




Article

Uncertainties in Atomic Data for Modeling Astrophysical Charge Exchange Plasmas

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Abstract: Relevant uncertainties of theoretical atomic data are vital to determining the accuracy of plasma diagnostics in a number of areas, including, in particular, the astrophysical study. We present a new calculation of the uncertainties on the present theoretical ion-impact charge exchange atomic data and X-ray spectra, based on a set of comparisons with the existing laboratory data obtained in historical merged-beam, cold-target recoil-ion momentum spectroscopy, and electron beam ion traps experiments. The average systematic uncertainties are found to be 35–88% on the total cross sections, and 57–75% on the characteristic line ratios. The model deviation increases as the collision energy decreases. The errors on total cross sections further induce a significant uncertainty to the calculation of ionization balance for low-temperature collisional plasmas. Substantial improvements of the atomic database and dedicated laboratory measurements are needed to obtain the current models, ready for the X-ray spectra from the next X-ray spectroscopic mission.

Keywords: charge exchange; X-ray astrophysics; atomic data; plasma diagnostics



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1. Introduction

Charge exchange plasma can be found in a broad range of astrophysical environments, including, in particular, the interfaces where the solar wind ions interact with neutrals in comets and planetary atmospheres [1–4], but potentially also in supernova remnants [5,6], star-forming galaxies [7,8], active galactic nuclei [9], and clusters of galaxies [10,11]. The modeling of the X-ray spectrum of charge exchange has become possible recently thanks to the efforts of Smith et al. [12] and Gu et al. [13]. These models are crucial to interpreting the observations, as well as to understanding the physical sources that power the plasma.

There is an increasing demand from the astronomical community that the plasma model should provide an estimate of the systematic uncertainties for the atomic data used. This is triggered by the accumulating evidence that the uncertainties from the atomic data, which are not accounted for at present, are as significant as the typical errors from instrumental calibration (see [14] for a recent example). So far, there is no systematic estimate of the uncertainties of the existing charge exchange models, making it difficult to assess the accuracy of the scientific results obtained with these models.

Most of the charge exchange reaction rates in existing models are obtained in theoretical calculations, with only a few laboratory benchmarks performed by several groups with various experimental methods (see, e.g., cross-beam/merged-beam neutral setups: [15–17]; tokamak and laser-produced plasmas: [18–20]; cold-target recoil-ion momentum spectroscopy (COLTRIM): [21–25]; electron beam ion trap (EBIT): [26–32]). A recent comparison using the data from the EBIT measurements [33] showed that the model and the laboratory spectra differ significantly in both line energies and strengths, for the L-shell charge

exchange between nickel ions and neutral particles. Another recent example is that the COLTRIMS measurement by Xu et al. [34] showed that the model calculations might differ from the measurements by 20–50% for the state-selective cross sections of Ne^{8+} and Ne^{9+} charge exchange. In this work, we compile a sample of existing laboratory measurements on charge exchange total cross sections, and state-selective cross sections, as well as characteristic X-ray line ratios, and put forward a systematic assessment of the model accuracy.

This paper is arranged as follows. In Section 2, we describe the sample and the results of the benchmark, and in Section 3, we discuss the potential improvement with future EBIT and COLTRIMS measurements. The benchmark is directly applied to the charge exchange model and atomic data [13] in the SPEX [35] software. Throughout the paper, the errors are given at a 68% confidence level.

2. Methods and Results

2.1. Total Cross Sections

First, we compare the SPEX calculations with existing laboratory results for a number of ions on their total cross sections for atomic hydrogen targets. The SPEX atomic data do not constitute one uniform set of theoretical calculations, but a mixture of three different types of approaches: (1) the rates derived with the empirical scaling reported in Gu et al. [13] (G16 hereafter), which was based on a numerical approximation to a collection of historical theoretical and experimental rates; (2) the multi-channel Landau-Zener method (hereafter MCLZ) reported in Mullen et al. [36]. The atomic data generated by MCLZ are also publicly available in the Kronos database (<https://www.physast.uga.edu/research/stancil-group/atomic-molecular-databases/kronos>, accessed on 1 December 2021); and (3) the recommended values (hereafter RCMD), based on dedicated calculations, including, in most cases, the quantum-mechanical and classical molecular-orbital close-coupling methods, and the atomic-orbital close-coupling method. The G16 approach can calculate, for any ions with a given atomic number and charge, the MCLZ data covering most of the H- and He- like ions with atomic number up to 30, and the RCMD rates are available for a small set of key ions, e.g., O VII [37], N VII [38], and C VI [39].

