STATUS OF OPERATION WITH NEGATIVE MOMENTUM COMPACTION AT KARA

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Abstract

For future synchrotron light source development novel operation modes are under investigation. At the Karlsruhe Research Accelerator (KARA) an optics with negative momentum compaction has been proposed, which is currently under commissioning. In this context, the collective effects expected in this regime are studied with an initial focus on the head-tail instability and the micro-bunching instability resulting from CSR self-interaction. In this contribution, we will present the proposed optics and the status of implementation for operation in the negative momentum compaction regime as well as a preliminary discussion of expected collective effects.

INTRODUCTION

Understanding collective effects and instabilities is critical to successful operation of synchrotron light sources and therefore they are subject to ongoing investigations. Most efforts focus on machines and optics with a positive momentum compaction factor, which is defined as

\[ \alpha_c = \frac{1}{L} \int DS(\sigma) \rho(\sigma) d\sigma, \] (1)

where \( L \) is the circumference of the ring, \( D_s(\sigma) \) is the dispersion and \( \rho(\sigma) \) the bending radius along the ring. Only a few experiments focus on a negative value of \( \alpha_c \) (e.g. [1–3]), however for a thorough understanding of collective effects and instabilities, studies in this regime are also necessary.

As KARA is a synchrotron light source also used as accelerator test facility, it is possible to change the optics and implement a negative momentum compaction factor. This allows to extend the investigations on collective effects and instabilities into the regime of a negative value of \( \alpha_c \).

LATTICE AND OPTICS

The injection energy into the KARA storage ring is 0.5 GeV. The storage ring itself is a ramping machine with a variable operation energy between 0.5 GeV and 2.5 GeV and a circumference of 110.4 m. The accelerating frequency is 500 MHz and the harmonic number is 184.

At KARA, two operation modes have already been established. Standard operation at 2.5 GeV with a moderate value of \( \alpha_c \approx 9 \times 10^{-3} \) and a short bunch mode at 1.3 GeV with a reduced value of \( \alpha_c \approx 1 \times 10^{-4} \) [4]. To reduce the value of \( \alpha_c \), the integral in Equation 1 and therefore the integrated dispersion has to be reduced. The dispersion and the beta functions are shown in Figure 1 for the normal operation and in Figure 2 for the short bunch operation. For the lower \( \alpha_c \) value the dispersion has negative parts, while for the normal operation it is strictly above zero. More information on different operation modes at KARA can be found in [5].

The method used for establishing a negative \( \alpha_c \) injection is described in the next chapter.

Figure 1: Calculated optics used for user operation at \( \alpha_c = 9 \times 10^{-3} \) and 0.5 GeV. The bottom depicts the magnets, quadrupoles in red, sextupoles in green and bends in blue.

Figure 2: Calculated optics for \( \alpha_c = 1 \times 10^{-4} \) at 0.5 GeV. In large parts the dispersion is negative.

Figure 3: Calculated optics for a negative value of \( \alpha_c = -8 \times 10^{-3} \). In large parts the dispersion is negative.
There have now been successful efforts to implement optics with a negative value of $\alpha_c$ at an energy of 0.5 GeV. To reach this the dispersion has to be even lower. Betafunctions and dispersion for the currently used negative alpha optics are shown in Figure 3. The dispersion is now stretched to almost $-2$ m in some areas, while for the low $\alpha_c$ optics it only reaches around $-1$ m.

**ESTABLISHING NEGATIVE ALPHA**

To establish operation with a negative momentum compaction factor the following strategy was applied.

As described in the previous chapter, in order to get a negative value of $\alpha_c$ the dispersion is stretched into negative values. Since this is also the way to reach a low positive value of $\alpha_c$, optics for a low and a negative value of $\alpha_c$ are similar to each other.

For the experiments, $\alpha_c$ is gradually reduced to a low value. Quadrupole strengths for a reduced value of $\alpha_c$ were simulated. Starting from an optics for standard injection at KARA (500 MeV, $\alpha_c \approx 8 \times 10^{-3}$) the strength of the third quadrupole family (Q3, in the center of the DBA structure), which has the most influence on $\alpha_c$, was changed stepwise in the direction of the values gained from the simulation.

