

## Structure of the negative parity states in $^{131}\text{Xe}$

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### Introduction

Non-axiality in nuclear shapes has long been a subject of interest as it causes a variety of structural phenomena. Xe isotopes are well known for their triaxial shapes [1] and hence become a good laboratory to test several processes to generate higher angular momentum. Unique parity  $\nu h_{11/2}$  orbital plays a crucial role on the structure of these nuclei. Systematically, two  $I^\pi = 13/2^-$  states were reported in  $^{119-127}\text{Xe}$ , as shown in FIG. 1. Rotational band built on the yrast  $13/2^-$  state is designated as the unfavoured signature partner of the  $\nu h_{11/2}$  band. On the other hand, the band built on yrare  $13/2^-$  state is described as the quasi- $\gamma$ -band coupled to an odd-quasineutron in  $h_{11/2}$  orbital. However, these interpretations fail to answer certain issues as discussed in Ref. [2]. Recently, different interpretations

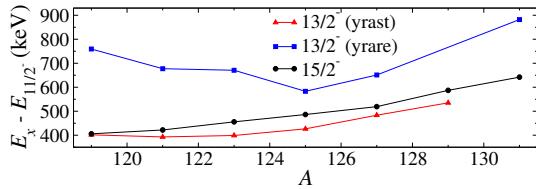


FIG. 1: Variation of the excitation energy, relative to  $E_{11/2^-}$ , of  $15/2^-$  and  $13/2^-$  (yrast and yrare) states in  $^{119-131}\text{Xe}$ , as a function of mass number.

are proposed for the negative parity bands in  $^{127}\text{Xe}$  based on the electromagnetic properties of the  $\Delta I = 1$  inter-band  $\gamma$ -transitions [3]. A detailed spectroscopic investigation on  $^{131}\text{Xe}$  is reported recently, but, only the favoured signature partner of  $\nu h_{11/2}$  band is observed [4]. Therefore, an attempt has been made in this work to search for the unfavoured signature partner of  $\nu h_{11/2}$  band in  $^{131}\text{Xe}$ .

### Experimental Details

Excited states in  $^{131}\text{Xe}$  were populated via  $^{130}\text{Te}(^4\text{He}, 3n\gamma)$  fusion-evaporation reaction at 38 MeV, using K-130 cyclotron at VECC. The target consisted of isotopically enriched  $^{130}\text{Te}$ , 2 mg/cm<sup>2</sup> thick, evaporated on a myler backing (600  $\mu\text{g}/\text{cm}^2$ ). Seven Compton suppressed Clover detectors of the Indian National Gamma Array [5], connected with a PIXIE based digital DAQ [6], were employed to detect and record the time stamped data in single and coincident modes.

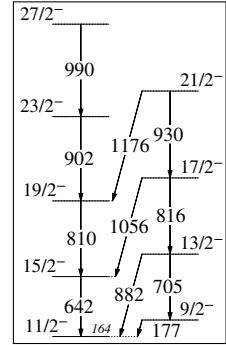


FIG. 2: The  $\nu h_{11/2}$  band in  $^{131}\text{Xe}$ .

### Results

Partial negative parity level scheme of  $^{131}\text{Xe}$ , as shown in FIG. 2, is developed in this work on the basis of  $\gamma\gamma$ -coincidence and inten-

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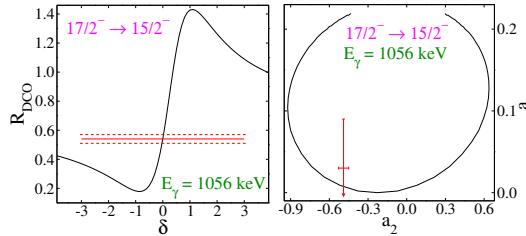