All the three datasets are tested when the corresponding theoretical cross sections (σ_{theo}) and experimental cross sections (σ_{exp} , see Table 1) are available. Examples are shown in Figure 1 for the C VI and O VII data. For C VI, the three calculations converge at the energy range from ~ 100 eV/amu to 4×10^4 eV/amu, while the MCLZ data do not cover higher energies, and the G16 and RCMD data miss the low energy part. For O VII, the difference between the three calculations becomes more significant than in the case of C VI. The cross section derivatives shown in Figure 1 indicate that the differences in the shapes of the three theoretical calculations become, in general, larger at lower collision energies.

In Figure 2, we plot the distributions of the absolute errors $\sigma_{\text{exp}} - \sigma_{\text{theo}}$ of the theoretical models. The standard deviations of the absolute errors are $1.4 \times 10^{-15} \text{ cm}^2$, $2.0 \times 10^{-15} \text{ cm}^2$, and $1.3 \times 10^{-15} \text{ cm}^2$, for the G16, MCLZ, and RCMD calculations, respectively. As shown in Figure 2b, the absolute errors of G16 become more scattered, and on average larger, at lower collision velocities. The standard deviations of the error distributions are $1.8 \times 10^{-15} \text{ cm}^2$ for $v < 600 \text{ km s}^{-1}$, and $0.9 \times 10^{-15} \text{ cm}^2$ for $v \geq 600 \text{ km s}^{-1}$.

We also summarize the relative deviations $(\sigma_{\text{exp}} - \sigma_{\text{theo}}) / \sigma_{\text{theo}}$ of the three calculations in Figure 2. The average absolute values of the fractional deviations are 55%, 88%, and 35% for the G16, MCLZ, and RCMD datasets. Similar to G16, the MCLZ calculation also has larger relative errors for low-velocity collisions, while the RCMD calculation shows fairly constant deviations for the velocity range considered. For high-energy collisions of $v > 3000 \text{ km s}^{-1}$, the three methods show reasonable agreement with the laboratory results within uncertainties $< 50\%$.

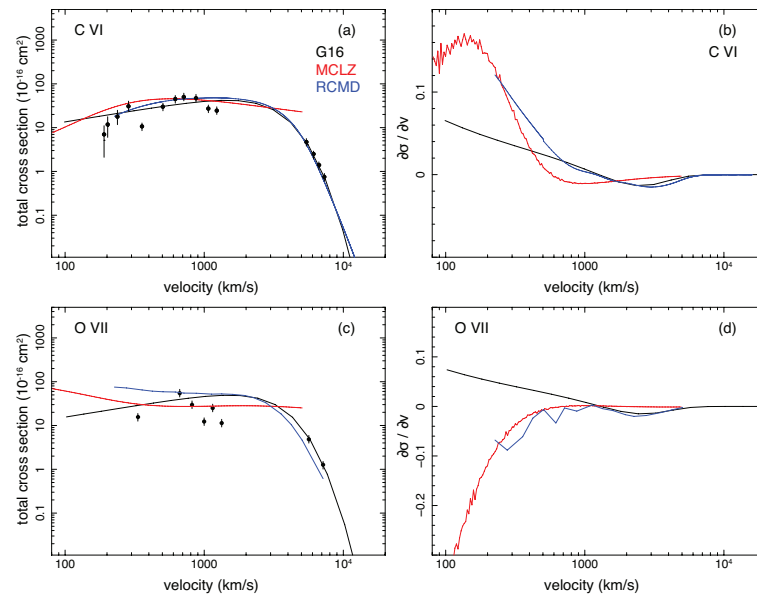


Figure 1. Total cross sections as a function of collision velocity and the cross section derivatives with respect to the velocity for C^{6+} (a,b) and O^{7+} (c,d) ions interacting with hydrogen atoms, resulting in C VI and O VIII ions. The data points are experimental results from Goffe et al. [40], Phaneuf et al. [41], Panov et al. [42], and Meyer et al. [43]. Approximate errors of 15% [43] are shown, except for the low energy ($<500 \text{ km s}^{-1}$) data of C VI, for which the actual errors were reported in the original paper. The solid lines are the model values from the calculations with the G16 (black), MCLZ (red), and RCMD (blue) methods. The abbreviations are explained in the text.

