problem and for compact restricted systems.
       Recomendation: $0.3 \leq \text{DLEXP} \leq 3$. Default: 1.
       TEMPO—is an initial temperature $T_0$ in Kelvin, defined in the region
       $4 \leq T_0 \leq 1800$. Default: 300.
       AINT—is a number of particles per beam $N_0$; necessary for calculations of
       a temperature rise ($N_0$ is assumed to be for a single instantaneous spill) and
       crude estimation of a contact residual dose (here $N_0$ is assumed to be an
       average beam intensity in protons per second). Default: $10^{12}$. 

16
**NMAT**

**NREMA**

NREMA- is a number of different materials to be included in the particular run. Materials can be single elements (built-in or defined by the user) or complex compounds. The subroutine **BEGINN** calculates the effective atomic masses and numbers, cross-sections, radiation lengths and some other quantities. \(1 \leq \text{NREMA} \leq 20\). Default: 1.

**MATR**

**AMA(20)**, **ROW(20)**, **ATW(20)**, **ZAW(20)**

This determines the NREMA specific materials. Special index \(\text{IM}=i\) is assigned to each material according to the input order. By definition \(\text{IM}=0\) corresponds to vacuum and \(\text{IM}=\text{NREMA}+1\) to outside of the global volume (blackhole) and the user does not need to include them in the MATR sequence. There are 22 other built in materials: H, LHE, LI, BE, CH2, CH, C, TISS, WATR, AIR, SOIL, CONC, Al, SI, LAR, TI, FE, CU, W, PB, U, BLA (see Table 5, Section 3.4). If \(\text{AMA}(i)\) coincides with one of these names, one does not need to define **ROW**(i), **ATW**(i), **ZAW**(i) for this i. If not, the user must define the material with its corresponding parameters as follows:

i) for single elements – arbitrary name \(\text{AMA}(i)\), density **ROW**(i) in \(g/cm^3\), atomic mass **ATW**(i) and atomic number **ZAW**(i).

ii) for each composite material – put in the MATR sequence the word ‘MIXT’ and density **ROW**(i) and provide the corresponding definition of its composition in a subroutine **MIXTUR** (see Section 4.1), according to the input order. Default: FE.

**NLNG**

**LZ, NLZ**

**LZ**– is a number of longitudinal sections in the system for the standard geometry sector. \(1 \leq \text{LZ} \leq 50\). Default: 1.

**NLZ**– multiplies number of sections, described by **ZSEC**, i.e. repeats these sections \(\text{NLZ}\) times; active only for \(\text{NLZ} \geq 2\). Default: 1.

**ZSEC**

**ZSE(50)**, **IZN(50)**, **IZI(50)**, **IZM(50)**

**ZSE**(i) – z-coordinate of the right boundary of i-section. Default: 100., 49*0.

**IZN**(i) – number of subsections in i-section. Default: 50*1.

**IZI**(i) – material index \(\text{IM}\) of i-section, active for \(\text{IND}(2) = \text{T}\). Default: 50*1.

**IZM**(i) – magnetic field index \(\text{MA}\) of i-section, active for \(\text{IND}(2) = \text{T}\) and \(\text{IND}(4) = \text{T}\). Default: 50*0.
Maximum total number of the longitudinal subdivisions in the standard sector of the current MARS version is $MZ=250$ and can be easily increased in the extended sector and/or in a user subroutine REG1. In the program by definition the smallest z-coordinate is $ZMIN$, most often $ZMIN=z_0=0$. The maximum longitudinal dimension of the system is $ZMAX$ – the largest z-coordinate.

**NLTR**

$LR$ is a number of radial sections in the system for the standard geometry sector. In the current version $LR \leq 20$. Default: 1.

**RSEC**

RSE(20), IRN(20), IRI(20), IRM(20)

RSE(i) – radius of i-th lateral section. Default: 5, 19*0.

IRN(i) – number of radial subsections in i-th lateral section. Default: 20*1.

IRI(i) – material index IM of i-th lateral section, active for IND(2) = F. Default: 20*1.

IRM(i) – magnetic field index MA of i-th lateral section, active for IND(2) = F and IND(4) = T. Default: 20*0.

Maximum total number of the radial subdivisions in the standard sector of the current MARS version is $MR=20$ and can be easily increased in the extended sector and/or in a user subroutine REG1. For IND(11) = F the overall restriction in the standard geometry sector is $MZ \times MR < 5000$. Total maximum number of regions is 10000. The maximum radial dimension of the system is $RMAX$.

**NAZM**

$NF$ is a number of azimuthal bins, $1 \leq NF \leq 60$. To be presented in the input sequence only if IND(11) = T. Default: 1.

For IND(11) = T the overall restriction for the standard geometry sector is $NFZP = MZ \times MR \times NF \leq 10000$. Total maximum number of regions is 10000.

**AZIM**

FIB(60)

The azimuthal grid, $0 \leq FIB(i) \leq 360$. To be presented in the input sequence only if IND(11) = T and NF $\geq 2$. Default: 60*0.

**NOBL**

NOB is a number of special regions in which particle energy spectra and other histograms will be scored, $0 \leq NOB \leq 3$. Default: 0.
**RZOB**  
**RZO(4,3)**  
Presented in the input sequence only if $NOB \geq 1$. Determines the sizes of $i$-region for particle energy spectra and other histogram scoring:
- **RZO(1,i)** - minimum radius $RMI(i)$, $0 \leq RMI \leq RMAX$.  
  Default: 0.
- **RZO(2,i)** - maximum radius $RMA(i)$, $RMI \leq RMA \leq RMAX$.  
  Default: 0.
- **RZO(3,i)** - left $z$-coordinate $ZMI(i)$, $ZMIN \leq ZMI \leq ZMAX$.  
  Default: 0.
- **RZO(4,i)** - right $z$-coordinate $ZMA(i)$, $ZMI \leq ZMA \leq ZMAX$.  
  Default: 0.

**PLOT**  
**RPLOT, ZPLO1, ZPLO2, ZPLO3**  
Control a simple plotting of the geometry cross-sections.  
- **RPLOT** - maximum $x$ or $y$ coordinate to be presented in plots;
- **ZPLO1** - $z$-coordinate for the first cross-section of geometry;
- **ZPLO2** - $z$-coordinate for the second cross-section of geometry;
- **ZPLO3** - maximum $z$-coordinate of geometry to be printed.
  If all these numbers are equal to zero there is no plotting.
  Default: 0 0 0 0.

**NEUS**  
**NG, NGOUT, NINT, NGMAX(20)**  
- **NG** - number of neutron energy groups in the BNAB/MARS cross-section library. Default: 28.
- **NGOUT** - cutoff group number (see Tables 7 and 8 below). Default: 26.
- **NINT** - number of intervals in the angular flux presentation. Default: 20.
- **NGMAX(i)** - cutoff group number with group-to-group transfers for a material with $IM=i$. Default: $20*11$.
  This card can be presented at $IND(14) = T$.

**NDET**  
**NDE1**  
- **NDE1** - number of point-like detectors for energy spectrum, flux and energy deposition of low-energy neutrons at $IND(14) = T$. $0 \leq NDE1 \leq 10$.
  Default: 0.

**FLOC**  
**RD, XD(10), YD(10), ZD(10)**  
- **RD** - radius of the point-like detectors for low-energy neutrons. Default: 0.5.
- **XD(i), YD(i), ZD(i)** - coordinates of $i$-detector. Default: 30*0.
  This card is presented at $IND(14) = T$ and $NDE1 \geq 1$. 
NCHO  NTYZ(6)
    NTYZ(i) - parameters of the systematic selection for low-energy neutrons 
at IND(14) = T. Default: 2, 2, 1, 3*0.

FLCH  FL(200)
    FL(i) - importance function for low-energy neutrons at IND(14) = T. De-
default: 200*0.

MUON  KPHA, IDNDX
    KPHA - number of hadron generations to follow, i.e. number of levels in hA 
    vertex tree. Default: 30.
    IDNDX - activates muon angular distribution scoring in the NOB special re-
geons at IND(10) = T, if IDNDX=1. Default: 0.

