

## AN UPPER LIMIT ON THE TAU NEUTRINO MASS

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We use the distribution of total charged energy in three-track  $\tau$  decay to place an upper limit on the mass of the  $\tau$  neutrino. Including both statistical and systematic uncertainties, the 95% confidence limit is  $m_\nu < 85$  MeV/c<sup>2</sup>.

## I. Introduction

Neutrino masses have assumed increasing prominence in recent years. Indeed, several of this workshop's major topics ; e.g., dark matter, the solar neutrino deficit, and double  $\beta$ -decay are deeply connected with the possibility of non-zero mass. Because the  $\tau$  lepton at 1784 Mev/c<sup>2</sup> is the most massive charged lepton, one might expect its neutral partner,  $\nu_\tau$ , to be the most massive neutrino. In fact, in the class of theories (see-saw mechanism) where the neutrino mass scales as the *square* of the charged-lepton mass<sup>1)</sup>, an electron-neutrino mass of 20 eV would imply a  $\nu_\tau$  mass of 125 MeV/c<sup>2</sup>.

Several earlier experiments have presented limits on  $m_\nu$ <sup>2-5)</sup>; here I report results based on a large sample of  $\tau^+\tau^-$  events collected at the Cornell Electron Storage Ring (CESR), using the CLEO detector.<sup>6)</sup> The data sample, acquired during 1982-1984, has been used in CLEO's recent determination of the Michel parameter<sup>7)</sup> in  $\tau$  decay. Both that work and an earlier one<sup>8)</sup> use the characteristic 1-vs.-3  $\tau\tau$  topology (see fig. 1) and discuss details of the event selection criteria. Because of space limitations, these will only be sketched herein.

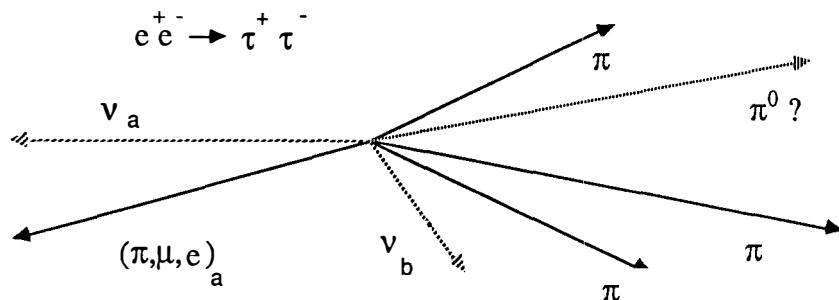


Figure 1 . 1-vs.-3 topology for  $\tau\tau$  events .

## II. Characteristics of the sample

An integrated luminosity of 128.3 pb<sup>-1</sup>, accumulated over a range of center-of-mass energy  $10.34 < W < 11.18$  GeV, yielded 9135  $\tau\tau$  events of the correct topology which passed the above-mentioned criteria, which are designed to reject QED and hadronic backgrounds. No attempt was made to reject events with one or more photons/ $\pi^0$ 's. This minimized the dependence upon CLEO's shower detector, which has c.a. 50% solid angle coverage and only modest resolution. Our primary experimental tools were the 1.0 Tesla solenoid and the CLEO drift chamber, which with >90% solid angle and  $\Delta P/P = 1.2\% * P(\text{GeV}/c)$ , provided high efficiency tracking

and good resolution in the total energy and invariant mass of the 3-track decay. A higher-purity subsample with an identified lepton ( $e$  or  $\mu$ ) as the single track was available to allow crosschecks and internal determination of the background fraction and 3-track invariant mass distribution (See Fig.2)

### III. Method of analysis and Monte Carlo simulation

The "neutral energy"  $E_n = E_{beam} - E_{3\pi}$  is computed for each 3-track decay, where  $E_{beam}$  is the beam energy and  $E_{3\pi}$  is the sum of the 3-track energies, assuming the tracks to be pions. The  $E_n$  spectrum is nearly independent of the particular beam energy. It does depend on the masses of the missing neutrino and of the hadronic system. The low  $E_n$  endpoint (the edge of phase space) varies directly with  $m_\nu$ ; for somewhat larger  $E_n$  the spectral shape depends on a higher power of  $m_\nu$ . For very large  $E_n$  the spectrum is featureless and more affected by missing pions than by  $m_\nu$ .