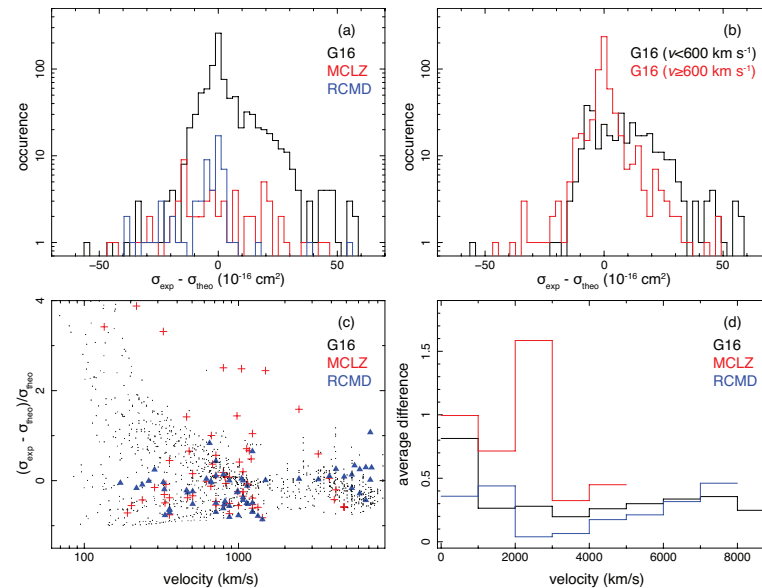


Figure 2. Distributions of the absolute (upper) and relative (lower) deviations of the theoretical cross sections from the experimental results obtained with the measurements summarized in Table 1. (a) Diagrams of the absolute errors for the G16 (black), MCLZ (red), and RCMD (blue) theories. (b) Diagrams of the absolute errors for G16 for low-collision velocities (black) and high velocities (red). (c) The relative deviations for the G16 (black points), MCLZ (red crosses), and RCMD (blue triangles) calculations. (d) The average deviations in absolute values for the three methods in each velocity interval.

Table 1. Experimental cross section data.

Reference	Type ^a	Ion	Theory Data
Shah et al. [44]	total	Li ^{q+} (q = 1–3)	G16
Seim et al. [45]	total	Li ^{q+} (q = 2–3), N ^{q+} (q = 2–5), Ne ^{q+} (q = 3–5)	G16
Goffe et al. [40]	total	B ^{q+} (q = 1–5), C ^{q+} (q = 1–4)	G16
Goffe et al. [40]	total	C ^{q+} (q = 5, 6), N ⁷⁺	G16, MCLZ, RCMD
McCullough et al. [46]	total	B ²⁺ , C ⁺ , N ⁺ , Mg ²⁺	G16
Crandall et al. [47]	total	B ^{q+} (q = 2–5), C ^{q+} (q = 3, 4), N ^{q+} (q = 3, 4), O ^{q+} (q = 5, 6)	G16
Gardner et al. [48]	total	B ^{q+} (q = 2–4), C ^{q+} (q = 2–4), N ^{q+} (q = 2–5), O ^{q+} (q = 2–5)	G16
Phaneuf et al. [49]	total	C ^{q+} (q = 1–4), N ^{q+} (q = 1–5), O ^{q+} (q = 1–5), Si ^{q+} (q = 2–7)	G16
Nutt et al. [50]	total	C ²⁺	G16
Phaneuf et al. [41]	total	C ^{q+} (q = 3, 4), O ^{q+} (q = 2–6)	G16
Phaneuf et al. [41]	total	C ^{q+} (q = 5, 6)	G16, MCLZ, RCMD
Sant’Anna et al. [51]	total	C ³⁺	G16
Ciric et al. [52]	total,nl	C ^{q+} (q = 3, 4), N ⁵⁺ , O ⁶⁺	G16
McCullough et al. [53]	total,nl	C ³⁺	G16
Panov et al. [42]	total	C ⁴⁺ , N ⁵⁺ , O ⁶⁺ , Ne ⁸⁺	G16
Panov et al. [42]	total	C ^{q+} (q = 5, 6), N ^{q+} (q = 6, 7), O ^{q+} (q = 7, 8), Ne ^{q+} (q = 9, 10)	G16, MCLZ, RCMD
Dijkkamp et al. [54]	total,nl	C ^{q+} (q = 3, 4), N ⁵⁺ , O ⁶⁺	G16
Fritsch & Lin [55]	total,nl	C ⁴⁺	G16
Hoekstra et al. [56]	total,nl	C ⁴⁺	G16
Stebbing et al. [57]	total	N ⁺ , O ⁺	G16
Fite et al. [58]	total	O ⁺	G16
Meyer et al. [43]	total	B ^{q+} (q = 2–5), C ^{q+} (q = 3, 4), N ^{q+} (q = 3, 4)	G16
Meyer et al. [43]	total	O ^{q+} (q = 3–6), Si ^{q+} (q = 4–9), Fe ^{q+} (q = 4–15)	G16
Meyer et al. [43]	total	O ^{q+} (q = 7, 8)	G16, MCLZ, RCMD
Havener et al. [59]	total	O ⁵⁺	G16
Huber [60]	total	Ne ^{q+} (q = 2–4), Ar ^{q+} (q = 2–4, 6)	G16
Kim et al. [61]	total	Si ^{q+} (q = 2–7)	G16
Beijers et al. [62]	nl	O ³⁺	G16
Rejoub et al. [63]	total	Ne ³⁺	G16
Havener et al. [64]	total	Ne ⁴⁺	G16
Bruhns et al. [65]	total	Si ³⁺	G16
Havener et al. [66]	total	C ³⁺	G16
Mroczkowski et al. [67]	total	Ne ²⁺	G16
Pieksma & Havener [68]	total	B ⁴⁺	G16
Folkerts et al. [69]	total	N ⁴⁺	G16