STOP  Stops the input session
    Any cards given after this are meaningless.

A few examples of the input sequence are presented in Section 3.5.
3.3 Extended Geometry (GEOM.INP)

As indicated above, \( \text{IND (18)} = T \) extends a user-friendly \( r-z-\phi \) geometry description in a file MARS.INP to a combination of boxes, cylinders, spheres and cones, placed into the global coordinate system (GCS). The code reads via the GREAD routine volume data defined in their own local coordinate systems (LCS) in the input file GEOM.INP. Data in the file are unformatted, separated by blanks. In the version 13(95) of MARS the maximum number of volumes defined in this file is equal to NVMAX=1000.

First line of the GEOM.INP file contains NVOLUM and NEXGM. Here NVOLUM is total number of the volumes in the extended geometry sector, and NEXGM is NFZP (total number of regions in the standard geometry sector), or some bigger number to start a volume numbering in the GEOM.INP file with NEXGM+1. Limitations:

\[ 1 \leq \text{NVOLUM} \leq \text{NVMAX} = 1000, \text{NFZP} \leq \text{NEXGM} \leq (10000 - \text{NVOLUM}). \]

Next NVOLUM lines in the input file describe in an arbitrary order each volume:

\[ \text{NV, NT, IM, MA, XR, YR, ZR, C1, C2, ...} \]

where NV is volume sequential number in the extended geometry sector;

\[ \text{NEXGM} \leq \text{NV} \leq \text{NEXGM} + 1000; \]

\[ \text{NT} \text{ - volume type; } 1 \leq \text{NT} \leq 4; \]

\[ \text{IM} \text{ - material number (index) of given volume; } 0 \leq \text{IM} \leq 20; \]

\[ \text{MA} \text{ - magnetic field index; } 0 \leq \text{MA}; \]

\[ \text{XR, YR, ZR} \text{ - coordinates of a reference point (RP) of the LCS for given volume;} \]

\[ \text{C1, C2, ...} \text{ - parameters of given volume.} \]

By convention, the first volume in the above list with NV=NEXGM+1 is the global volume defined in MARS.INP. It contains all other volumes in standard, extended and user-supplied (REG1) geometry sectors. Usually, this is a "big" cylinder (most often with IM=0). A current set of the volume types and description parameters are given in Table 4.
Table 4: Current set of volume types in GEOM.INP.

<table>
<thead>
<tr>
<th>NT</th>
<th>Name</th>
<th>Description</th>
</tr>
</thead>
</table>
| 1  | BOX  | $C_1$ and $C_2$—half-sizes along $x$ and $y$ axes,  
     |      | $C_3$—box length along $z$ axis;  
     |      | $RP$—center of the lowest $z$ plane;  
     |      | $x$ and $y$ LCS axes are parallel to two perpendicular  
     |      | edges of the lowest $z$ plane;  
     |      | $z$ axis points from center of the lowest $z$ plane to  
     |      | center of the highest $z$ plane;  
     |      | $x$, $y$, $z$ axes of the LCS are parallel to those of  
     |      | the GCS. |
| 2  | CYLINDER | $C_1$ and $C_2$—inner and outer radii of the cylinder;  
     |      | $C_3$—cylinder length along $z$ axis;  
     |      | $RP$—center of the lowest $z$ plane;  
     |      | $z$ axis points from center of the lowest $z$ plane to  
     |      | center of the highest $z$ plane;  
     |      | $z$ axis of the LCS is parallel to that of the GCS;  
     |      | $x$ and $y$ axes of the LCS orientation is arbitrarily. |
| 3  | SPHERE | $C_1$ and $C_2$—inner and outer radii of the sphere;  
     |      | $RP$—sphere center;  
     |      | LCS axes oriented arbitrarily. |
| 4  | CONE  | $C_1$ and $C_2$—inner and outer radii at the lowest $z$;  
     |      | $C_3$ and $C_4$—inner and outer radii at the highest $z$;  
     |      | $C_5$—cone length along $z$ axis;  
     |      | $RP$—center of the lowest $z$ plane;  
     |      | $z$ axis points from center of the lowest $z$ plane to center  
     |      | of the highest $z$ plane;  
     |      | $z$ axis of the LCS is parallel to that of the GCS;  
     |      | $x$ and $y$ axes of the LCS orientation is arbitrarily. |
3.4 Materials, Defaults and Energy Groups

For user convenience this section includes information on the MARS built-in materials, on the default parameters and on the low-energy groups in two representations.

Table 5: Built-in materials. Material number (index) IM is assigned to each material according to the input order in the file MARS.INP. By definition IM=0 corresponds to vacuum and IM=NREMA+1 to the black hole outside the global volume.

<table>
<thead>
<tr>
<th>NAME</th>
<th>ρ(g/cm³)</th>
<th>A</th>
<th>Z</th>
<th>Material</th>
</tr>
</thead>
<tbody>
<tr>
<td>H</td>
<td>0.071</td>
<td>1.00</td>
<td>1.</td>
<td>Hydrogen</td>
</tr>
<tr>
<td>LHE</td>
<td>0.125</td>
<td>4.00</td>
<td>2.</td>
<td>Liq. Helium</td>
</tr>
<tr>
<td>LI</td>
<td>0.534</td>
<td>6.94</td>
<td>3.</td>
<td>Lithium</td>
</tr>
<tr>
<td>BE</td>
<td>1.85</td>
<td>9.01</td>
<td>4.</td>
<td>Beryllium</td>
</tr>
<tr>
<td>CH2</td>
<td>0.94</td>
<td>8.96</td>
<td>4.62</td>
<td>Polyethylene</td>
</tr>
<tr>
<td>CH</td>
<td>1.03</td>
<td>10.24</td>
<td>5.2</td>
<td>Polystyrene</td>
</tr>
<tr>
<td>C</td>
<td>2.265</td>
<td>12.01</td>
<td>6.</td>
<td>Graphite</td>
</tr>
<tr>
<td>TISS</td>
<td>0.99</td>
<td>12.37</td>
<td>6.3</td>
<td>Tissue</td>
</tr>
<tr>
<td>WATR</td>
<td>1.00</td>
<td>12.41</td>
<td>6.32</td>
<td>Water</td>
</tr>
<tr>
<td>AIR</td>
<td>0.00121</td>
<td>13.05</td>
<td>6.53</td>
<td>Air</td>
</tr>
<tr>
<td>SOIL</td>
<td>1.90</td>
<td>20.80</td>
<td>10.40</td>
<td>Soil</td>
</tr>
<tr>
<td>CONC</td>
<td>2.35</td>
<td>20.88</td>
<td>10.44</td>
<td>Concrete</td>
</tr>
<tr>
<td>AL</td>
<td>2.70</td>
<td>26.98</td>
<td>13.</td>
<td>Aluminium</td>
</tr>
<tr>
<td>SI</td>
<td>2.33</td>
<td>28.09</td>
<td>14.</td>
<td>Silicon</td>
</tr>
<tr>
<td>LAR</td>
<td>1.40</td>
<td>39.95</td>
<td>18.</td>
<td>Liq. Argon</td>
</tr>
<tr>
<td>TI</td>
<td>4.54</td>
<td>47.88</td>
<td>22.</td>
<td>Titanium</td>
</tr>
<tr>
<td>FE</td>
<td>7.86</td>
<td>55.85</td>
<td>26.</td>
<td>Iron</td>
</tr>
<tr>
<td>CU</td>
<td>8.92</td>
<td>63.55</td>
<td>29.</td>
<td>Copper</td>
</tr>
<tr>
<td>W</td>
<td>19.30</td>
<td>183.85</td>
<td>74.</td>
<td>Tungsten</td>
</tr>
<tr>
<td>Pb</td>
<td>11.35</td>
<td>207.19</td>
<td>82.</td>
<td>Lead</td>
</tr>
<tr>
<td>U</td>
<td>18.95</td>
<td>238.03</td>
<td>92.</td>
<td>Uranium</td>
</tr>
<tr>
<td>BLAl</td>
<td>10000.</td>
<td>240.00</td>
<td>94.</td>
<td>Black Hole</td>
</tr>
</tbody>
</table>
Table 6: Optional input variables in MARS.INP and their defaults.

