



# GW200105 and GW200115 are compatible with a scenario of primordial black hole binary coalescences

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**Abstract** Two gravitational wave events, i.e. GW200105 and GW200115, were observed by the Advanced LIGO and Virgo detectors recently. In this work, we show that they can be explained by a scenario of primordial black hole binaries that are formed in the early Universe. The merger rate predicted by such a scenario could be consistent with the one estimated from LIGO and Virgo, even if primordial black holes constitute a fraction of cold dark matter. The required abundance of primordial black holes is compatible with the existing upper limits from microlensing, caustic crossing and cosmic microwave background observations.

## 1 Introduction

Based on a second half of the third observing run, the LIGO Scientific and Virgo Collaborations (LVC) [1] reported two gravitational wave events, namely GW200105 and GW200115, which are compatible with neutron star-black hole (NSBH) binaries. At 90% confidence level, the primary components were found to be black holes with masses of  $8.9^{+1.2}_{-1.5} M_{\odot}$  and  $5.7^{+1.8}_{-2.1} M_{\odot}$ , respectively, while the secondary ones  $1.9^{+0.3}_{-0.2} M_{\odot}$  and  $1.5^{+0.7}_{-0.3} M_{\odot}$ , respectively. Compared with the maximal mass of neutron stars, both of the secondaries are compatible with neutron stars with a probability of  $\sim 90\%$ . When the events were assumed to be representative of the entire NSBH population, the merger rate densities of GW200105 and GW200115 were inferred to be  $16^{+38}_{-14} \text{ Gpc}^{-3} \text{ year}^{-1}$  and  $36^{+82}_{-30} \text{ Gpc}^{-3} \text{ year}^{-1}$ , respectively. When a broader distribution of component masses was assumed, they became  $130^{+112}_{-69} \text{ Gpc}^{-3} \text{ year}^{-1}$ . The above rates were found to be consistent with the model predictions of NSBH formation in isolated binaries or young star clusters. For either of the two events, however, LVC found no

evidence of tides or tidal disruption meanwhile identified no electromagnetic counterparts. Therefore, there is no direct evidence that the secondaries are neutron stars.

In this work, we show that a scenario of primordial black hole (PBH) binary coalescences can explain the origin of GW200105 and GW200115. Gravitational collapse of enhanced overdensities could produce PBHs in the early Universe [2–12]. PBHs can form binaries through several formation channels [13, 14]. Amongst these channels, the most efficient one was originally proposed in Refs. [14, 15], and recently extensively studied in Refs. [16–18] by considering such as a generic mass distribution of PBHs and the tidal forces from all the neighboring PBHs and linear density perturbations. To determine if GW200105 and GW200115 are primordial, we have to check whether their merger rate densities can be accounted for with the corresponding PBH abundance allowed by the existing upper limits (see review in Ref. [19] and references therein). In the following, we will compute the merger rate density of PBHs in two specific models by following Ref. [20]. The first one assumes that the mass function is determined by GW200105 and GW200115. The second one assumes that all the black hole binary coalescences observed by LVC are primordial. The predicted rates will be compared with the estimated ones from LVC.

The rest of the paper is arranged as follows. In Sect. 2, we show the formula for the merger rate of PBH binaries. In Sect. 3, we briefly review the log-normal mass function of PBHs. Our results and conclusions are shown in Sects. 4 and 5, respectively.

## 2 Merger rate of PBH binaries

We consider the formation channel of PBH binaries in the early Universe [14]. It is known that this channel makes a

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dominant contribution to the PBH merger rate [17]. This means neglects of the binary formation mechanisms, such as dynamical captures and three-body interactions, in the late Universe [21]. The late-Universe formation mechanism was originally proposed in Ref. [14] and revisited recently in Ref. [15]. Two neighboring PBHs could form a binary due to torque from a third neighboring PBH, and merge with each other within the age of the Universe due to the energy loss via gravitational emissions. However, the mass distribution of PBHs was assumed to be monochromatic in the original literature. More recently, the merger rate of PBHs with an extended mass function was studied in Refs. [16–18]. To be specific, Ref. [16] took into account the tidal force from a PBH which is closest to the center of mass of PBH binary. Reference [17] assumed a flat mass function of PBHs within a narrow mass range. To generalize the above two works, Ref. [18] took into account the tidal forces from all neighboring PBHs and linear density perturbations, and meanwhile considered a generic mass function of PBHs. Therefore, we follow Ref. [18] to compute the merger rate of PBH binaries with a log-normal mass distribution.