^a: total = total cross section, nl = nl-resolved cross section.

The laboratory results should have their own uncertainties; however, these values are available for only a part of the measurements. Here, we provide a rough estimate of the combined measurement uncertainty. The mean systematic uncertainties on the cross sections measured in, e.g., Meyer et al. [43], Draganić et al. [70], Cabrera-Trujillo et al. [71], are approximately 15% for the energy range considered. Assuming that this value can be applied to the other laboratory results, the measurement uncertainties are about 1% for the sample used in testing the G16 calculation, and ~4% for the MCLZ and RCMD results. These relatively minor uncertainties can be accepted as the errors of the theoretical deviations obtained above (e.g., 55%, 88%, and 35% for the G16, MCLZ, and RCMD approaches).

The total charge exchange cross section is needed not only for calculating the charge exchange emission, but also to derive the ionization concentration for general cosmic plasmas in collisional ionization or photoionization equilibrium. The uncertainties in the theoretical calculation would introduce systematic uncertainties to the charge state distribution for the low-temperature plasmas where ions and neutral atoms coexist. As shown in Figure 3, we present two test cases on the concentration calculations of N and O ions in collisional ionization equilibrium (CIE). Here, we assume uncertainties of 50% on the charge exchange recombination rates. The induced errors on the charge distributions of N I and O I would become 10% and 60% at an equilibrium temperature of 1.2 eV. The difference between N I and O I errors reflects the different relative contribution of charge exchange to the total recombination in the concentration calculation. This result suggests that the charge exchange atomic data are vital to the modeling accuracy of lowly ionized species for collisional plasmas. It is expected that similar uncertainties would also apply to photoionization modeling, which includes the charge exchange component in the same way.

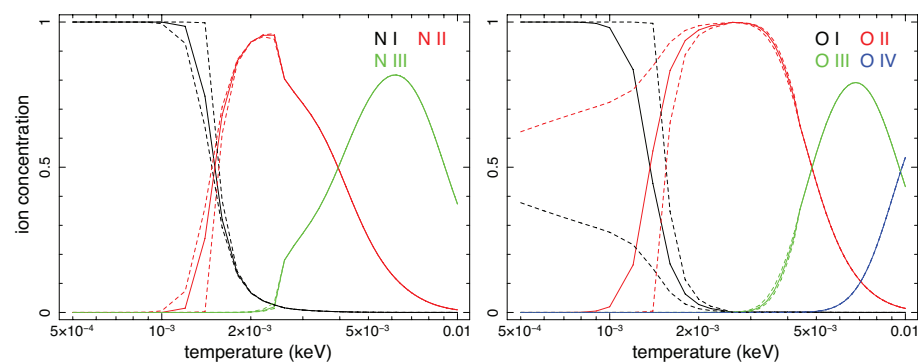


Figure 3. Charge state distributions of N (left) and O (right) as a function of equilibrium temperature for the CIE plasma, calculated with SPEX version 3.06.01. The dashed lines show the calculations when the charge exchange recombination rates are changed by 50%, while the other ionization and recombination data are kept intact.

2.2. Cross Sections for the Peak nl Shells

Next, we examine the state-resolved cross sections. The selective population of high- n levels of the recombining ions is known to be a characteristic property of the charge exchange reaction. The distribution functions on the quantum numbers n and l are key to the calculation of the spectrum, though the present theory still cannot fully reproduce the nl distributions measured in the laboratory [27,33,72].