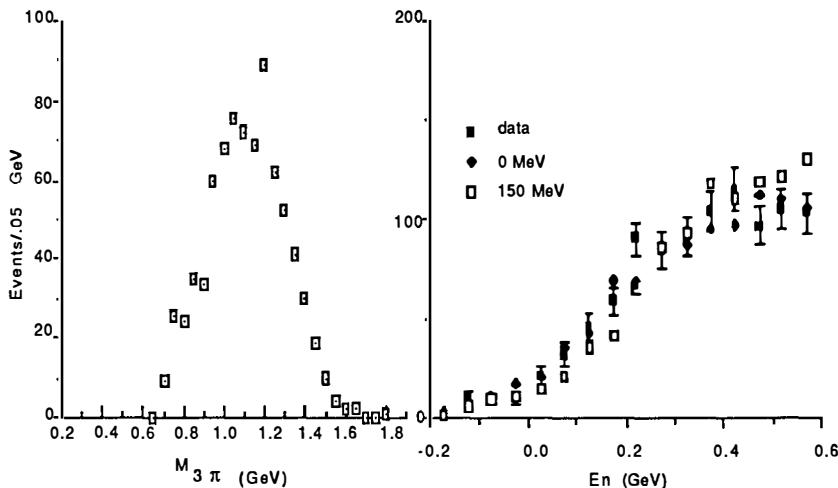


Figure 2. Invariant mass distribution for lepton-tagged events with all particles treated as pions.

Figure 3.  $E_n$  spectra for data and for two values of  $m_\nu$ : 0 (diamonds) and 150 (closed squares).

The  $E_n$  spectrum from the data is compared to Monte Carlo simulations in which  $m_\nu$  is assigned the values 0, 25, 50, 100, and 150 MeV/c<sup>2</sup>. These simulations include realistic branching fractions into  $\rho$ 's,  $A_1$ ; etc., initial state radiation down to 5 MeV, and a 22% hadronic background, determined from the relative  $R$  values and trigger efficiencies. Various distributions (decay angle, invariant mass) are checked for conformity to the data. The events are weighted so as to yield

3-track invariant mass spectra which are constrained to agree with that derived from the lepton-tagged subsample (Fig. 2). In the comparisons, the "theoretical" distributions are normalized to the total data in the fitted region. Figure 3 shows the two cases of  $m_\nu = 0, 150 \text{ MeV}/c^2$  compared to the data, for the range  $E_n < 0.6 \text{ GeV}/c^2$ .

#### IV. Results

The data are fitted for four ranges of  $E_n$ . For each case  $\chi^2$  is plotted versus  $m_\nu$  and the minimum is found from a parabolic fit. The statistical error in  $m_\nu$  is found from the curvature of the parabola in the usual way. Averaging these results gives  $m_\nu = 31 \pm 25 \text{ MeV}/c^2$  (statistical). Dominant systematic errors are indicated as follows: 1) variation with fitted region (over and above that expected from statistics),  $\pm 15 \text{ MeV}/c^2$ , 2) absolute momentum scale,  $\pm 8 \text{ MeV}/c^2$ , 3)  $\pi^0$  fraction,  $\pm 7 \text{ MeV}/c^2$ , 4) errors in hadronic mass distribution,  $\pm 5 \text{ MeV}/c^2$ . Combining errors in quadrature yields  $m_\nu = (31 \pm 25_{\text{stat}} \pm 20_{\text{sys}}) \text{ MeV}/c^2$ . Following the Particle Data Group procedures, we extract a 95% confidence limit  $m_\nu < 85 \text{ MeV}/c^2$ . This value is of comparable sensitivity and in agreement with other recent measurements of  $m_\nu$ .

#### V. References

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