<table>
<thead>
<tr>
<th>Keyword</th>
<th>Variables and Arrays</th>
<th>Default</th>
</tr>
</thead>
<tbody>
<tr>
<td>MTEXT</td>
<td></td>
<td>Format(A80)</td>
</tr>
<tr>
<td>INDX</td>
<td>IND(20)</td>
<td>20*F</td>
</tr>
<tr>
<td>NEVT</td>
<td>NSTOP, NTIME</td>
<td>200, 0</td>
</tr>
<tr>
<td>ENRG</td>
<td>E0, EM, PSTAM, EMI, EMNU</td>
<td>100...0145..3..005..5</td>
</tr>
<tr>
<td>IPIB</td>
<td>I0, IBEAM</td>
<td>1, 0</td>
</tr>
<tr>
<td>BEAM</td>
<td>STXX, STYY, STTX, STY</td>
<td>4*0.</td>
</tr>
<tr>
<td>INIT</td>
<td>XINI, YINI, ZINI, DXIN, DYIN, DZIN</td>
<td>5*0., 1.</td>
</tr>
<tr>
<td>SMIN</td>
<td>STEPEM, STEPH</td>
<td>0.3, 10.</td>
</tr>
<tr>
<td>VARS</td>
<td>EFF, DLEXP, TEMPO, AINT</td>
<td>0..1..300..,10^12</td>
</tr>
<tr>
<td>NMAT</td>
<td>NREMA</td>
<td>1</td>
</tr>
<tr>
<td>MATR</td>
<td>AMA(20), ROW(20), ATW(20), ZAW(20)</td>
<td>FE, 79*0</td>
</tr>
<tr>
<td>NLNG</td>
<td>LZ, NLZ</td>
<td>1, 1</td>
</tr>
<tr>
<td>ZSEC</td>
<td>ZSE(50), IZN(50), IZI(50), IZM(50)</td>
<td>100..49<em>0..,100</em>1,50*0</td>
</tr>
<tr>
<td>NLTR</td>
<td>LR</td>
<td>1</td>
</tr>
<tr>
<td>RSEC</td>
<td>RSE(20), IRN(20), IRI(20), IRM(20)</td>
<td>5.., 19<em>0.., 40</em>1, 20*0</td>
</tr>
<tr>
<td>NAZM</td>
<td>NF</td>
<td>1</td>
</tr>
<tr>
<td>AZIM</td>
<td>FIB(60)</td>
<td>60*0.</td>
</tr>
<tr>
<td>NOBL</td>
<td>NOB</td>
<td>0</td>
</tr>
<tr>
<td>RZOB</td>
<td>RZO(4,3)</td>
<td>12*0.</td>
</tr>
<tr>
<td>PLOT</td>
<td>RPLLOT, ZPLO1, ZPLO2, ZPLO3</td>
<td>4*0.</td>
</tr>
<tr>
<td>NEUS</td>
<td>NG, NGOUT, NINT, NGMAX(20)</td>
<td>28, 26, 20, 20*11</td>
</tr>
<tr>
<td>NDET</td>
<td>NDE1</td>
<td>0</td>
</tr>
<tr>
<td>FLOC</td>
<td>WJ, XD(10), YJ(10), ZD(10)</td>
<td>0.5, 30*0.</td>
</tr>
<tr>
<td>NCHO</td>
<td>NTYZ(6)</td>
<td>2, 2, 1, 3*0</td>
</tr>
<tr>
<td>FLCH</td>
<td>FL(200)</td>
<td>200*0.</td>
</tr>
<tr>
<td>MUON</td>
<td>KPHA, IDNIDX</td>
<td>30, 0</td>
</tr>
<tr>
<td>STOP</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
### Table 7: Neutron energy group numbers NG and lower energy boundaries EG (MeV) in the group at $E \leq 14.5$ MeV in the 28-group representation. Used with $\text{IND(14)} = T$ and $\text{IND(17)} = F$.

<table>
<thead>
<tr>
<th>NG</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
<th>6</th>
<th>7</th>
</tr>
</thead>
<tbody>
<tr>
<td>EG</td>
<td>14.</td>
<td>10.5</td>
<td>6.5</td>
<td>4.0</td>
<td>2.5</td>
<td>1.4</td>
<td>0.8</td>
</tr>
<tr>
<td></td>
<td>8</td>
<td>9</td>
<td>10</td>
<td>11</td>
<td>12</td>
<td>13</td>
<td>14</td>
</tr>
<tr>
<td>EG</td>
<td>0.4</td>
<td>0.2</td>
<td>0.1</td>
<td>4.65E-2</td>
<td>2.15E-2</td>
<td>1.E-2</td>
<td>4.65E-3</td>
</tr>
<tr>
<td>NG</td>
<td>15</td>
<td>16</td>
<td>17</td>
<td>18</td>
<td>19</td>
<td>20</td>
<td>21</td>
</tr>
<tr>
<td>EG</td>
<td>2.15E-3</td>
<td>1.E-3</td>
<td>4.65E-4</td>
<td>2.15E-4</td>
<td>1.E-4</td>
<td>4.65E-5</td>
<td>2.15E-5</td>
</tr>
<tr>
<td>NG</td>
<td>22</td>
<td>23</td>
<td>24</td>
<td>25</td>
<td>26</td>
<td>27</td>
<td>28</td>
</tr>
<tr>
<td>EG</td>
<td>1.E-5</td>
<td>4.65E-6</td>
<td>2.15E-6</td>
<td>1.E-6</td>
<td>4.65E-7</td>
<td>2.15E-7</td>
<td>2.15E-9</td>
</tr>
</tbody>
</table>

### Table 8: Neutron energy group numbers NG and lower energy boundaries EG (MeV) in the group at $E \leq 18$ MeV in the 49-group representation. Used with $\text{IND(14)} = T$ and $\text{IND(17)} = T$.

<table>
<thead>
<tr>
<th>NG</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
<th>6</th>
<th>7</th>
</tr>
</thead>
<tbody>
<tr>
<td>EG</td>
<td>16.90</td>
<td>15.80</td>
<td>14.80</td>
<td>14.00</td>
<td>13.03</td>
<td>12.12</td>
<td>11.28</td>
</tr>
<tr>
<td></td>
<td>8</td>
<td>9</td>
<td>10</td>
<td>11</td>
<td>12</td>
<td>13</td>
<td>14</td>
</tr>
<tr>
<td>EG</td>
<td>10.50</td>
<td>9.312</td>
<td>8.260</td>
<td>7.327</td>
<td>6.500</td>
<td>5.757</td>
<td>5.099</td>
</tr>
<tr>
<td>NG</td>
<td>15</td>
<td>16</td>
<td>17</td>
<td>18</td>
<td>19</td>
<td>20</td>
<td>21</td>
</tr>
<tr>
<td>EG</td>
<td>4.516</td>
<td>4.000</td>
<td>3.640</td>
<td>3.320</td>
<td>2.881</td>
<td>2.500</td>
<td>2.162</td>
</tr>
<tr>
<td>NG</td>
<td>22</td>
<td>23</td>
<td>24</td>
<td>25</td>
<td>26</td>
<td>27</td>
<td>28</td>
</tr>
<tr>
<td>EG</td>
<td>1.870</td>
<td>1.618</td>
<td>1.400</td>
<td>1.210</td>
<td>1.058</td>
<td>0.920</td>
<td>0.800</td>
</tr>
<tr>
<td>NG</td>
<td>29</td>
<td>30</td>
<td>31</td>
<td>32</td>
<td>33</td>
<td>34</td>
<td>35</td>
</tr>
<tr>
<td>EG</td>
<td>0.400</td>
<td>0.200</td>
<td>0.100</td>
<td>4.65E-2</td>
<td>2.15E-2</td>
<td>1.00E-2</td>
<td>4.65E-3</td>
</tr>
<tr>
<td>NG</td>
<td>36</td>
<td>37</td>
<td>38</td>
<td>39</td>
<td>40</td>
<td>41</td>
<td>42</td>
</tr>
<tr>
<td>EG</td>
<td>2.15E-3</td>
<td>1.00E-3</td>
<td>4.65E-4</td>
<td>2.15E-4</td>
<td>1.00E-4</td>
<td>4.65E-5</td>
<td>2.15E-5</td>
</tr>
<tr>
<td>NG</td>
<td>43</td>
<td>44</td>
<td>45</td>
<td>46</td>
<td>47</td>
<td>48</td>
<td>49</td>
</tr>
<tr>
<td>EG</td>
<td>1.00E-5</td>
<td>4.65E-6</td>
<td>2.15E-6</td>
<td>1.00E-6</td>
<td>4.65E-7</td>
<td>2.15E-7</td>
<td>2.53E-8</td>
</tr>
</tbody>
</table>