As shown in Figure 4, we compare the laboratory measurements of four reactions with theoretical calculations, using the G16 method. G16 is the only calculation available in SPEX for the ions tested. It defines empirically n of the most populated levels as functions of the collision velocity, charge, and ionization potential. For the four test cases, G16 successfully predicts the peak n : $n = 3$ for C IV and O III, $n = 4$ for N V and O VI. The cross sections of the peak n levels, however, show deviations from the G16 values at the low energies. For C IV and O III, the measured values for $v = 100 \text{ km s}^{-1}$ are higher by a factor of ~ 2.5 than the theoretical ones. This is probably because the G16 method underestimates the total cross sections at low energies, as already shown in Figure 2. For $v > 500 \text{ km s}^{-1}$, the G16 calculations become consistent, with the measurements within 40% for the peak n .

To assess the l -distribution function, in Figure 4, we also compare the cross sections of the np subshells. The l -distribution defined in G16 is a smooth function that switches as a function of velocity between the different empirical l distributions introduced in Janev & Winter [73] (see also Equations (4)–(8) and Appendix B in [13]). The G16 cross sections on the np shells are lower, by a factor of 2–5, than the experimental values for $v < 500 \text{ km s}^{-1}$. The deviations again become much smaller at higher velocities. To summarize above, the G16 method could reproduce the nl -resolved cross sections for the test cases with an accuracy of $\sim 40\%$ for $v > 500 \text{ km s}^{-1}$, while for the low-velocity collision, the G16 cross sections, as well as the line intensities calculated based on the atomic data, are much less reliable.

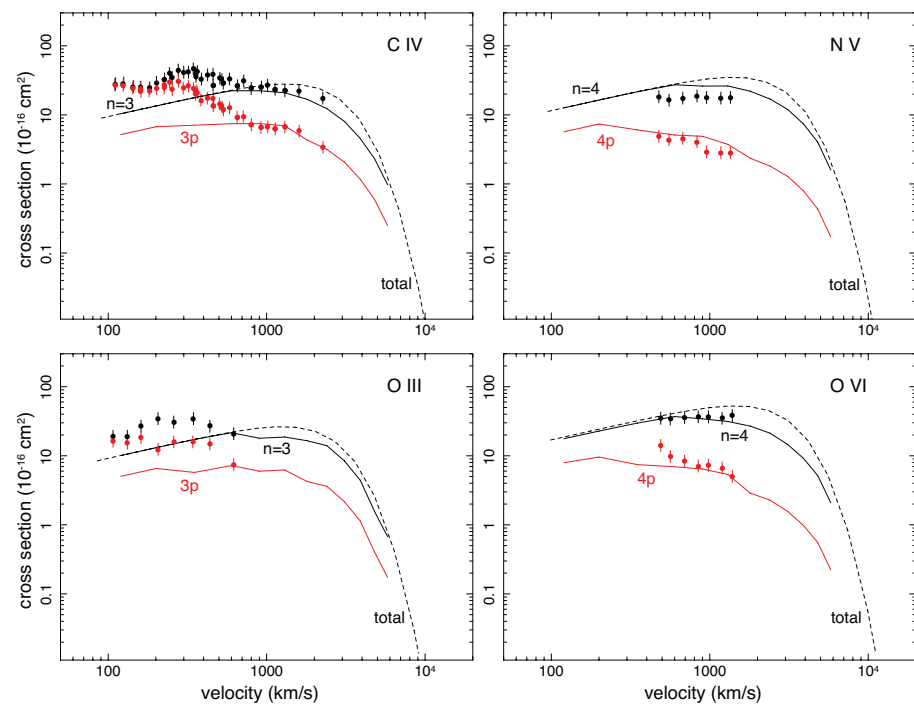


Figure 4. State-selective cross sections as a function of collision velocity for C IV, N V, O III, and O VI. The data points are taken from merged-beam experiments (see Table 1 for detail), for the peak n shells (black) and the np subshells (red). The approximate errors of 15% [43] are shown. The black solid lines are the G16 calculations of the peak n shells, and the red lines are the G16 data for the np subshells. The dash lines are the G16 calculations of the total cross sections.

2.3. Line Ratios

The large ratios between $1s - np$ ($n > 2$) and $1s - 2p$ lines are often used as characteristic diagnostics of the highly charged charge exchange plasma [10,11]. It is known that the line ratios would decrease with increasing collision velocity, because a high-speed collision might yield captures on high angular momentum states, producing more $1s - 2p$ transitions through cascade. So, the line ratios can often be utilized as a probe of collision velocity [27,74]. The accuracy of the velocity measurement is therefore determined by the quality of the atomic data.