FIG. 3: (Left) Variation of the theoretical DCO ratio as a function of  $\delta$ , and (Right) contour plot of the angular distribution coefficients,  $a_2$  vs.  $a_4$ , for different  $\delta$ . Experimental data point for the 1056 keV  $\gamma$  ray is marked in red.

sity relationships. The band above  $I^\pi = 9/2^-$  state is established by placing three new  $\gamma$  rays, *viz.*, 882, 930 and 1176 keV in addition to the earlier reported 705, 816 and 1056 keV  $\gamma$ -transitions [4, 7]. The dipole (quadrupole) character of 882, 1056 and 1176 keV (705, 816 and 930 keV)  $\gamma$ -rays is confirmed from the present angular correlation measurement ( $R_{DCO}$ ). To estimate the  $E2/M1$  multipole mixing ratio ( $\delta_{E2/M1}$ ) of the inter-band  $\Delta I = 1$   $\gamma$ -rays, the experimental  $R_{DCO}$  is compared with its theoretical values, estimated using the computer code ANGCOR. A representative plot of  $R_{DCO}$  versus  $\delta$  for 1056 keV transition is shown in FIG. 3. It is evident from this figure that this  $\gamma$  ray has a very low  $E2$  admixture ( $\delta \approx 0$ ). Earlier reported angular distribution coefficients are also found in agreement with such low  $\delta \approx -0.2$  value (FIG. 3) [7].

## Discussions

Structure of the  $\alpha = +1/2$  partner of  $h_{11/2}$  band in odd- $A$  nuclei has become a topic of investigation nowadays. In a recent study, the earlier reported signature partner of  $\pi h_{11/2}$  band in  $^{135}\text{Pr}$  has been reinterpreted in terms of wobbling excitation mostly, based on the large  $\delta$  value of the connecting transitions [8].

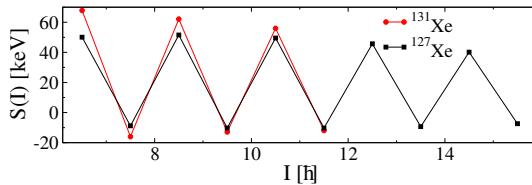


FIG. 4: Energy staggering [ $S(I)$ ] between the  $\alpha = \pm 1/2$  partners of  $\nu h_{11/2}$  band in  $^{127,131}\text{Xe}$ .

Later, the wobbling motion is also confirmed for the so-called unfavoured signature partner of  $\pi h_{11/2}$  ( $\nu h_{11/2}$ ) band in  $^{133}\text{La}$  ( $^{127}\text{Xe}$ ,  $^{105}\text{Pd}$ ) [3, 9, 10]. Therefore, proper identification of the unfavoured signature partner of  $h_{11/2}$  band is important to understand the structure of this band.

Systematically, two bands were reported above yrast and yrare  $13/2^-$  states in  $^{119-127}\text{Xe}$ . However, no such bands is found in the latest studies on  $^{131}\text{Xe}$  [4, 11]. The spin-parity assignment, and the decay pattern of the newly established band above  $I^\pi = 9/2^-$  state make it a suitable candidate for the unfavoured signature partner of  $\nu h_{11/2}$  band. Low  $E2$  admixture in the  $\Delta I = 1$  inter-band transitions provide further support to this interpretation. The energy staggering [ $S(I)$ ] between  $\alpha = \pm 1/2$  partners of  $\nu h_{11/2}$  band in  $^{131}\text{Xe}$  is found to be similar to that reported in  $^{127}\text{Xe}$  (FIG. 4).

## Summary

Low-lying negative parity states in  $^{131}\text{Xe}$  were studied following an  $\alpha$ -induced reaction at VECC, Kolkata. The unfavoured signature partner of  $\nu h_{11/2}$  band has been identified with firm spin-parity assignment. The  $S(I)$  value is found to be large and comparable in magnitude with that reported in  $^{127}\text{Xe}$ . The detailed results along with the theoretical calculations under the framework of particle-rotor model will be presented.

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