### Table 9: Photon energy group numbers NG and lower energy boundaries EG (MeV) in the group at $E \leq 11$ MeV. Used with $\text{IND(14)} = T$ and $\text{IND(17)} = T$.

<table>
<thead>
<tr>
<th>NG</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
<th>6</th>
<th>7</th>
<th>8</th>
</tr>
</thead>
<tbody>
<tr>
<td>EG</td>
<td>11.</td>
<td>9.</td>
<td>7.</td>
<td>5.5</td>
<td>4.5</td>
<td>3.5</td>
<td>2.5</td>
<td>1.75</td>
</tr>
<tr>
<td>NG</td>
<td>9</td>
<td>10</td>
<td>11</td>
<td>12</td>
<td>13</td>
<td>14</td>
<td>15</td>
<td></td>
</tr>
<tr>
<td>EG</td>
<td>1.25</td>
<td>0.75</td>
<td>0.35</td>
<td>0.15</td>
<td>0.08</td>
<td>0.04</td>
<td>0.02</td>
<td>0.01</td>
</tr>
</tbody>
</table>
3.5 Input Examples

**Example 1.** Calculate antiproton production with 9-cm long 1-cm diameter copper target irradiated with 120-GeV proton beam. Beam R.M.S spot size is $\sigma_x=0.005$ cm and $\sigma_y=0.007$ cm. In addition to forced antiproton production the user is interested in energy deposition calculation including knock-on electron and $e^+e^-$-pair production by hadrons. To create a file of antiprotons generated on the target in a given phase space, the user adds in **MAIN**:

```
OPEN (UNIT=9,FILE='PBAR.OUT',STATUS='UNKNOWN')
```

and a few statements in the **LEAK** routine (see Sect. 4.5). The geometry and scoring is described in the **standard** mode. The **MARS.INP** file can look as:

```
Pbar Cu Target, sigx=0.005, 02-Nov-1994
INDX T 5=T 12=T
NEVT 100000
ENRG 120.
IPIB 1 2
BEAM 0.005 0.007
SMIN 0.001 3.
MATR 'CU'
ZSEC 9. 51=3
NLTR 5
RSEC 0.002 0.005 0.01 0.1 .5
STOP
```

The geometry is very simple here, completely adequate to the **standard** mode. Instead of the above bining and writing to a file PBAR.OUT, one can use HBOOK (see **MAIN**) for analyses of energy deposition and generated antiprotons. Then the solid target can be described just as:

```
ZSEC 9.
RSEC 0.5
```

**Example 2.** Intense 800-GeV proton beam hits a graphite dump followed by aluminum and steel absorbers, by 210-cm vacuum gap, by 5-cm polyethylene slab, and finally by 1 meter of a wet dirt. The user is interested in energy deposition calculations, in maximum amount of the output, including partial distributions, temperature rise and estimation of a residual dose rate, in intermediate result dumps. A compound material is defined in the **MIXTUR** routine (Sect. 4.1). The routines
REG1 and REG2 (Sect. 4.3) define the dump as surrounded with a steel shield sitting at the axis of a hypothetical cylindrically symmetrical tunnel with 30-cm thick concrete walls. With these three user routines a standard MARS.INP file can look as (Option 1):

```
Tevatron C0 Dump, 02/28/95
INDX T T 8=T
NEVT 500000 5
IPIB 1 2
ENRG 800.
SMIN 0.01 8.
BEAM 0.0416 0.0944
VARS 0. 1. 300. 4.E13
NMAT 6
MATR 'C' 'AL' 'FE' 'CONC' 'CH2' 'MIXT'
26=2.1
NLNG 7
51=14 5 10 1 3 2 5 101=1 2 3 2 0 4 5
NLTR 9
RSEC 0.03 0.1 0.3 1. 3.5 10. 30. 150. 180.
PLOT 5. 300. 450. 805.
STOP
```

Option 2. The geometry description is simpler and the output is more sophisticated if one uses HBOOK activated in MAIN. The last lines in MARS.INP can be re-written as:

```
NLNG 7
101=1 2 3 2 0 4 5
NLTR 4
RSEC 10. 30. 150. 180.
```

Option 3. It is worthwhile to use extended geometry description in this example. The following changes have to be done to the MARS.INP file:

```
INDX T T T 8=T
NLNG 1
ZSEC 805.
NLTR 1
RSEC 180.
```
A GEOM.INP file can look as:

```
12 1
2 2 0 0 0. 0. 0. 0. 0. 180. 805.
3 2 1 0 0. 0. 0. 0. 0. 10. 350.
4 2 2 0 0. 0. 350. 0. 0. 10. 75.
5 2 3 0 0. 0. 425. 0. 0. 10. 50.
6 2 2 0 0. 0. 475. 0. 0. 10. 15.
7 2 0 0 0. 0. 490. 0. 0. 150. 210.
8 2 5 0 0. 0. 700. 0. 0. 150. 5.
9 2 6 0 0. 0. 705. 0. 0. 150. 100.
10 2 3 0 0. 0. 0. 10. 30. 490.
11 2 0 0 0. 0. 0. 30. 150. 490.
12 2 4 0 0. 0. 0. 150. 180. 805.
```

User subroutines REG1 and REG2 are not needed in this case. Detailed histogramming can be done via HBOOK (as in Option 2) or by adding a required number of regions into the files MARS.INP or GEOM.INP. Depending on the application an appropriate combination of the standard and extended volumes, of the standard and HBOOK histogrammings etc., can improve the outcome.

### 4 User Subroutines

All the user subroutines described below are collected in the file m13.f and are called if IND (3) = T. By default they are the dummy and must be the dummy until the user changes them for his/her particular application.

#### 4.1 Compounds (MIXTUR)

The user can define any composite material not built into the current version of MARS (see Table 5) with the help of a subroutine MIXTUR (I, M, A, Z, W). For each non-standard material marked in the input sequence with 'MIXT' this subroutine has to define number of components in the compound M, atomic masses A(M), atomic numbers Z(M) and relative fractions of components W(M) in the compound. Material index I is determined according to the input order of material cards MATR. An example of the subroutine (used in the second example of Sect. 3.5) for the case when the first five materials are the pre-defined elements (from Table 5 or described on the input card MATR), and the sixth is SiO2 with 5% of water by weight, is presented below:
SUBROUTINE MIXTUR(I,M,A,Z,W)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C INPUT: I-MATERIAL INDEX. OUTPUT:
C M - NUMBER OF COMPONENTS, 2≤M≤6
C A, Z - ATOMIC MASSES AND NUMBERS OF COMPONENTS
C W - RELATIVE FRACTIONS OF COMPONENTS
DIMENSION A(l),Z(l),W(l)
IF(I.LT.6) RETURN
M=3
A(1)=28.
A(2)=16.
A(3)=1.
Z(1)=14.
Z(2)=8.
Z(3)=1.
S=0.95*(28.+32.)+0.05*18.
W(1)=0.95*28./S
W(2)=(0.95*32+0.05*16.)/S
W(3)=0.05*2./S
RETURN
END