In Figure 5, we plot the comparison of the line ratio calculations and experiments for C VI and O VIII. The experimental data are taken from the beam-gas measurements by Andrianarijaona et al. [75] for C VI and Seely et al. [76] for O VIII. A caveat of the comparison is that these experiments used the Kr atom as a target, while the original theoretical calculations are based on capture from H atom. As reported in Leung & Kirchner [77], the line ratios from Kr and H collisions are somewhat different, in particular for the low-energy regime, even though the ionization potentials of Kr and H atoms are nearly the same. To compensate this discrepancy, we calculate the H-to-Kr scalings as a function of velocities on both C VI and O VIII line ratios, using the theoretical results reported in Leung & Kirchner [77] (in their Figures 3 and 6), and apply the scalings to the G16, MCLZ, and RCMD line ratios. The scaled line ratios should represent a better approximation to the collisions with the Kr target.

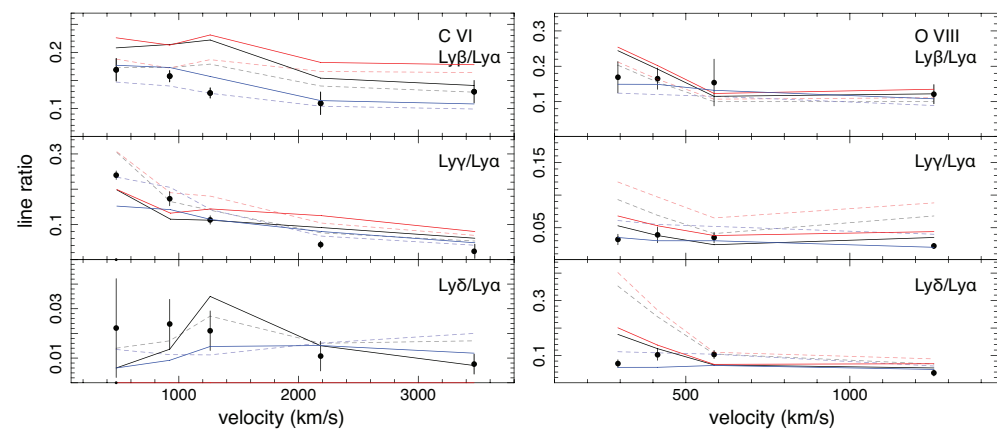


Figure 5. Comparison of experimental and theoretical line ratios for the C^{6+} (left) and O^{8+} (right) reactions. The experimental data from Andrianarijaona et al. [75] and Seely et al. [76] are plotted as data points, and the predictions from G16, MCLZ, and RCMD are shown in black, red, and blue curves. The dashed curves show the calculations of C^{6+} and O^{8+} collisions with H atoms, and the solid lines show the collisions with Kr atoms obtained using the scaling of Leung & Kirchner [77].

As shown in Figure 5, the experiments and calculations yield the same peak n , $n = 4$ for C VI and $n = 5$ for O VIII, though the line ratios still differ at several velocities. One of the main discrepancies occurs between the scaled G16/MCLZ and the lab data for the C VI $Ly\beta/Ly\alpha$ line ratio, where the two theoretical values exceed the measured one by about 70% at $v = 1000 \text{ km s}^{-1}$. The RCMD calculation shows better agreement with the lab values on this line ratio.

A more extensive comparison can be seen in Table 2. It is a compilation of several laboratory efforts, including the recent electron beam ion trap devices with X-ray spectral analysis carried out at both low and high resolutions. The EBIT devices simulate charge exchange reactions only at low collision energies. The average relative discrepancies (experiment-theory/theory) of the line ratios are 0.63, 0.77, and 0.54 for the G16, MCLZ, and RCMD calculations, respectively. For the peak n shell, the average discrepancies are 0.61, 0.81, and 0.56 for the three models. These differences are significantly larger than those on the modeling of collisional ionization equilibrium plasmas ($\sim 10\text{--}40\%$, [14,78,79]), suggesting that the state-of-the-art charge exchange spectral models, even with dedicated theoretical calculations, are still less reliable than those for the CIE plasma.

Table 2. Experimental and theoretical Line ratios.

