

# BEVATRON EXTERNAL PROTON BEAM\*

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## EXTRACTION SYSTEM

The achromatic extraction system for the Bevatron employs an energyloss (jump) target and two internal deflecting magnets, each accompanied by a small Panofsky-type quadrupole lens. The beam emerges from the accelerator

approximately one betatron wavelength after the target. This kind of system [1] has several advantages over the corresponding single-magnet extraction system [2, 3].

A. Because the vertical and horizontal betatron periods are nearly equal ( $\nu_h \approx 0.7 \nu_v \approx 0.8$ ), the internal optics can be adjusted to obtain a triple focus near the exit window. As a consequence:

1. The dispersion produced by the energy

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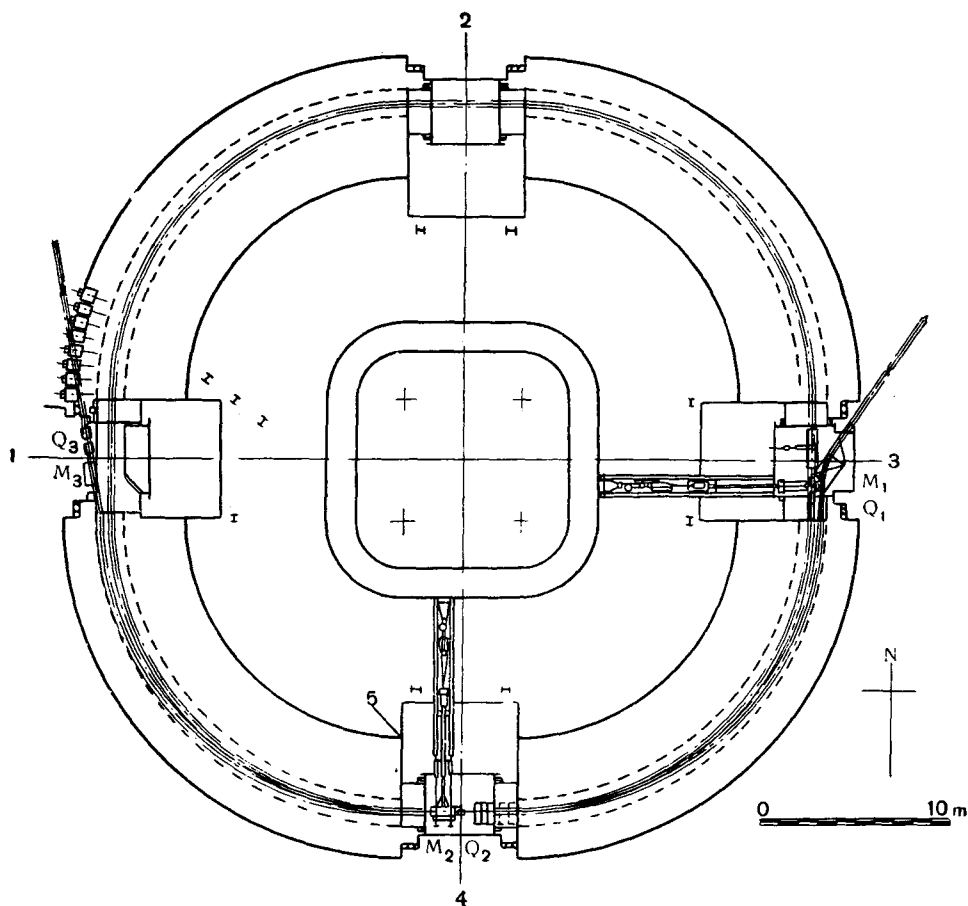


Fig. 1. Layout of the internal deflection system. The circulating beam travels in a clockwise direction:

1 — west straight section; 2 — north straight section; 3 — east straight section; 4 — south straight section; 5 — energy-loss target.

spread in the beam emerging from the target is minimized.

