

# Potential of $^{124}\text{Sn}$ isotope in the search for neutrinoless double beta decay

M. K. Singh<sup>1,2,\*</sup>, H. T. Wong<sup>1</sup>, and V. Singh<sup>3</sup>

<sup>1</sup>*Institute of Physics, Academia Sinica, Taipei 115201, Taiwan.*

<sup>2</sup>*Department of Physics, Banaras Hindu University, Varanasi 221005, India. and*

<sup>3</sup>*Physics Department, School of Physical and Chemical Sciences, Central University of South Bihar, Gaya 824236, India.*

## Introduction

Globally, extensive research efforts are focused on the experimental detection of neutrinoless double beta decay ( $0\nu\beta\beta$ ). The rationale is that the search for  $0\nu\beta\beta$  represents the most sensitive experimental avenue for determining if neutrinos are Majorana particles, further providing insights into the absolute neutrino mass scale and the mechanism of mass generation [1]. Next-generation  $0\nu\beta\beta$  experiments target the inverted mass hierarchy (IH) and strive to reach the normal hierarchy (NH) sensitivity domain.

The pursuit of  $0\nu\beta\beta$  is a key focus in nuclear and particle physics, targeting 35 isotopes where  $0\nu\beta\beta$  is allowed while single  $\beta$ -decay is forbidden, highlighting their experimental importance [2]. Among the isotopes used in  $0\nu\beta\beta$  searches, the natural isotopic abundance of  $^{124}\text{Sn}$  is relatively moderate, at approximately 5.79%. This is advantageous for experiments since a higher natural abundance reduces the need for isotopic enrichment, making it a potentially attractive candidate for  $0\nu\beta\beta$  experiments. However, its  $Q_{\beta\beta}$ -value (2287.7 keV – see Fig. 1, top panel) is lower compared to some other isotopes, which can pose a challenge in terms of background discrimination. Around the  $Q_{\beta\beta}$ -value of  $^{124}\text{Sn}$ , the main contributors to ambient  $\gamma$ -radiation are the 2.2 MeV line from  $^{214}\text{Bi}$  and the 2.6 MeV line from  $^{208}\text{Tl}$ . Excellent energy resolution ( $\Delta$ ) can significantly minimize the Compton edge contamination from such back-

ground structures within the selected region of interest (RoI), which is further beneficial for selecting a narrower RoI around the  $Q_{\beta\beta}$ .

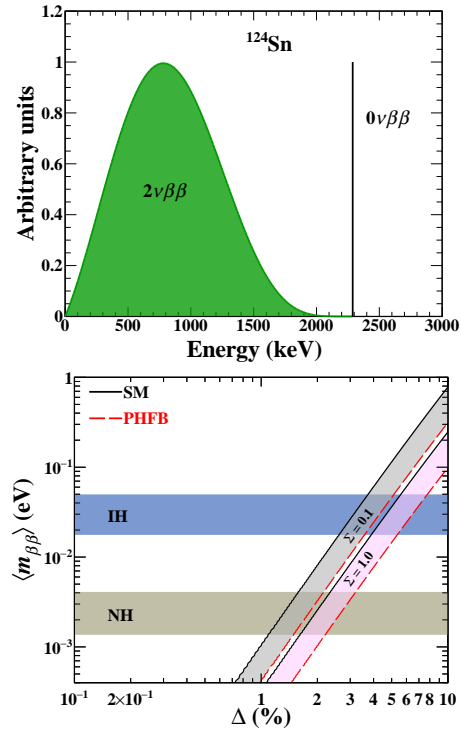


FIG. 1: (Top) The combined energy spectrum of two electrons emitted in both  $2\nu\beta\beta$  and  $0\nu\beta\beta$  decay modes for  $^{124}\text{Sn}$ . (Bottom) Contamination of  $2\nu\beta\beta$  events within the RoI is represented in the  $\langle m_{\beta\beta} \rangle$  versus  $\Delta$  space for exposure  $\Sigma=0.1$  and 1.0 ty (ton-year), where the uncertainty in  $|M^{0\nu}|$  causes a corresponding uncertainty in  $\langle m_{\beta\beta} \rangle$ .

## Sensitivity projection

The INdia's TIN (TIN.TIN) is a proposed experiment aimed at searching for  $0\nu\beta\beta$  decay using 90% enriched  $^{124}\text{Sn}$  as the target iso-

\*Electronic address: [manu@gate.sinica.edu.tw](mailto:manu@gate.sinica.edu.tw)

tope. Our research aims to anticipate the discovery potential of the TIN.TIN experiment, which will employ cryogenic bolometer technology [3].

