

# The most correct $\rho^0(770)$ meson mass and width values

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**Abstract.** By application of the pion electromagnetic form factor and also the P-wave isovector  $\pi\pi$  scattering phase shift parametrizations, following from the first principles, like unitarity and analyticity, for a description of the corresponding accurate data, it is clearly demonstrated that the  $\rho^0(770)$  meson mass and width values obtained by the Gounaris-Sakurai model in a description of the same data can not be accepted as correct ones.

## 1 Introduction

The initial state radiation method gives very precise information on the pion electromagnetic (EM) form factor (FF)  $F_\pi^{EM,I=1}(t)$  [1, 2], measuring total cross-section

$$\sigma_{tot}(e^+e^- \rightarrow \pi^+\pi^-) = \frac{\pi\alpha^2(0)}{3t} \beta_\pi^3(t) \left| F_\pi^{EM,I=1}(t) + R e^{i\phi} \frac{m_\omega^2}{m_\omega^2 - t - im_\omega\Gamma_\omega} \right|^2 \quad (1)$$

with the pion velocity  $\beta_\pi(t) = \sqrt{1 - \frac{4m_\pi^2}{t}}$ , the  $\rho - \omega$  interference phase  $\phi$  to be expressed as  $\phi = \arctg \frac{m_\rho\Gamma_\rho}{m_\rho^2 - m_\omega^2}$  and the amplitude  $R$  as a free parameter.

On the other hand, the most accurate up to now P-wave isovector  $\pi\pi$  scattering phase shift  $\delta_1^1(t)$  data at the elastic region with theoretical errors, to be generated from the existing inaccurate experimental information by the Garcia-Martin-Kamiński-Peláez-Yndurain (GKPY) Roy-like equations with an imposed crossing symmetry condition [3], appeared.

Further we demonstrated that the the mass and width values of the  $\rho^0(770)$  meson following from the application of the Gounaris-Sakurai model to a description of these data are not trustworthy.

## 2 Analysis

The extensively quoted pion EM FF G.-S. model [4] has been constructed by assuming that for a wide energy of the elastic region up to  $t = 1 \text{ GeV}^2$  the P-wave isovector  $\pi\pi$  scattering

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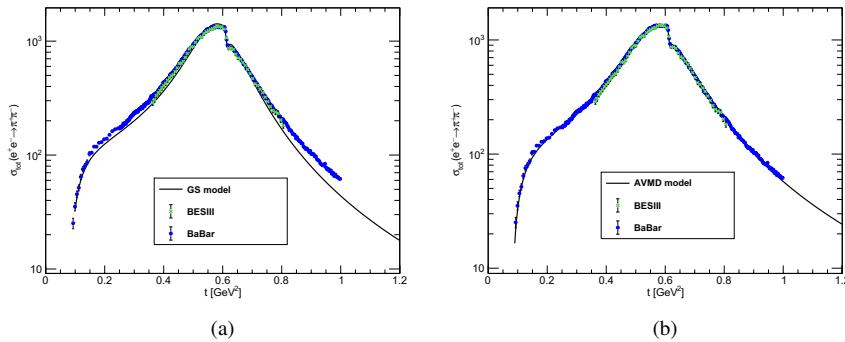


Figure 1: Optimal description of the unified BESIII-BaBar data on  $\sigma_{tot}(e^+e^- \rightarrow \pi^+\pi^-)$  at the elastic region: (a) by pion EM FF G.-S. model (b) by pion EM FF U&A model.

phase shift  $\delta_1^1(t)$  satisfies a two parametric effective-range formula of the Chew-Mandelstam type [5]

$$\frac{q^3}{\sqrt{t}} \cot \delta_1^1(t) = a + b q^2 + q^2 h(t); \quad h(t) = \frac{2}{\pi} \frac{q}{\sqrt{t}} \ln\left(\frac{\sqrt{t} + 2q}{2m_\pi}\right); \quad q = [(t - 4m_\pi^2)/4]^{1/2} \quad (2)$$

and takes the following form

$$F_\pi^{GS}(t) = \frac{m_\rho^2 + m_\rho \Gamma_\rho \left( \frac{3}{\pi} \frac{m_\pi^2}{q_\rho^2} \ln\left(\frac{m_\rho + 2q_\rho}{2m_\pi}\right) + \frac{m_\rho}{2\pi q_\rho} - \frac{m_\pi^2 m_\rho}{\pi q_\rho^3} \right)}{(m_\rho^2 - t) + \Gamma_\rho \left( \frac{m_\rho^2}{q_\rho^3} \right) (q^2(h(t) - h(m_\rho^2)) + q_\rho^2 h'(m_\rho^2)(m_\rho^2 - t)) - i m_\rho \Gamma_\rho \left( \frac{q}{q_\rho} \right)^3 \frac{m_\rho}{\sqrt{t}}}. \quad (3)$$

Its application to an optimal description of the unified BESIII-BaBar data [1, 2] at the elastic region (see Fig.1a) with  $\chi^2/ndf = 40.6341$  gives parameters

$$\begin{aligned} m_\rho &= (775.73 \pm 0.10) MeV \\ \Gamma_\rho &= (126.51 \pm 0.13) MeV \end{aligned} \quad (4)$$

not in a coincidence with values presented by Review of Particle Physics [6].