2. There is a reduction in the effect of the fringing field, which not only defocusses strongly horizontally (and focusses vertically), but also is very aberrational.

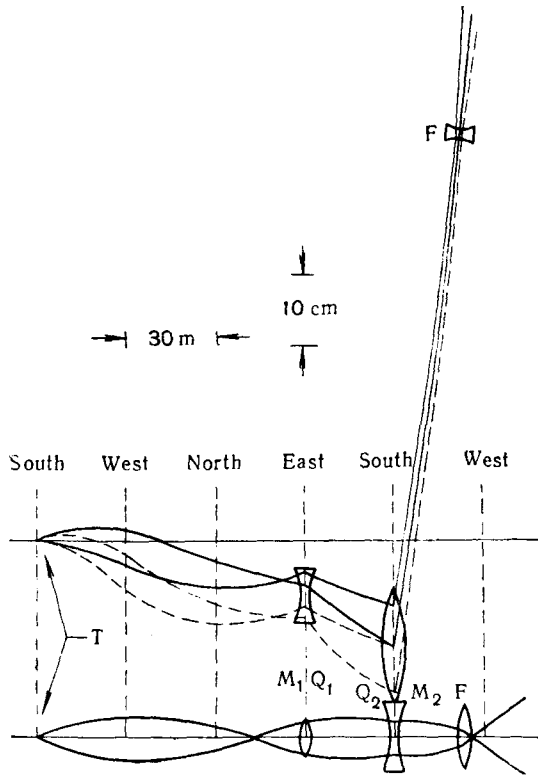


Fig. 2. Optical properties of the internal deflected beam. Here  $M_1$  and  $M_2$  are  $C$  magnets, and  $Q_1$  and  $Q_2$  are «Panoisky-type» quadrupoles.  $F$  represents the effect of the fringing field which focuses strongly vertically and defocusses horizontally.  $T$  is the energy-loss target of beryllium or polyethylene. Note that there is a difference of a factor of 240 between the radial and azimuthal scales.

3. The aperture requirements for the exit channel are reduced.

B. Because only a small deflection is required of the first magnet, a smaller septum, and therefore a smaller target, may be used. This leads to a reduction in beam disorder produced in the targeting process.

The layout of the internal system is shown in Fig. 1. The beryllium target in or near the south straight section is 13 to 40 mm long (depending upon the extraction energy) and  $\approx 6$  mm high. The maximum inward deflection of the beam in

terms of the momentum loss  $\Delta P$  is given by

$$\Delta R_{\max} \approx \frac{2R}{v_h^2} \frac{\Delta P}{P}, \quad (1)$$

where  $R = 15.2$  m is the initial radius of the beam. Normally  $\Delta R_{\max}$  is  $\approx 100$  mm.

Magnets  $M_1$  and  $Q_1$  are located in the east straight section just over one-half betatron period from the target near the point of maximum inward deflection and maximum dispersion of the deflected beam. Consequently, a major function of  $Q_1$  is to control the spatial recombination of beam particles of different energies. Magnet  $M_1$  deflects the beam inward to  $M_2$  and  $Q_2$  located in the south straight section at a radius smaller than that of the target;  $M_2$  deflects the

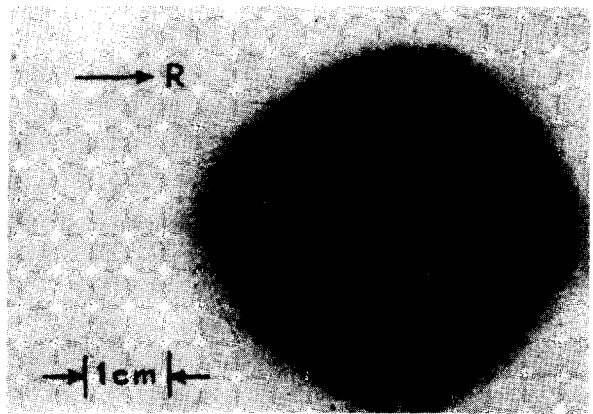


Fig. 3. Pattern of the deflected beam ahead of  $M_1Q_1$ , made from exposure of X-ray film to an irradiated sheet of 3-mm polyethylene. Corners of the aperture of  $M_1$  are shown for reference. External proton beam—6 GeV ahead of  $M_1$ .

