

Figure 2: Internal view of the aluminum DTL.

The interfaces of the ports are the same as in the real DTL Tank, duplicating the actual geometry (i.e. diameter of the port, distance of the interfaces surfaces from the tank geometrical axis, position of the screw holes and sealing housing). This was done in order to carry out the bead pulling on the mock-up with the same movable tuner which will be used in the bead pulling of the DTL Tank in the DTL assembly workshop before the transfer and integration of the assembled and tested tank to the ESS accelerator.

The system for bead-pulling mainly consists of 2 aluminum plates directly fixed to the end plates of the DTL tank and supported to the ground by 2 adjustable aluminum bar. A step motor controlled by LABVIEW moves the dielectric thread through a system of pulleys, pulling a metallic bead along the beam axis line of the cavity. The same routine launches the Network Analyzer acquisition. The plates are qualified in CMM, and the precise positioning of the bead wire in the DTL geometrical center is provided by laser tracker references located on each plate. The tension of the thread can be varied by acting on a screw and it is monitored during the measurement by a load cell. The interfaces of this measurements apparatus were studied to be compatible with the real DTL tank geometry (Fig. 3).



Figure 3: Aluminum DTL with bead pull system, movable tuners and post couplers (courtesy of CERN-Linac4) and VNA.

RF MEASUREMENTS : STABILIZATION AND TUNING PROCEDURES

The stabilization procedures for the DTL foresee the consolidated method of the on-purpose perturbation of the cavity at one end of the structure and then to the other one; this provokes a frequency shift Δf and a field perturbation at each DTL cell. Therefore, it makes sense to introduce a Tilt Sensitivity parameter

$$TS_i = \frac{E_{0i}^{(he)} - E_{0i}^{(le)}}{\Delta f}$$

Where $E_{0i}^{(he)}$ and $E_{0i}^{(le)}$ refer to axial field values of the i -th cell in the case of perturbation at the high and low energy end respectively. This parameter is an index of the response of the cavity to perturbations and therefore the stabilization procedure consists of the determination of the set of PC lengths that minimize such parameter. In this case, the perturbing object is a hollow cylinder of 11 mm diameter, whose axis coincides with the beam axis and which is inserted in such a way to have the same Δf at both cavity ends (typically in the order of about -120 kHz). In the initial measurements, all the PCs were set at the length of 190 mm and with all the tuners set at 55 mm insertion. In this configuration, the TM and the PC mode bands are well separated (about 30 MHz) and therefore the TM modes only contribute to perturb the operational one. The TM_{010} frequency is equal to 352.1 MHz and the loaded Q is equal to about 2700, while the 1st upper mode is about 700 kHz higher than the TM_{010} one. The related TS is shown in Fig. 4. As one can expect, the TS is rather large (about $\pm 50 \text{ MHz}^{-1}$) and it is dominated by the 1st upper TM_{011} mode, moreover the E_0 varies in a range of $\pm 11\%$ about its mean value.

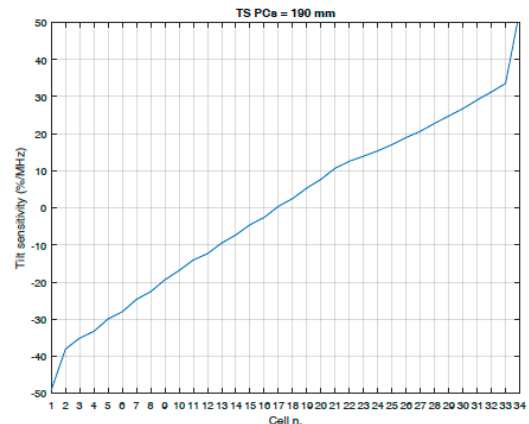


Figure 4: Tilt sensitivity for the initial PC setting.

At first, the optimum PC lengths were sought looking at the $TS'_i = TS_i - TS_{i-1}$, according to [3]: this procedure requires the determination of the resonance length of each PC looking to the discontinuity of resonance TM_{010} frequency. Nonetheless, in our case, the cavity 3dB bandwidth of about 240 kHz is too high to reliably distinguish such resonance. Therefore, it was decided to directly vary the PC lengths in order to reduce the TS, by trials and errors. This step took about 30 successive iterations. The