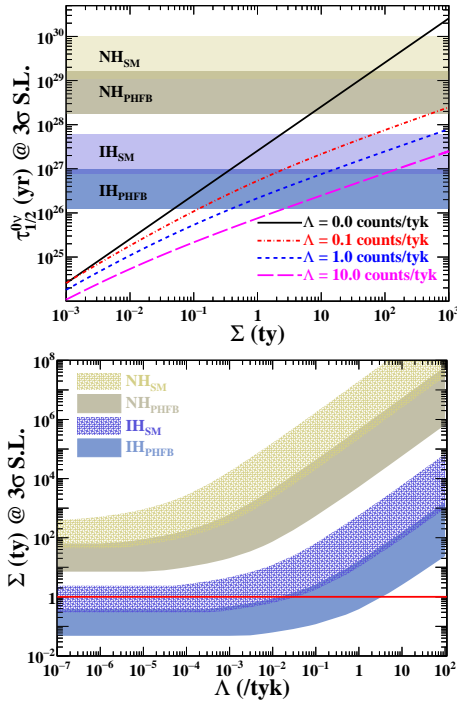


FIG. 2: (Top) The required sensitivity at a  $3\sigma$  significance level (SL) in  $\tau_{1/2}^{0\nu}$  versus  $\Sigma$  at  $\Delta_0$  is determined for  $\Lambda=(0,0.1,1.0,10)$ /tyk [ty-keV]. The IH and NH bands are shown for both PHFB and SM models. (Bottom) Sensitivity requirements for both fully covering and just entering the hierarchy at  $3\sigma$ SL in the  $\Sigma$  versus  $\Lambda$  space at  $\Delta_0$ .

The  $0\nu\beta\beta$  search is intrinsically affected by the  $2\nu\beta\beta$  background. Finite detector resolution ( $\Delta$  [FWHM] in %) introduces irreducible  $2\nu\beta\beta$  events, contaminating the RoI, as shown in the bottom panel of Fig. 1. The sensitivity of effective mass of Majorana neutrino  $\langle m_{\beta\beta} \rangle$  derived from the conversion of half-life  $\tau_{1/2}^{0\nu}$  is influenced by theoretical uncertainties in the nuclear matrix element  $|M^{0\nu}|$ . This results in a band structure for  $\langle m_{\beta\beta} \rangle$  sensitivity due to the  $2\nu\beta\beta$  background alone (in addition to the IH and NH bands), with maximal uncertainty observed between the shell model (SM) and Projected Hartree-Fock-Bogoliubov (PHFB) model. The maximum uncertainty

range of  $|M^{0\nu}|$  for  $\Sigma = 1.0$  ty shows that the region free from  $2\nu\beta\beta$  contamination begins at  $\Delta = 1.6\%$  for the SM and  $\Delta = 2.2\%$  for the PHFB model in the NH case. In the IH case, this uncontaminated region starts at  $\Delta = 3.9\%$  for SM and  $\Delta = 5.3\%$  for PHFB.

The dependence of  $\tau_{1/2}^{0\nu}$  at  $3\sigma$ SL on  $\Sigma$  for projected  $\Delta_0$  ( $\equiv 0.5\%$ ) with varying background rates ( $\Lambda$ ) is presented in top panel of Fig. 2. It is found that to enter the  $\text{IH}_{\text{PHFB}}$  region with a background rate of  $\Lambda=0.1$ /tyk, the TIN.TIN experiment must achieve a sensitivity of  $\Sigma=0.12$  ty, while a sensitivity of  $\Sigma=1.7$  ty is needed to enter the  $\text{IH}_{\text{SM}}$  region at  $3\sigma$ SL.

The parameter space of  $\Sigma$  versus  $\Lambda$  at  $\Delta_0$ , along with the uncertainty bands for  $|M^{0\nu}|$  for both IH and NH, is used to explain the need of precise calculation of  $|M^{0\nu}|$  and the potential for background improvement (see bottom panel, Fig. 2). At  $\Lambda=0.1$ /tyk, achieving coverage of  $\text{IH}_{\text{PHFB}}$  necessitates  $\Sigma=2.6$  ty, while for  $\text{IH}_{\text{SM}}$ ,  $\Sigma=66$  ty is required to identify the signal at  $3\sigma$ SL.

## Summary and prospects

Investigating  $0\nu\beta\beta$  decay in multiple nuclei is essential, and among the double-beta decay candidates,  $^{124}\text{Sn}$  has received limited attention [4], though it is expected to be a focus in upcoming experiments such as TIN.TIN [3]. To achieve full coverage of the  $\text{IH}_{\text{PHFB}}$ , the TIN.TIN experiment with  $\Delta_0=0.5\%$  requires a minimum exposure of  $\Sigma_{\text{min}}=0.38$  ty. In contrast, a conservative approach for covering  $\text{IH}_{\text{SM}}$  necessitates a  $\Sigma_{\text{min}}=2.4$  ty. For a comprehensive investigation into sensitivity, readers are recommended to Ref. [2].

## References

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