beam outwards so that it emerges from the Bevatron at the entrance to the west straight section. A schematic of the internal optics is shown in Fig. 2. The alternating-gradient character of  $Q_1$ ,  $Q_2$ , and the fringing field «lens» is required to limit the aperture required for the beam both inside and outside the accelerator. The design of the internal magnets (see Table 1) was guided by the results of an orbit integration program which used the Bevatron field measurements at both 5 and 7 GeV/c. Predictions of this program have been verified by direct measurement of the beam characteristics.

Figs. 3, 4 and 5 show radiographs made from foils exposed to the deflected beam ahead of  $M_1$  and  $Q_1$ , ahead of  $M_2$  and  $Q_2$ , and at the west

Table 1

Bevatron external-beam magnets

Magnet			Aperture (mm)		Length, m	Design field, kGs	Field gradient, kGs/cm	Field <sup>(a)</sup> uniformity, %	Design <sup>(b)</sup> power, kW	Lamination thickness, mm	Weight, kGs
name	type	no.	width	height							
$M_1$	C	1	76	53	0.32	5.5		$\pm 0.25$	35	6	120
$M_2$	C	1	152	110	1.22	6.3		$\pm 0.1$	84	6	2040
$M_3$	H	1	127	51	1.37	15		$\pm 0.05$	154	6	1860
$M_{30}$	H	1	125	110	0.25	2		$\pm 1$	5.5	—	45
$M_4^{(c)}$	H	4	380	110	2.13	14		$\pm 0.1$	143	13	26400
$Q_1$	Panofsky Quad.	1	82	82	0.64		0.4	$\pm 0.25^{(d)}$	47	6	17
$Q_2$	Panofsky Quad.	1	164	164	0.64		0.2	$\pm 0.25^{(d)}$	29	6	180
$Q_3$	Quad.	2	112 (diam)		0.74		2.4	$\pm 0.25^{(d)}$	155	3	780
$Q_4^{(c)}$	Quad.	6	209 (diam)		1.22		1.2	$\pm 0.25^{(d)}$	203	0.8	5630

a) Based on line integral of field.

b) Peak value for pulsed operation or the dc value corresponding to a beam momentum of 7 GeV/c.

c) Magnets designed but not yet constructed.

d) Maximum relative amplitude of sixth and tenth field harmonics at 85% of the pole-tip radius.

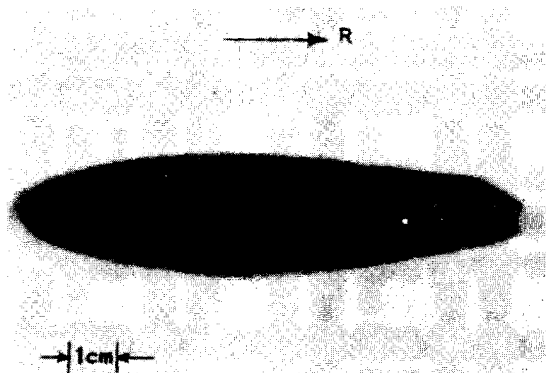


Fig. 4. Pattern of deflected beam ahead of  $M_2Q_2$  made from exposure of X-ray film to an irradiated sheet of 3-mm polyethylene. Corners of the aperture of  $M_2$  are shown for reference. External proton beam—6 GeV ahead of  $Q_2$ .