4.2 Source (BEG1)

A source term of almost any complexity can be described in a user-written subroutine BEG1 (JJ, W, E, X, Y, Z, DX, DY, DZ). This is especially useful when the problem should be solved in two (or more) consecutive steps. BEG1 is used for more efficient calculations in the following cases:

- separate event generator (DTUJET, ISAJET etc.);
- very extended source (e.g., beam loss distributions in accelerators, coupling with the STRUCT code);
- geometrically and energetically extended source (e.g., thermal neutron flux studies in the multi-section labyrinths at accelerators);
- any other pre-calculated source;
- tagging of certain regions in the source particle phase-space.
Any part of an external source can be read. Each call to a subroutine BEG1 must provide a new set of initial particle parameters. In this subroutine the user may re-define any or all of 9 parameters in case he/she wants them to be different from the ones defined in the input sequence. The parameters are particle type (JJ), statistical weight of the current event ($W$ with a default $W=1$), kinetic energy of initial particle ($E$), initial coordinates ($X$, $Y$, $Z$) and initial direction cosines ($DX$, $DY$, $DZ$). Before each call to a subroutine BEG1 (once per event) all 9 parameters are equal to input ones and any can be changed inside.

For example, to have incident particles distributed uniformly on a 3.3-cm aperture and along a 600-cm region longitudinally, the user should provide the following subroutine:

```fortran
SUBROUTINE BEG1(JJ,W,E,X,Y,Z,DX,DY,DZ)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C RE-DEFINES EACH OR ANY OF THE 9 PARAMETERS
C OF INITIAL SOURCE PARTICLES
C PARTICLE TAGGING IN 'MTAG' SOURCE ZONES OF 'EG'
C ENERGY INTERVALS FOR 'NTAG' DETECTOR ZONES
C DEFAULTS: MTAG=0, INTG=1, IETG=4
C
COMMON/MATINT/IM
COMMON/TAGG/ET(60,4,6),EG(4),NTAG(6),MTAG,INTG,IETG
DATA R,PI/3.5,3.1415926/
C
A=2.*PI*RNDM(-1.)
X=R*SIN(A)
Y=R*COS(A)
Z=600.*RNDM(-1.)
RETURN
END
```

4.3 Geometries (REG1, REG2, REG3)

If one wants to study cascades in a very complex geometry not embraced by the above options, user subroutines REG1 and REG2 must be provided. MARS, version 13(95), allows the user to place geometrical objects of almost any complexity inside the pre-defined standard ($r$-$z$-$\phi$) or extended geometry. By convention the total number of standard regions $NFZP$ must be $NFZP \geq 1$. With the help of his/her own subroutines REG1 ($X$, $Y$, $Z$, $N$, $NIM$) and REG2 ($N$, $IM$, $MAG$) the
user can describe arbitrary physical regions numbered from the $N_{\text{MIN}} \leq N \leq N_{\text{MAX}}$ interval. Here $N_{\text{MAX}}=10000$; $N_{\text{MIN}}=NFZP+1$ for the standard geometry sector and $N_{\text{MIN}}=NEXGM+NVOLUM$ for the extended geometry sector. Each region can be divided into any number of arbitrary subregions with index $N_{\text{IM}}$ with a default value $N_{\text{IM}}=0$. This feature provides the possibility of distinguishing geometrical zones without scoring results there. The user can successfully substitute this with the HBOOK histogramming (see MAIN and m13hist.f).

For each call, a subroutine $\text{REG1}(X, Y, Z, N, N_{\text{IM}})$ finds the position of the given point $(X, Y, Z)$ in the system: it determines the corresponding physical region number $N$ (each with its own material index) and, if one wishes, the subregion number $N_{\text{IM}}$. The second subroutine $\text{REG2}(N, IM, MAG)$ attributes to a given $N (\geq N_{\text{MIN}})$ the material index $IM$ and, if $\text{IND}(4)=T$, the magnetic index $MAG$. The last parameter can be used in a subroutine $\text{FIELD}$ to determine the type of magnetic field (uniform, dipole, quadrupole etc.) in the region $N$. Default: $MAG=0$ (no magnetic field in the region).

By convention the region outside of the global volume has a number $N=0$ and properties of the black hole. In some applications it is useful to tag and to score the leakage out of the non-standard regions into the black holes labeled with $N \leq -1$ to use these negative tags in a user routine $\text{LEAK}$. The user must pay a special attention to careful programming of the subroutine $\text{REG1}$, because the geometrical modules consume usually about 80% of the CPU time, as is typical of cascade Monte Carlo programs.

Simple examples of the subroutines $\text{REG1}$ and $\text{REG2}$ (used in the second example of Section 3.5) are:
SUBROUTINE REG1(X,Y,Z,N,NIM)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C NON-STANDARD GEOMETRY MODULE
C FINDS THE PLACE OF GIVEN POINT IN THE SYSTEM
C INPUT: X, Y, Z
C OUTPUT:
C N - PHYSICAL REGION NUMBER, NMIN≤N≤NMAX
C NMIN=NFZP+1 (STANDARD)
C OR
C NMIN=NEXGM+NVOLUM (EXTENDED)
C NMAX=10000
C N<1 DEFINES NUMBERED LEAKAGE OUT OF THE SYSTEM
C IN NON-STANDARD SECTOR
C NIM - GEOMETRICAL SUBREGION NUMBER, 0≤NIM
C
C REVISION: 28-FEB-1995
C
R=SQRT(X*X+Y*Y)
IF(R.LT.10.) RETURN
M=500
IF(R.LT.150.) THEN
IF(Z.LT.490.) THEN
STEEL SHIELDING:
N=M+1
C TUNNEL:
IF(R.GT.30.) N=M+2
END IF
ELSE
C CONCRETE SHELL:
N=M+3
END IF
RETURN
END
SUBROUTINE REG2 (N, IM, MAG)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C FINDS MATERIAL AND MAGNETIC INDICES
C INPUT: N - PHYSICAL REGION NUMBER
C OUTPUT:
C IM - MATERIAL INDEX
C MAG - MAGNETIC FIELD INDEX, IF IND(4)=T
C
C REVISION: 28-FEB-1995
C IM=1(C), 2(AL), 3(FE), 4(CONC), 5(CH2), 6(DIRT)
C
DIMENSION IMUN(3)
DATA IMUN/1, 0, 6/
M=N-500
IM=IMUN(M)
RETURN
END

In some cases it is convenient to re-define a few material and/or magnetic indices assigned to the standard or extended regions at the initialization stage. This can be easily done in a user routine REG3, e.g.:

SUBROUTINE REG3 (N, IM, MAG)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C RE-DEFINES IM AND MAG FOR STANDARD SECTOR
IF(N.EQ.137) IM=14
IF(N.EQ.2503) IM=2
RETURN
END

4.4 Magnetic and Electrical Fields (FIELD, SUFI)

To describe the magnetic field components (BX, BY, BZ) in the regions with parameter MAG≠0 the user puts IND(4)=T and provides a subroutine FIELD. The same routine is used to describe an electrical field. One can use a corresponding map or analytical expressions to find the field components in the point (X, Y, Z). Parameter MAG defined in REG2 can speed up the search. It indicates the type of the field in the region. The unit for magnetic field is Tesla. A 2-D or 3-D field map is read in a user routine SUFI at the initialization stage and transferred to the routine FIELD via appropriate COMMON block.
An example of the FIELD subroutine is:

```fortran
SUBROUTINE FIELD(X,Y,Z,BX,BY,BZ,BBB)
  IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
  COMMON/BLINT2/JJ,KK,MAG
  DATA PAQ,BQUA,G1/8.,2.,0.53/
  BX=0.
  RY=0.
  BZ=0.
  IF(Z.LT.622.) THEN
    BX=5.
  IF(MAG.EQ.2) BX=-5.
  ELSE
    CALL QUAD(X,Y,R,RAQ,G1,BX,BY)
    CALL DIPOLE(X,Y,R,RAQ,BQUA,BX1,BY1)
    BX=BX+BX1
    BY=BY+BY1
  END IF
  RETURN
END
```

4.5 Leakage (LEAK)

A user subroutine LEAK (N, K, JJ, W, E, X, Y, Z, DX, DY, DZ) handles particles which escape from the system (N=0) or from the non-standard regions tagged with N≤−1. Parameter K is the tree vertex level (generation number) for the hadronic part of the calculated cascade. If IND (14) =T, then K has the same meaning as the above if given particle was generated in hA vertex at E>0.0145 GeV, otherwise it is forced to be K=−205, if the leaked neutron was produced in sequent interactions below 0.0145 GeV. JJ is particle type from Table 1. If K=−205, the
particle type is neutron by default and $JJ$ is its energy group number from Tables 7 or 8. $W$ is leaked particle statistical weight, $E$ is its kinetic energy, $(X, Y, Z)$ and $(DX, DY, DZ)$ - its coordinates and direction cosines, respectively.