Ion	$v \text{ (km s}^{-1}\text{)}$	Ratio	Experiment	G16	MCLZ	RCMD	Reference ^a
N VII	794	$Ly\beta/Ly\alpha$	0.76	0.12	0.13	0.10	1
		$(Ly\gamma + Ly\delta)/Ly\alpha$	0.62	0.18	0.20	0.29	
O VII	724	$Ly\beta/Ly\alpha$	0.19	0.07	0.09	0.07	
		$(Ly\gamma + Ly\delta)/Ly\alpha$	0.24	0.11	0.47	0.07	
O VIII	774	$Ly\beta/Ly\alpha$	0.13	0.11	0.10	0.11	
		$(Ly\gamma + Ly\delta)/Ly\alpha$	0.17	0.15	0.18	0.14	
Ne IX	743	$Ly\beta/Ly\alpha$	0.04	0.04	0.12	—	
		$(Ly\gamma + Ly\delta)/Ly\alpha$	0.05	0.05	0.18	—	
Ne X	783	$Ly\beta/Ly\alpha$	0.12	0.08	0.08	0.08	
		$(Ly\gamma + Ly\delta)/Ly\alpha$	0.11	0.06	0.08	0.04	
O VIII	293	$Ly\beta/Ly\alpha$	0.169 ± 0.044	0.244^b	0.254	0.149	2
		$Ly\gamma/Ly\alpha$	0.032 ± 0.008	0.053	0.068	0.035	
		$Ly\delta/Ly\alpha$	0.071 ± 0.014	0.177	0.201	0.057	
		$Ly\epsilon/Ly\alpha$	0.0065 ± 0.003	0.054	0.0061	0.027	
O VIII	414	$Ly\beta/Ly\alpha$	0.165 ± 0.030	0.192	0.202	0.149	
		$Ly\gamma/Ly\alpha$	0.039 ± 0.012	0.038	0.053	0.030	
		$Ly\delta/Ly\alpha$	0.103 ± 0.02	0.125	0.138	0.057	
		$Ly\epsilon/Ly\alpha$	0.005 ± 0.0076	0.031	0.0024	0.019	
O VIII	586	$Ly\beta/Ly\alpha$	0.154 ± 0.006	0.115	0.123	0.132	
		$Ly\gamma/Ly\alpha$	0.035 ± 0.008	0.024	0.038	0.030	
		$Ly\delta/Ly\alpha$	0.104 ± 0.015	0.066	0.068	0.064	
		$Ly\epsilon/Ly\alpha$	0.0048 ± 0.0061	0.015	0.00086	0.014	

Table 2. Cont.

Ion	v (km s ^{−1})	Ratio	Experiment	G16	MCLZ	RCMD	Reference ^a
O VIII	1256	Ly β /Ly α	0.121 ± 0.027	0.122	0.135	0.108	3
		Ly γ /Ly α	0.022 ± 0.004	0.035	0.044	0.020	
		Ly δ /Ly α	0.037 ± 0.011	0.055	0.071	0.050	
C VI	477	Ly ϵ /Ly α	0.0048 ± 0.0028	0.023	0.00045	0.0090	
		Ly β /Ly α	0.169 ± 0.023	0.208 ^b	0.226	0.177	
		Ly γ /Ly α	0.240 ± 0.012	0.198	0.199	0.152	
C VI	924	Ly δ /Ly α	0.022 ± 0.020	0.0062	2.8 × 10 ^{−6}	0.0061	
		Ly β /Ly α	0.157 ± 0.012	0.214	0.213	0.173	
		Ly γ /Ly α	0.173 ± 0.023	0.115	0.132	0.142	
C VI	1262	Ly δ /Ly α	0.024 ± 0.009	0.014	2.7 × 10 ^{−6}	0.0091	
		Ly β /Ly α	0.128 ± 0.009	0.222	0.231	0.157	
		Ly γ /Ly α	0.113 ± 0.012	0.112	0.144	0.114	
C VI	2185	Ly δ /Ly α	0.021 ± 0.008	0.035	2.2 × 10 ^{−6}	0.015	
		Ly β /Ly α	0.109 ± 0.019	0.154	0.182	0.114	
		Ly γ /Ly α	0.043 ± 0.011	0.091	0.125	0.080	
C VI	3466	Ly δ /Ly α	0.011 ± 0.006	0.015	6.7 × 10 ^{−7}	0.015	
		Ly β /Ly α	0.130 ± 0.021	0.141	0.178	0.108	
		Ly γ /Ly α	0.024 ± 0.018	0.061	0.080	0.048	
O VII	low	Ly δ /Ly α	0.0076 ± 0.004	0.007	5.0 × 10 ^{−7}	0.012	4
		He _{high} /He α	0.167	0.168	0.152	0.058	
		He _{high} /He α	0.162	0.161	0.133	—	
Ar XVII	low	He _{high} /He α	0.191	0.133	—	—	
Fe XXV	low	He _{high} /He α	0.267	0.156	0.079	—	
O VIII	low	Ly _{high} /Ly α	1.006	0.786	0.887	0.366	5
Ne X	low	Ly _{high} /Ly α	1.207	0.690	0.865	0.210	
Mg XII	low	Ly β /Ly α	0.227 ± 0.040	0.179	0.205	—	
Mg XII	low	Ly γ /Ly α	0.133 ± 0.022	0.070	0.083	—	
Mg XII	low	Ly δ /Ly α	0.044 ± 0.015	0.038	0.046	—	
Mg XII	low	Ly ϵ /Ly α	0.095 ± 0.015	0.028	0.030	—	
Mg XII	low	Ly ζ /Ly α	0.030 ± 0.018	0.221	0.120	—	
Mg XII	low	Ly η /Ly α	0.080 ± 0.014	0.091	0.287	—	
S XVI	low	Ly β /Ly α	0.203 ± 0.070	0.153	0.171	—	
S XVI	low	Ly γ /Ly α	0.082 ± 0.016	0.055	0.064	—	
S XVI	low	Ly δ /Ly α	0.053 ± 0.011	0.028	0.033	—	5
S XVI	low	Ly ϵ /Ly α	0.053 ± 0.008	0.017	0.020	—	
S XVI	low	Ly ζ /Ly α	0.016 ± 0.005	0.012	0.014	—	
S XVI	low	Ly η /Ly α	0.029 ± 0.008	0.024	0.014	—	
S XVI	low	Ly θ /Ly α	0.111 ± 0.019	0.149	0.101	—	
S XVI	low	Ly ι /Ly α	0.165 ± 0.024	0.058	0.165	—	