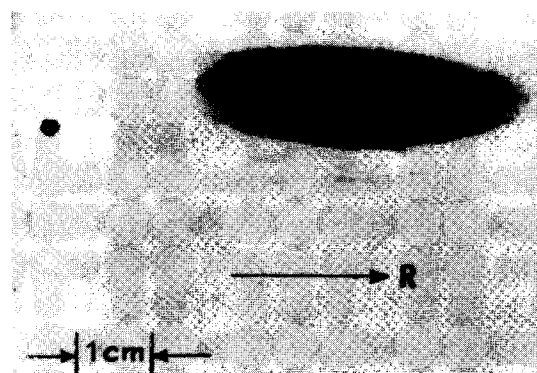


Fig. 5. Pattern of beam emerging at west straight section, made by direct exposure of X-ray film to the 7 GeV/c beam of approximately  $10^8$  protons. External proton beam —.6 GeV emerging from Bevatron.

straight section, respectively. The limiting apertures of each following deflecting magnet is indicated. Additional tests made by simultaneous activation of aluminum foils at the three positions indicate that, as suggested by Figs. 3 and 4, very little ( $\approx 10\%$ ) of the deflected beam is lost on the internal magnets.

#### PROPERTIES OF THE EXTRACTED BEAM

The beam emerges from the west straight section at an average angle to the initial orbit

which varies from 3.2 deg at 3 GeV/c to 4.0 deg at 7 GeV/c. The divergence of the beam indicates vertical and horizontal images 1.5 and 6 m, respectively, upstream from the point of initial exit from the Bevatron. At this location the beam is well confined vertically (Fig. 5). About 75% of the emerging beam is contained in a horizontal width of 50 mm. There is evidence that the beam outside this width is of poor optical quality, for reasons that are not completely understood. Collimation inside the Bevatron has improved somewhat