For PBHs within mass intervals of  $(m_i, m_i + dm_i)$  and  $(m_j, m_j + dm_j)$ , the merger rate per unit volume within temporal interval  $(t, t + dt)$  is defined as  $R(t) = \mathcal{R}(m_i, m_j, t) dm_i dm_j$  in units of  $\text{Gpc}^{-3} \text{year}^{-1}$ . The comoving merger rate density is obtained as follows [18]

$$\begin{aligned} \mathcal{R}(m_i, m_j, t) &= 3.9 \times 10^6 \times \left(\frac{t}{\tau}\right)^{-\frac{34}{37}} f^2 (f^2 + \sigma_{\text{eq}}^2)^{-\frac{21}{74}} \min\left(\frac{P(m_i)}{m_i}, \frac{P(m_j)}{m_j}\right) \\ &\times \left(\frac{P(m_i)}{m_i} + \frac{P(m_j)}{m_j}\right) (m_i m_j)^{\frac{3}{37}} (m_i + m_j)^{\frac{36}{37}}, \end{aligned} \quad (1)$$

where  $f = \Omega_{\text{pbh}}/\Omega_{\text{m}}$  is the total fraction of non-relativistic matter in PBHs,  $\sigma_{\text{eq}} \approx 0.005$  is the variance of overdensities of the rest of dark matter on scales of order  $\mathcal{O}(10^{-2} - 10^5)M_{\odot}$  at equality [17],  $\min(x_1, x_2)$  selects the minimal value between  $x_1$  and  $x_2$ ,  $P(m)$  is the mass function of PBHs,  $t$  and  $\tau$  denote the cosmic time and the present age of the universe, respectively. The abundance of PBHs in dark matter is  $f_{\text{pbh}} = \Omega_{\text{pbh}}/\Omega_{\text{dm}} \simeq f \Omega_{\text{m}}/\Omega_{\text{dm}}$ . Here,  $\Omega_{\text{pbh}}$ ,  $\Omega_{\text{dm}}$  and  $\Omega_{\text{m}}$  denote the energy density fractions of PBHs, dark matter and non-relativistic matter, respectively, in the critical energy density of the Universe at present. In addition, we would not consider the redshift evolution of the merger rate density, since only the low-redshift events are detectable for LVC. Throughout this work, we fix all cosmological parameters to be the best-fit values from the Planck 2018 results [22].

### 3 Mass function of PBHs

The mass function of PBHs can be determined by the production mechanisms of PBHs in the early Universe. The most likely one is based on the gravitational collapse of overdensities in the radiation dominated epoch of the Universe [2–8]. Therefore, the mass function of PBHs depends on the properties of primordial curvature perturbations. It is usually parametrized to be a log-normal function as follows

$$P(m) = \frac{1}{\sqrt{2\pi}\sigma m} \exp\left(-\frac{\ln^2(m/m_*)}{2\sigma^2}\right), \quad (2)$$

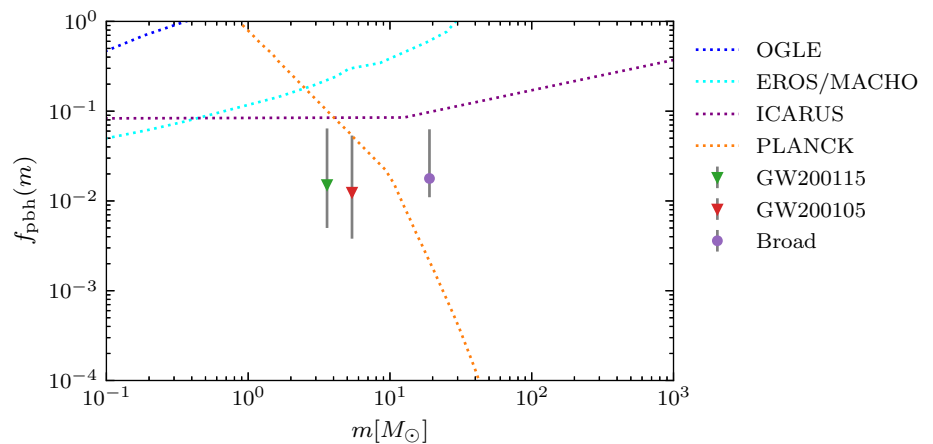
where  $m_*$  and  $\sigma$  stand for the center mass and width, respectively [18, 20, 23]. In this work, we consider two different methods to determine values of  $m_*$  and  $\sigma$ . First, we assume that all observed black hole binaries are primordial. The results have been shown to be  $m_* = 19M_{\odot}$  and  $\sigma = 0.97$  [20], which will be used in this work. There are also alternative choices of mass function [18], which would not significantly change our results of the abundance of PBHs. Second, we assume that only GW200105 and GW200115 are of primordial origin and they determine the mass function of PBHs. Typically, we choose one half of the total mass to determine the value of  $m_*$ , namely  $m_* = 5.4M_{\odot}$  for GW200105 and  $m_* = 3.6M_{\odot}$  for GW200115. Meanwhile, we set the width to be  $\sigma = m_*/10M_{\odot}$ , implying  $\sigma = 0.54$  for GW200105 and  $\sigma = 0.36$  for GW200115. In such a case, slightly different choices would change our predictions on the merger rate density by a factor of  $\mathcal{O}(1)$ . Based on Eq. (1), they would not significantly alter our results of the abundance of PBHs [20].