^a: references 1 = Greenwood et al. [72]; 2 = Seely et al. [76]; 3 = Andrianarijaona et al. [75]; 4 = Wargelin et al. [80]; 5 = Betancourt-Martinez [81]. ^b: H-to-Kr scaling has been applied to the theoretical line ratios for O VIII and C VI; see text for details.

3. Discussion and Ending Remarks

Based on a large sample of laboratory measurements, we have systematically compared the commonly used charge exchange atomic data to the experimental results. The G16, MCLZ, and RCMD calculations utilized in the SPEX code do not fully reproduce the measurements, with notable, and likely velocity-dependent discrepancies in both total cross sections, state-resolved cross sections, and line ratios in the X-ray spectra. While the ease of the use of the present CX model is beneficial for the X-ray astronomical community, it should be used with caution, in particular for non-charge-exchange experts. The unresolvable disagreements call for advanced theoretical calculations for especially the low collision energy regime, in combination with more laboratory measurements with, in particular, EBIT and COLTRIMS facilities.

The previous EBIT experiments have provided relevant benchmarks to the predicted cross sections for electron capture into specific principal quantum number states n . However, a comparison with the angular-momentum l -resolved cross sections is challenging, as they depend on the collision energy; and the EBIT measurements are limited to low collision energies (<10 eV/u) [27]. Besides, the charge exchange process not only produces X-ray lines, but also generates lines in the ultraviolet and optical band as the Rydberg levels populated by charge exchange relax through radiative cascades to the ground state of the ion. Thus, the simultaneous measurements of EUV and optical charge exchange cascade

photons at the EBIT would be of interest, and they could provide additional information on the population of nl -states for plasma modeling [29]. Furthermore, possible multi-electron capture contributions from the molecular targets used in the EBIT measurements can also be avoided by using an atomic hydrogen target, where only single-electron capture can occur [82]. Atomic hydrogen is of particular interest as it is also the most abundant neutral element in the universe, and it makes a comparison between laboratory measurements and astrophysical observations more reliable.

Besides EBIT, the COLTRIMS and beam-gas experiments have been providing reliable measurements on velocity-dependent total and state-resolved cross sections. The improvement in the momentum measurement technique allows nl selectivity, and for a few cases, it might even be able to resolve the spin state. The state-of-the-art measurement accuracy is about 11% for both the total and nl -resolved cross sections [83].

A systematic measurement of the cosmic abundant ions with the COLTRIMS facilities, in combination with simultaneous EBIT X-ray spectroscopy, is desirable for the astronomical community. A consistent and continuous effort will be needed to ensure that the charge exchange atomic data will be ready for the high-resolution X-ray spectra taken with next-generation missions, XRISM (launch due in 2023, [84]) and Athena (early 2030s, [85]).

Assessing uncertainties carried out by the theoretical atomic data is also vital to the success of the upcoming missions. The atomic physics and plasma code community has already begun this work, with a persistent effort on the evaluation of the errors on electron impact excitation and transition probability data [14,86–90], as well as errors on photon impact data and modeling [91]. One implication from the aforementioned works, including the present work on the charge exchange modeling, is that the classical assumption of constant model uncertainty (e.g., 20% on line emissivity) is no longer valid, since the uncertainties are proven to vary significantly with the underlying model and its key parameters.

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