The routine LEAK used in the first example of Sect. 3.5 to collect generated antiprotons in the file PBAR.OUT looks like:

```fortran
SUBROUTINE LEAK(N, K, JJ, W, E, X, Y, Z, DX, DY, DZ)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C PARTICLE LEAKAGE SPECIAL SCORING
C IF(K=-205): JJ=NGROUP, E=0.0005*(EN(JJ)+EN(JJ+1))
C FOR L.E.NEUTRONS (E<0.0145 GEV)
C REVISION: 02-NOV-1994
COMMON/STAZIU/ZMAX,REXT
DATA POPT,DPO,DZ0/8.9,0.02,2,0.99/
C
IF(JJ.NE.12) RETURN
IF(DZ.LE.DZ0) RETURN
P=SQRT(E*(E+1.87656))
DP=ABS((P-POPT)/POPT)
IF(DP.GT.DP0) RETURN
WRITE(9)E,W,X,Y,Z,DX,DY,DZ,K
RETURN
END
```

4.6 Fictitious Scattering (ALIGN, SAGIT)

In some applications components of the considered system can be turned or shifted with respect to each other. Say, the arc of any circular accelerator is built of the magnets turned by a fixed angle. MARS, version 13(95), allows an elegant way to handle such systems. The user describes the geometry as the straight along the $z$-axis. Then, with the help of a routine ALIGN, he/she creates the angular or space kicks at the required boundaries in the directions opposite to the real ones. One can easily see that the resulting coordinates and angles of any particle are identical to those in the real bent geometry. It is convenient to use spare NIM parameters to control such “a fictitious scattering” at the required boundaries with NIB≠NIM as that parameter in a previous region. Of course, any coordinate can be used to locate the needed boundary.

The following example shows use of the ALIGN routine for a single kick:
SUBROUTINE ALIGN(X,Y,Z,DX,DY,DZ,NIM,NIB)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C FICTITIOUS SCATTERING : DISCRETE AT BOUNDARIES
C RE-DEFINES EACH OR ANY OF THE FIRST 6 PARAMETERS
C
DATA ANG/1.5E-3/
IF(NIB.EQ.1050.AND.NIM.EQ.2010) THEN
  DY=DY-ANG
END IF
RETURN
END

The same method is used in MARS to handle objects with saggita, i.e. continuously bent. "Fictitious scattering" defined in a user subroutine SAGIT as an anti-kick at every step along charged and neutral particle trajectories, simplifies the geometry description and makes a precise correspondence to the real bent objects.

4.7 Edge Scattering (EDGEUS)

If the "edge scattering problem" is considered with IND(9) = TRUE and the geometry includes non-standard insertions defined with REG1 and REG2, the user should supply a subroutine EDGEUS(X,Y,Z,DX,DY,DZ,U,V). It finds the distance U from the given trajectory point (X, Y, Z) to a non-standard surface, and the projection V of the direction vector (DX, DY, DZ) to the normal to the surface which passes in its positive direction through the point (X, Y, Z).

This example shows use of the EDGEUS routine for the surface Y = YPL > 0 and for the particle with Y > YPL and DY < 0:

SUBROUTINE EDGEUS(X,Y,Z,DX,DY,DZ,U,V)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C EDGE-SCATTERING PROBLEM
U=Y-YPLANE
V=DY
RETURN
END
4.8 Special Volumes (VFAN)

In a standard geometry sector the volumes of the regions needed, for example, to compute energy deposition density in GeV/g, are calculated at the final stage in the SERVN subroutine. The array of volumes $V_V$ of the non-standard regions $M$ should be provided by the user in a subroutine VFAN ($N, V$). The same subroutine can be used in specific activity and residual dose rate calculations at IND (13) = T. The VFAN routine can look as:

```fortran
SUBROUTINE VFAN(N,V)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C FINDS VOLUME V, CM**3 OF REGION N>NFZP
DIMENSION VV(2)
DATA VV/1.53E4,7.34E6/
M=N-750
V=VV(M)
RETURN
END
```

4.9 Intermediate Dumps (DUMP)

A parameter $NTIME \geq 1$ in the input sequence of MARS.INP creates dumps of the intermediate results $NTIME$ times. The user defines in the subroutine DUMP (NDUMP) what to display on screen and to store in the file DUMP:

```fortran
SUBROUTINE DUMP(NDUMP)
IMPLICIT REAL (A-H,O-Z), INTEGER (I-N)
C DUMP OUTPUT NTIME \geq 1 TIMES, NUPRI=NSTOP/NTIME
NDUMP = 14
OPEN (NDUMP,FILE='DUMP',STATUS='UNKNOWN')
CALL DATIMH(MDATE,MTIME)
WRITE(G)MDATE,MTIME,NI,NSTOP,NUPRI,NDUMP
WRITE(NDUMP)MDATE,MTIME,NI,NSTOP,NUPRI,NDUMP
C WRITE NEEDED RESULTS TO 'DUMP' FILE HERE:
C
CLOSE (NDUMP)
RETURN
END
```
5 Output

There are three levels of the MARS, version 13(95), output which are essentially self-explainatory:

- general output in the file MARS.OUT;
- a file MARS.HIST defined in MAIN and filled by HBOOK for further analyses with PAW and other tools;
- a few files to be used by graphics packages PAW and others (.GRA, .PLOT) and to be used in a consecutive run in a multi-stage case; the files are defined in the MAIN subroutine (see Sect. 2.4).

5.1 General Output (MARS.OUT)

A file MARS.OUT consists of two major sections: input status and printout of the Monte Carlo session results.

A. Input Status

The subroutine BEGINN prints the cards it has read and some calculated quantities to be used in the Monte Carlo session. Every value is provided with the corresponding keyword and title. The following is printed consecutively: code version number, date and time, title of the problem, options logical statement, requested numbers of incidents (events) and of intermediate dumps, incident particle and beam types, incident kinetic energy, particle fluence cutoff energies, star production threshold, incident particle type, initial coordinates and direction cosines of the beam spot center, R.M.S. beam spot sizes and angular spread, accuracy of boundary localization STEPEM and control parameter STEPH, point-like target efficiency, exponential conversion factor, initial temperature, beam intensity, number of hA generations to follow.

Then for each material the following values that have been either read in or calculated are printed: material index, material name, averaged atomic mass and atomic number, averaged electron density, ionization potential, critical energy and radiation length, contents of composite materials (if any), threshold energies for δ-electrons and direct e⁺e⁻ pair production by charged hadrons, calculated lengths for inelastic hadron-nuclear interactions at incident energy E₀ and ionization ranges at the cutoff energy EM.
If $\text{IND}(14) = \text{T}$ information on low-energy neutrons is printed: content of the NEUS card, group cross-sections for each material, number of point-like detectors and their coordinates (if any).

If $\text{IND}(11) = \text{T}$ there is a printout of the azimuthal grid and if $\text{NOB} \geq 1$ a printout of radial and longitudinal boundaries of macro-regions for particle spectra scoring.