On the other hand an application of the pion EM FF model (see Fig.1b)

$$F_\pi^{EM,I=1}(q) = \frac{(q - q_Z)(q_N - q_P)}{(q - q_P)(q_N - q_Z)} \frac{(q_N - q_\rho)(q_N - q_{\bar{\rho}})}{(q - q_\rho)(q - q_{\bar{\rho}})} (f_{\rho\pi\pi}/f_\rho). \quad (5)$$

following from the first principles [7] to a description of the same data gives slightly lower values of parameters

$$\begin{aligned} m_\rho &= (763.026 \pm 0.10) MeV \\ \Gamma_\rho &= (144.233 \pm 0.13) MeV \end{aligned} \quad (6)$$

than in [6].

Which of these three sets of  $\rho^0$  meson parameters can be considered to be correct ones? For a solution of this problem the most accurate up to now P-wave isovector  $\pi\pi$  scattering phase shift  $\delta_1^1(t)$  data [3] have been exploited.

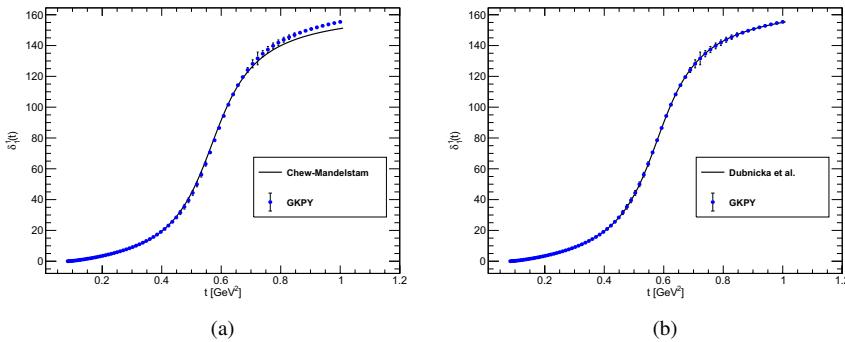


Figure 2: Description of accurate  $\delta_1^1(t)$  data: (a) by effective-range formula of Chew-Mandelstam type with  $\chi^2/ndf=2.4499$  (b) by the model independent parametrization following from the first principles with  $\chi^2/ndf=0.0244$ .

If the effective-range formula of the Chew-Mandelstam type (2) is used to a description of  $\delta_1^1(t)$  data (see Fig.2a) the following  $\rho^0$  meson parameter values

$$\begin{aligned} m_\rho &= (772.42 \pm 0.03) \text{MeV} \\ \Gamma_\rho &= (153.85 \pm 0.11) \text{MeV}. \end{aligned} \quad (7)$$

are found [8], which do not coincide, neither with the values obtained by description of the BESIII-BaBar data on  $\sigma_{tot}(e^+e^- \rightarrow \pi^+\pi^-)$  by G.-S. charged pion EM FF model, nor with values of PDG.

On the other hand, if fully solvable mathematical scheme through the phase representation of the pion EM FF

$$F_\pi^{EM,I=1}(t) = P_n(t) \exp \left[ \frac{t}{\pi} \int_4^\infty \frac{\delta_1^1(t')}{t'(t'-t)} dt' \right], \quad (8)$$

elaborated in [9, 10], is applied, a perfect description of  $\delta_1^1(t)$  data (see Fig.2b), by a model independent phase shift  $\delta_1^1(t)$  representation [8]

$$\delta_1^1(q) = \operatorname{arctg} \frac{A_3 q^3 + A_5 q^5 + \dots}{1 + A_2 q^2 + A_4 q^4 + \dots} \quad (9)$$

with  $\chi^2/ndf = 0.0244$  is achieved and the  $\rho^0$  meson mass and width

$$\begin{aligned} m_\rho &= (763.56 \pm 0.51) \text{MeV} \\ \Gamma_\rho &= (143.09 \pm 0.82) \text{MeV}. \end{aligned} \quad (10)$$

are found to be consistent with (6).

The same values of  $\rho^0$  meson parameters are found also by a generalization of formula (5) to the region of the excited states of the  $\rho^0$  meson [8], what can not be said about results obtained by a generalization of the pion EM FF G.-S. model [4] to the same region.

This all convinces us that the  $\rho^0(770)$  meson mass and width values obtained by the Gounaris-Sakurai model can not be accepted as correct ones.

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