### 4 Limits on PBH abundance

Our results are shown in Fig. 1. We depict the limits on  $f_{\text{pbh}}$  estimated from the merger rates of GW200105 (red triangle) and GW200115 (green triangle), and the broad mass function (purple circle), respectively. The gray solid lines denote the error bars at 90% CL, due to uncertainties in the measurements of the merger rates [1]. The dotted curves stand for the existing upper limits at 90% CL from OGLE (blue) [24], EROS/MACHO (cyan) [25, 26], ICARUS (purple) [27], and Planck (orange) [28], which are shown here for comparison.

In the case of single PBH event, the mass function is determined by the considered event, i.e. GW200105 or GW200115. Given the observational uncertainty, we derive the allowed abundance of PBHs which matches the observed merger rate of the event [1]. We find that both of the events are well explained by such a scenario, when we take  $f_{\text{pbh}} \simeq \mathcal{O}(10^{-2})$  that is compatible with the existing best constraint  $f_{\text{pbh}} \simeq \mathcal{O}(10^{-1})$  from caustic crossing [27]. We also find that the inferred upper limits are close to the existing upper limit, implying that such a scenario may be further tested

**Fig. 1** Constraints on  $f_{\text{pbh}}$  with 90% CL error bars (gray solid) inferred from the merger rates for GW200105 (red triangle) and GW200115 (green triangle), and the broad mass function (purple circle). The existing upper limits on  $f_{\text{pbh}}$  at 90% CL from OGLE (blue), EROS/MACHO (cyan), ICARUS (purple), and Planck (orange) are shown for comparison



in the near future. Furthermore, it is interesting to find that the mass function with  $m_* = 5.4M_\odot$  and  $\sigma = 0.54$  can almost precisely explain GW200105 and GW200115 simultaneously. But this coincidence may be accidental.

In the case of the broad mass distribution, the mass function is determined by all the observed black hole binaries, since all of them are assumed to be primordial. This has been done by using maximum-likelihood analysis in an existing literature [20]. In our work, the merger rate is predicted via integration over  $m_i \in [2.5, 40]M_\odot$  and  $m_j \in [1, 3]M_\odot$ , following a method consistent with Ref. [1]. Given  $f_{\text{pbh}} \simeq \mathcal{O}(10^{-2})$ , we conclude that the merger rates predicted by this model are compatible with those of GW200105 and GW200115 reported by LVC [1], although the predicted constraint on  $f_{\text{pbh}}$  is in tension with that inferred from the cosmic microwave background measured by Planck [28]. In addition, by combining the broad mass function from Ref. [20] with the two recently observed events, i.e. GW200105 and GW200115, we obtain a joint result of  $m_* = 16.4M_\odot$  and  $\sigma = 1.0$ , which is slightly different with that of Ref. [20].

## 5 Conclusions

In this work, we have shown that the scenario of PBH binary coalescences can account for the observed merger rates of GW200105 and GW200115. We computed the merger rates in two different models. The first one assumed that only GW200105 and GW200115 have a primordial origin. The second one assumed that all the observed black hole binaries are primordial. Given  $f_{\text{pbh}} \simeq \mathcal{O}(10^{-2})$ , we found that both of the two models can explain the recently reported events. For the former, the required abundance of PBHs were found to be compatible with the existing upper limits from other observations. The inferred upper limits were shown to be close to the existing upper limits, implying that such scenarios can be further tested by observations of (SGWB) in the near future [29–33]. However, the latter was found

to be in tension with the existing observational limits. In addition, we found that the log-normal mass function of PBHs with  $m_* = 5.4M_\odot$  and  $\sigma = 0.54$  almost simultaneously explains the observed merger rates of GW200105 and GW200115, including the center values and error bars. In summary, depending on the abundance of PBHs, the scenario of PBH binaries formed through the early-Universe channel could well explain GW200105 and GW200115 events.

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