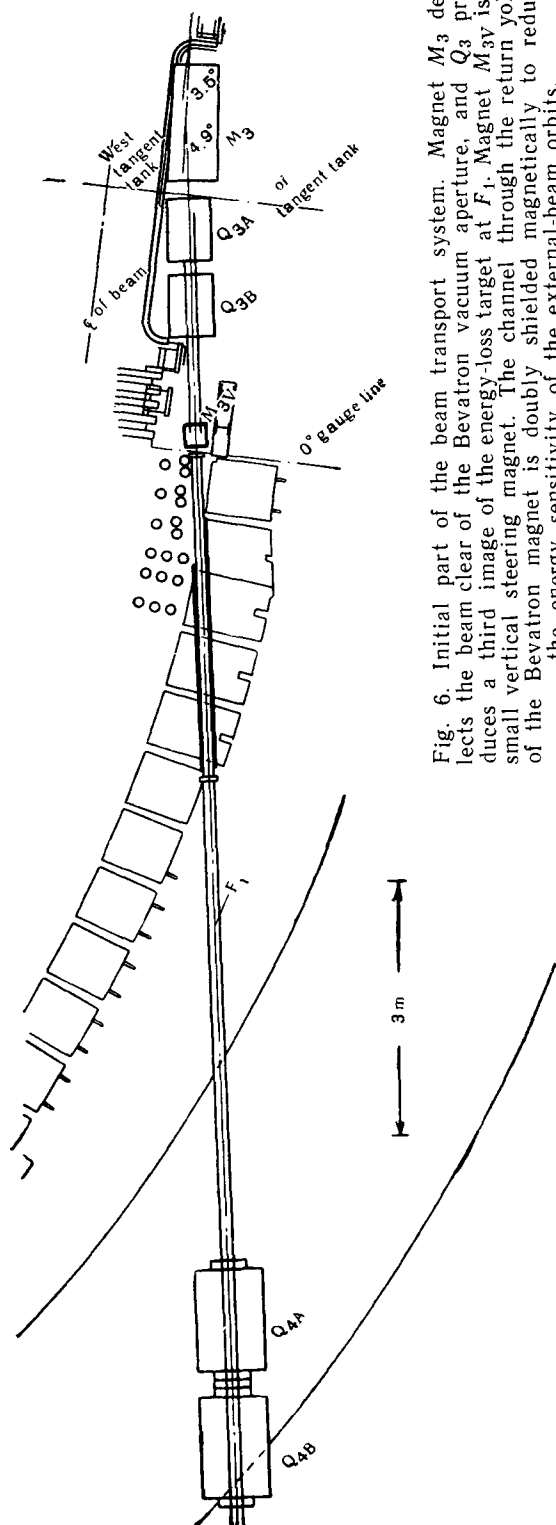


Fig. 6. Initial part of the beam transport system. Magnet  $M_3$  deflects the beam clear of the Bevatron vacuum aperture, and  $Q_3$  produces a third image of the energy-loss target at  $F_1$ . Magnet  $M_{3V}$  is a small vertical steering magnet. The channel through the return yoke of the Bevatron magnet is doubly shielded magnetically to reduce the energy sensitivity of the external-beam orbits.

the quality of the beam available for experimentation. The extracted beam is monitored with a secondary-emission chamber and an ionization chamber, and by activation methods. The intensity varies from about one-third to one-half the circulating-beam intensity, depending upon the degree of collimation. Because, as was stated above, very little deflected beam is lost on the magnets themselves, it is assumed that the major loss of beam during extraction occurs in targeting.

The emittance of the extracted beam has been measured with the help of the deflecting magnet  $M_3$  and the quadrupole doublet  $Q_3$ , which produce a third image of the internal target at  $F_1$  (Figs 6 and 7). At 7 GeV/c the minimum image is 3 mm vertical by 10 mm horizontal (full width at half maximum). The emittance is about 30 mrad·mm vertical and 60 mrad·mm horizontal. These measured values are consistent with the emittances calculated from the initial beam characteristics and the estimated disorder introduced in targeting. The larger horizontal emittance is attributed to the greater horizontal size of the circulating beam, incomplete elimination of the dispersion, and larger magnetic aberrations.

#### PROPOSED TRANSPORT SYSTEM

The installation of the beam up to the main shielding wall of the Bevatron (Figs. 6 and 7) is now complete, and an experimental program using the beam at  $F_1$  is under way. Fig. 7 shows the proposed transport system through the experimental area to a backstop outside the building. Design of most of the major components for this system is now finished. Completion of the installation is expected by the end of 1963. In addition to the provision for two new experimental areas utilizing possible focal points  $F_2$  and  $F_3$ , this system permits the solution of a major shielding problem because the backstop will be displaced beyond the presently most useful experimental areas. It is expected that secondary beams emerging from the backstop will play an important part in the future experimental program. The proposed backstop switching magnet will provide for simultaneous experimental setups in that area.

It is intended that the transport system remain reasonably flexible under experimental demands. The arrangement shown in Fig. 7 is achromatic between  $F_1$  and  $F_3$ . The performance of this system is illustrated schematically in