If all of $\text{RPL0T}$, $\text{ZPLO1}$, $\text{ZPLO2}$, $\text{ZPLO3}$ are not equal to zero, two lateral and one longitudinal cross-sectional views will be drawn.

**B. Printout**

First, for the standard $(r, z, \phi)$ part of the considered geometry, a table is printed which indicates correspondence of region boundaries and index $N$, assigned to each region. Also printed are corresponding material and magnetic field indices. The calculational results are referenced to this table by a region number $N$.

The MARS.OUT file contains the results of the Monte Carlo session as a number of tables and single quantities. All results are normalized to one incident particle (to one event), only temperature rise and residual dose distributions are normalized to $\text{AINT}$ incidents (events). The following tables are printed here consecutively:

1. Longitudinally integrated lateral distributions of charged and total star density and hadron flux, of partial and total energy deposition.

2. Laterally integrated longitudinal distributions of partial and total energy deposition, corresponding cumulative distribution and total energy deposited via the different channels: neutron reactions below $EM$, low-energy particles from the nuclear de-excitation processes, electromagnetic showers and ionization losses of charged hadrons and muons.

3. Three-dimensional star density distribution (*stars per cubic centimeter*) induced by charged hadrons and by all hadrons for momenta $\geq PSTM$ with corresponding statistical errors (one R.M.S.). Collision estimator.

4. Three-dimensional total hadron flux distribution (*particles per centimeter squared*) for kinetic energy $\geq EM$ with corresponding statistical errors. Track-length estimator.

5. Three-dimensional charged hadron flux distributions for kinetic energy above two thresholds, 0.0145 and 0.5 GeV as a default.
6. Three-dimensional $e^+e^-$ flux distributions for kinetic energy above two thresholds, 0.001 and 0.005 GeV as a default, if $\text{IND}(1)=T$.

7. Three-dimensional photon flux distributions for kinetic energy above two thresholds, 0.001 and 0.005 GeV as a default, if $\text{IND}(1)=T$.

8. Three-dimensional muon flux distributions for kinetic energy above two thresholds, 0.005 and 0.5 GeV as a default, if $\text{IND}(10)=T$.

9. Three-dimensional neutron ($E<0.0145$ GeV) flux distribution (particles per centimeter squared) with corresponding statistical errors, if $\text{IND}(14)=T$.

10. Total number of stars produced in the system.

11. Laterally integrated longitudinal distributions of star density, of total hadron flux, and of charged hadron, $e^+e^-$ and muon fluxes.

12. Three-dimensional distribution of energy deposition density brought by low-energy particles from de-excitation of nuclei, in GeV/g, if $\text{IND}(1)=T$.

13. Three-dimensional distribution of energy deposition density of electromagnetic showers produced by $\pi^0$ decays, by high energy $\delta$-rays and by prompt $e^+e^-$ pairs from hadrons, in GeV/g, if $\text{IND}(1)=T$.

14. Three-dimensional distribution of energy deposition density from hadron and muon electromagnetic losses with limited energy transfer, in GeV/g, if $\text{IND}(1)=T$.

15. Three-dimensional distribution of energy deposition density due to neutron interactions below $E<0.0145$ GeV, in GeV/g, if $\text{IND}(14)=T$.

16. Three-dimensional distribution of total energy deposition density with corresponding statistical errors (one R.M.S.), in GeV/g, if $\text{IND}(1)=T$.

17. Three-dimensional distribution of dose equivalent, in Rem, if $\text{IND}(1)=T$.

18. Crude estimation of three-dimensional distribution of residual dose rate, in Rad/hr, after 30-day irradiation at the mean beam intensity $A_{\text{INT}}$ particles per sec and 1-day cooling. These are valid only for sufficiently thick systems, beyond a lateral thickness of $\geq \lambda_{\text{in}}$. 

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19. Three-dimensional distribution of instantaneous temperature rise at given initial temperature \( T_0 = \text{TEMP} \) and number of particles per a single beam pulse \( N_0 = A\text{INT} \), if \( \text{IND (1)} = \text{T} \).

20. Longitudinal distribution of relative energy deposition by charged particles falling below the thresholds. Two distributions presented if \( \text{IND (1)} = \text{T} \), are for the first smallest radial bin and for the rest of the system.

21. Three-dimensional distribution of the relative statistical errors of star densities, fluxes and energy deposition densities.

22. Leakage data: number and energy of albedo hadrons, of punchthrough hadrons and of hadrons that escaped the sides of the system; leakage energy of low-energy neutrons and of electromagnetic showers; total leakage energy; energy balance.

23. Leakage energy spectra of different particles for the upstream plane, downstream plane and for the rest of the system.

24. If \( \text{NOB} \geq 1 \), energy spectra of different particles in the pre-determined special regions.

25. Tables of particle spectra, of star density, of hadron, muon and low-energy fluence, of total energy deposition density, and of temperature rise, ready for use with graphics packages. The same tables are saved in separate .GRA files if activated in \text{MAIN}.

If \( \text{IND (8)} = \text{F} \), the tables No. 5-7, 12-14, 17, 18, 20, 21 will be absent in the output.

If \( \text{IND (18)} = \text{T} \), most of the above values will be printed in a compact form for the extended geometry regions as with \( \text{IND (8)} = \text{F} \).

If the user defines geometrically complex insertions defined with \text{REG1} and \text{REG2} routines as a supplement to the standard and extended geometries, additionally the program prints in a compact form most of the above quantities with corresponding statistical errors for all non-standard regions.
5.2 Histograms

In many cases it is worthwhile to initialize in the MAIN program the HBOOK package to use the output file MARS.HIST for interactive analysis with the PAW system. Then all the powerful features of that system as described in [56] can be used for comprehensive physics analysis of the run session. By default MARS sorts histograms by following classes:

- histogram type (vertex, fluence, energy deposition and energy spectrum):
- particle class (hadron, electromagnetic and muon);
- charge (neutrals, charged and total).

These histograms are filled for the whole system or for the NOB special regions defined in the MARS.INP file. The histogram list in the file MARS.HIST with their IDs and functional contents is obtained with the PAW command hi/ list. The list is self-explainatory. Default histogram IDs are shown in Table 10. A few examples of corresponding analyses are given in the next section.

Table 10 : Default histogram ID at NOB=1.

<table>
<thead>
<tr>
<th>Particle Type</th>
<th>Vertex</th>
<th>Fluence</th>
<th>Energy Dep.</th>
<th>Spectrum</th>
</tr>
</thead>
<tbody>
<tr>
<td>n</td>
<td>1</td>
<td>101</td>
<td></td>
<td>301</td>
</tr>
<tr>
<td>h±</td>
<td>2</td>
<td>102</td>
<td></td>
<td>302</td>
</tr>
<tr>
<td>Total h</td>
<td>3</td>
<td>103</td>
<td>203</td>
<td></td>
</tr>
<tr>
<td>γ</td>
<td>4</td>
<td>104</td>
<td></td>
<td>304</td>
</tr>
<tr>
<td>e±</td>
<td>5</td>
<td>105</td>
<td></td>
<td>305</td>
</tr>
<tr>
<td>Total EM</td>
<td>6</td>
<td>106</td>
<td>206</td>
<td></td>
</tr>
<tr>
<td>μ±</td>
<td>8</td>
<td>108</td>
<td>208</td>
<td>308</td>
</tr>
</tbody>
</table>
6 Physics Analyses: Examples

Results of MARS.OUT can be used as a final product. In addition, the created files .GRA and .PLOT are ready for acceptance by the popular graphics packages PAW, TOPDRAWER, GNUPLOT, XMGR and KALEIDAGRAPH to create high resolution plots for most of the distributions listed in the previous section. Figure 1 is a typical plot obtained via such an interface. The figure shows energy deposition rate in the C-layer of the SAMUS/WAMUS muon spectrometer of the D0 detector at Fermilab for two shielding configurations. Results are obtained with 4000 DTUJET93 events (0.9×0.9 TeV $p\bar{p}$ collisions).