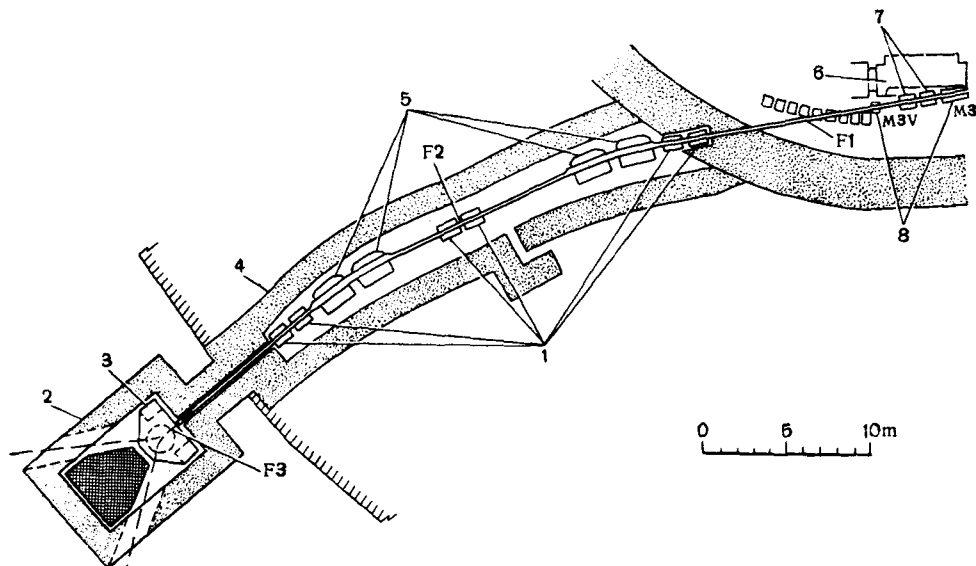


Fig. 7. The entire external system is shown without detail. Here  $F_1$ ,  $F_2$  and  $F_3$  are possible target locations for secondary beams. The system is now completed as far as the Bevatron main shielding wall; the remainder is under construction:

1 — quadrupoles  $Q_4$ ; 2 — back stop; 3 — switching magnet  $M_3$ ; 4 — shielding; 5 — bending magnets  $M_4$ ; 6 — west straight section; 7 — quadrupoles  $Q_3$ ; 8 — bending magnets.

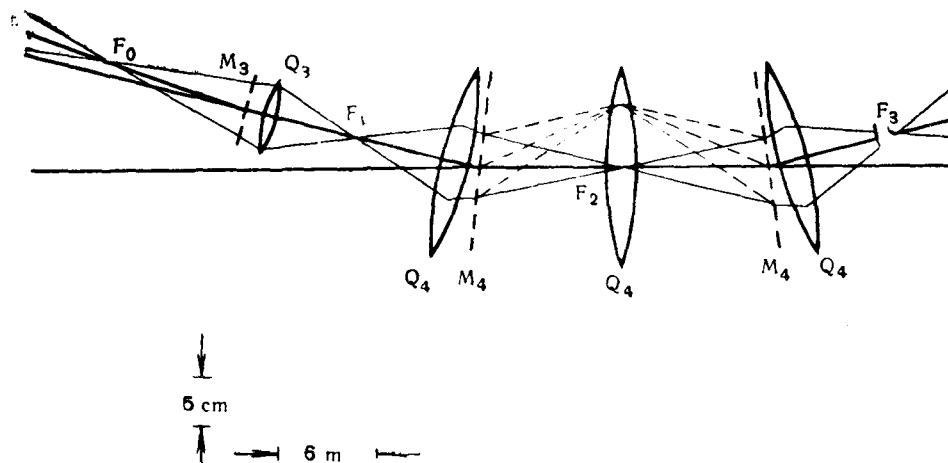


Fig. 8. Optical properties of the transport system. Here  $F_0$  is the second image of the energy-loss target which occurs inside the Bevatron. This arrangement of the deflecting and focussing magnets leads to a system which is achromatic between  $F_1$  and  $F_3$ . The dashed rays are for a momentum slightly below that for which the system is tuned.

Figs 8 and 9. A small momentum error, scattering in a target at  $F_1$ , or a deflection following  $F_1$ , do not necessarily degrade the image at  $F_3$  [4]. Because of the conservatively designed aperture (200 mm in diameter), a momentum

error of several percent does not lose the beam. For this reason it is believed that shielding along the transport system need not be extensive. The proposed 1.5-m-thick concrete shielding wall is formed almost entirely from old

blocks replaced during the reconversion of the Bevatron. The backstop itself will contain, in addition to the switching magnet, an iron core surrounded by concrete blocks in modules that facilitate experimental setup.