![Figure 1](image)

**FIGURE 1.** Dose rate in the forward muon system of the D0 detector at Fermilab.

Figure 2 (histogram ID=203) represents energy deposition density in the forward region of the D0 detector at Fermilab. Interactions of particles, produced in the $p\bar{p}$ collisions, with detector and accelerator components in this region (4 to 10 meters from the collision point) are the major source of backgrounds in the forward muon spectrometer. All the details of geometry and magnetic field were taken into account in MARS coupled here with DTUJET93. The figure highlights the hottest objects and immediately indicates if the levels in the chamber are above the tolerable one. A similar plot for particle fluxes identifies the channels for radiation to reach critical regions and helps to find an appropriate absorber to plug them.
Recently a first-pass study [45] showed that the electromagnetic component of the backgrounds from $\mu \rightarrow e\nu\bar{\nu}$ decays has the potential of killing the very attractive concept of a high-energy high-luminosity $\mu^+\mu^-$ collider unless there is significant suppression via various shielding and collimators in the detector vicinity. All simulations in the whole 80-m long inner triplet and in the detector were done with MARS, version 13(95). The simplified detector geometry used in the calculations is shown in Fig. 3.
In this application, MARS, allows an elegant way to handle the following processes:

- forced $\mu \rightarrow e\nu\bar{\nu}$ decays in the beam pipe;
- tracking of created electrons in the beam pipe under influence of the magnetic field with emission of synchrotron photons along the track;
- simulation of electromagnetic showers in the triplet and detector components induced by electrons and synchrotron photons hitting the beam pipe;
- simulation of muon interactions (bremsstrahlung, direct $e^+e^-$ pair production, ionization, deep inelastic nuclear interactions) and decays along the tracks in the lattice and detector;
- simulation of electromagnetic showers in the triplet and detector components created at the above muon interaction vertices;
- histogramming and analysis of particle energy spectra, fluences and energy deposition in various detector regions as well as in the whole IR.
Figure 4 shows calculated $e^+e^-$ energy spectrum in the accelerator components in the vicinity of the detector. The huge peak sitting around 1 TeV represents the $\mu \rightarrow e\nu\bar{\nu}$ decay spectrum with a tail at lower energies enriched by electrons and positrons of electromagnetic showers induced in the beam pipe and superconducting coils. Photons emitted due to synchrotron radiation along $e^+e^-$ tracks in a strong 8 T magnetic field have an average energy around 1 GeV. The number of photons is about 300 times that for electrons and positrons.

![Figure 4](image.png)

**FIGURE 4.** Photon (left) and electron/positron (right) energy spectra in the inner triplet accelerator components.

References for numerous comparisons of MARS code predictions with data and with other codes are given in Section 2.1. A few rather interesting recent comparisons are presented below. Figure 5 shows attenuation of a dose rate along the 30-inch waveguide conduit adjacent to a Linac tunnel where an accidental beam loss (70 MeV, $6 \times 10^{13}$ protons per sec) takes place. The results obtained with MARS and LAHET [58] are in a pretty good agreement. It is worthwhile to notice that the LAHET code has one of the best physics model for the intermediate energy range (20 to 600 MeV). An arsenal of non-analog techniques and a very detailed geometry description are necessary in these calculations making the agreement even more remarkable.
Another example is a recent intercomparison of four popular codes for high energy hadronic cascades studies: FLUKA [13], LAHET [58], GCALOR (marriage of GEANT and CALOR codes) [59] and MARS. A simple configuration was chosen: 100-GeV proton beam irradiating a 3-m long iron cylinder ($\rho = 7.4\text{g/cm}^3$) with a radius of 50 cm followed by 200 $\text{g/cm}^2$ of ordinary concrete ($\rho = 2.35\text{g/cm}^3$). Table 11 gives the calculated total number of neutrons leaked through the front, side and back surfaces. Corresponding leakage neutron spectra calculated with FLUKA using its most sophisticated physics model and with MARS are also in a rather good agreement (Fig. 6).

Table 11: Neutron leakage out of iron cylinder with a concrete shell.

<table>
<thead>
<tr>
<th>Code</th>
<th>Front</th>
<th>Side</th>
<th>Back</th>
</tr>
</thead>
<tbody>
<tr>
<td>MARS</td>
<td>153</td>
<td>4.2</td>
<td>0.96</td>
</tr>
<tr>
<td>FLUKA</td>
<td>118</td>
<td>3.6</td>
<td>1.02</td>
</tr>
<tr>
<td>LAHET</td>
<td>176</td>
<td>5.3</td>
<td>0.67</td>
</tr>
<tr>
<td>GCALOR</td>
<td>163</td>
<td>5.7</td>
<td>0.97</td>
</tr>
</tbody>
</table>

FIGURE 5. Dose rate along the Linac waveguides calculated with the LAHET (diamonds) and MARS (circles) codes.
Comparisons for cascades induced by 20-TeV proton beam loss along the superconducting magnets in the SSC tunnel are shown in Fig. 7. Neutron spectra in the tunnel cross-section, obtained with the FLUKA and MARS codes, agree very well in the energy range spanning 10 decades. MARS calculations are also in a good agreement with GCALOR and FLUKA92 predictions in designs of shield configurations to reduce the background rates in the GEM detector at the SSC [34] and in the CMS muon system for the LHC project [60].

It is worthwhile to notice a recent remarkable achievement with the D0 detector at Fermilab [54]. A system to suppress backgrounds in the forward muon spectrometer, designed with the help of MARS coupled with the STRUCT code, has been installed in the vicinity of the experimental hall and has provided substantial reduction of the accelerator related particle fluxes in the detector, in an excellent agreement with MARS predictions.
FIGURE 7. Neutron energy spectra (cm$^{-2}$ per 10$^4$ p/m) in the SSC tunnel for the routine 20-TeV proton beam loss in the superconducting magnets, calculated with the MARS (histogram) and FLUKA92 (diamonds) codes.

7 Rules-Of-Thumb

It is obvious that the quality of the typical output of any code, for given physics model, depends on how the user built the calculational model and how he/she handled the code. The geometry description is of primary importance. Then, as stated above, the most essential parameters to control the calculational accuracy for given physics model are the number of incidents (primary events) NEVT and the accuracy of boundary localization in iterative transport algorithm STEPEM. Naturally, the higher NEVT and the smaller STEPEM, the better the result will be. But here we come to a contradiction with both CPU time $t_1$ and time $t_2$ allotted for the whole problem. The strategy would be to keep both $t_1$ and $t_2$ as small as possible. There are a few rules for a user to get the best from the MARS code. These rules are:
1. Required NEVT is determined by a statistical error in a phase-space or geometrical region of interest. The calculated results in a given region $N$ are statistically valid only if a R.M.S. statistical error $\delta \leq 20\%$. So, run until this condition is satisfied. Do short runs first to estimate required NEVT and play with DUMP output.

2. Keep $\text{STEPEM} \leq 0.3 \times t_{\text{min}}$, where $t_{\text{min}}$ is a smallest linear size of the smallest region in the considered geometry.

3. Use as few as possible geometrical regions described in all the standard, extended and user—supplied sectors: CPU time grows almost linearly with the number of regions in the direction of predominant propagation of the particles.

4. Versions 13.1(95) and 13.2(95) are built such a way that in many typical applications they give identical results and only at very high energies (TeV region) and/or for very fine geometrical structures the user should use version 13.1(95) with $-r8$ option.

5. Use IND(6) = T and DLEXP=1 options with a great care. Do short tests first.

6. Use IND(8) = F in a routine run to reduce amount of the output.

7. Use IND(1) = F, if you do not need energy deposition related distributions or photon and electron fluxes.

8. Use IND(9) = F, if you are not studying an extremely fine geometrical structure.

9. Use IND(10) = F, if you are not interested in muon production.

10. Use IND(12) = F, if you are not interested in antiproton production.

11. Use IND(14) = F, if you are not specifically interested in low-energy neutron fluxes.

In the last five rules, the code takes care of those components in some effective manner anyway, but the user can reduce the CPU time drastically if he/she turns off the corresponding options.
8 Acknowledgements

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References