#### OPERATIONAL USE OF THE EXTERNAL PROTON BEAM

Maximum intensity for the extracted beam requires that the internal magnets be plunged

magnet is powered with a solid-state power supply, regulated with either silicon-controlled rectifiers or a magnetic-amplifier regulator supplemented with a fast transistor loop. Reproducibility in the current on a pulse-to-pulse basis is at least as good as one part in  $10^3$  for each magnet.

Targeting is expected to be the same for the external beam as for secondary beams, in the sense that the circulating beam can be divided among several targets during a given pulse.

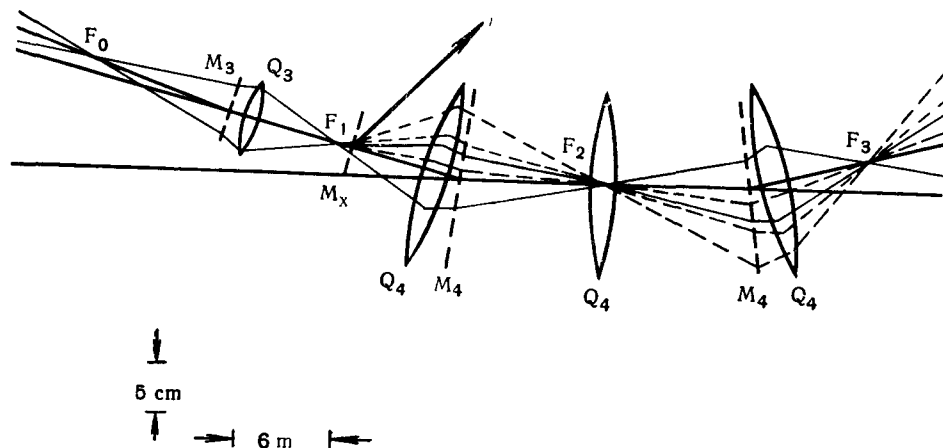


Fig. 9. Optical properties of the transport system. Scattering at  $F_1$  (or deflection by a small magnet  $M_x$  near  $F_1$ ) does not destroy the image at  $F_2$  and  $F_3$ . The angular error at  $F_2$  and  $F_3$  can be corrected, if desired, with another magnet like  $M_x$  located just ahead of  $F_2$ :  
I — secondary beam.

into the aperture after injection. This is done by two hydraulic plunging mechanisms which accurately position  $M_1Q_1$  and  $M_2Q_2$  in 0.7 s with strokes of 0.7 m. Hence full intensity in the extracted beam can be obtained for proton momenta down to 3 GeV/c. The beam has been extracted for momenta as low as 2 GeV/c with some loss of intensity because of the aperture restriction. A small fraction of the beam can be injected and accelerated with the internal magnets fixed in their final positions. To provide for extraction of the beam over a range of energies during a given Bevatron pulse, the radial position of  $M_1Q_1$  can be programmed mechanically to compensate for the change in the «jump» distance with energy.

The current in each magnet is pulsed. To a first approximation, the slope follows that of the magnetic field of the Bevatron, but small corrections to this program are available if needed. In this way the extraction system remains «tuned» over a range of momenta. Each

Operation with flattop and with either a long or short beam spill is feasible. Because of the orbit distortions caused by the field of the rapid beam deflector, the extraction efficiency for the fast spill is reduced (by a factor of order 2) with the present system.

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#### DISCUSSION

M. G. Meshcheryakov

What is the pulse repetition rate for the current of the accelerated protons in the bevatron in your latest experiments?

W. A. Wenzel

1. Eleven pulses per minute.
2. A telegram has come with new information from Berkeley. The Bevatron circulating beam intensity has been raised from  $2.3 \times 10^{12}$  protons per pulse as reported by Lambertson on tuesday to  $3.0 \times 10^{12}$  protons per pulse. The significant fact is that this happened as a consequence of some alterations in the radio frequency programming at injection.