To avoid resonances and stay at the well known stable transverse working point used for standard operation, at each step in Q3, the “field lens”, the tunes were adjusted using the other quadrupole families. This was done in two to four steps. First, families 2 and 4 were adjusted to keep the vertical tune. Second, the horizontal tune was adjusted using families 1 and 5. If necessary, the third step was to change families 2 and 4 to compensate the residual effect of families 1 and 5 on the vertical tune. In some cases this again made adjustments to families 1 and 5 necessary. That way it was possible to keep the tunes more or less constant during the decrease of $\alpha_c$ and also move the quadrupole strengths only in one direction to avoid hysteresis effects. The sextupole magnet strengths were adjusted according to the changes in Q3 starting from the standard injection settings.

After an optics with a low value of $\alpha_c$ was successfully established, it was chosen to extrapolate the magnet settings to negative values of $\alpha_c$ based on OPA [6] simulations as crossing $\alpha_c = 0$ did not seem feasible.

At this point, on each injection cycle (1 Hz) the beam was lost after a few hundred turns, which was visible on the first turn beam position data. Therefore kicker magnets, injection septum and some corrector magnets were adjusted. After successful injection and storage of the beam, the horizontal tunes were moved to the working point of standard operation at KARA, where injection efficiency and maximum beam current were improved.

**STATUS OF OPERATION**

Injection into negative values of $\alpha_c$ has been established. At present, the maximum current stored in the ring is limited and the highest achieved beam current was 17 mA while for the normal injection optics a maximum value of 150 mA is typical. Multiple tests have been performed and various factors have been identified that affect the limit.

One such factor is the horizontal injection orbit. Using the frequency of the accelerating voltage (RF frequency) the orbit was changed and the beam current limit was higher for larger orbit deviations. But above a certain threshold larger deviations resulted in little to no injection rate. The deviations of the currently used settings are shown in Figure 4. The maximum deviation is around 6.2 mm, which is large compared to the 3 mm maximum orbit deviations during normal user optics injection.

A second observation is that a reduced absolute value of $\alpha_c$ seemed to have a beneficial effect on the maximum injected current. $\alpha_c$ was changed from a value of $-1 \times 10^{-2}$ to $-8 \times 10^{-3}$. Due to this decrease the dispersion was stretched less and this resulted in smaller orbit deviations. Restoring the previous orbit by shifting the RF frequency improved the possible injected current which could indicate that a bigger momentum offset between injected and stored beam is beneficial. On the other hand modifying the bending magnets to change the energy of the stored relative to the injected beam did not result in the same effect.

Changes to the sextupoles also seemed to affect the total injected current. There were indications that lowering the vertical sextupole strengths and therefore lowering the vertical chromaticity increases the total possible beam current. Since the orbit deviations are quite large, the beam does not traverse the sextupoles in the center and therefore they add additional quadrupole components. Changes of the tune due to variations of the sextupole strengths were roughly compensated by adjusting the quadrupole strengths for these tests. Nonetheless, more systematic tests are necessary to determine the effect of the chromaticity on the maximum injectable beam current.

At $\alpha_c \approx -8 \times 10^{-3}$ the chromaticity $\xi_v$ has been measured for the vertical and horizontal plane and the results are shown in Figure 5. The horizontal chromaticity is $\xi_h = -0.460(\pm 0.138)$ and the vertical is $\xi_v = -6.624(\pm 0.484)$. Only a positive change in the momentum was carried out, due to the fact, that the orbit deviations were already quite big. Any larger deviations due to decreases in the RF frequency would have resulted in a beam loss.

**Figure 4: Orbit deviations during injection into optics with a negative value of $\alpha_c$. Large deviations in the horizontal plane are visible, with the largest value of roughly 6.2 mm, shortly after the injection septum.**

**Figure 5: Orbit deviations during injection into optics with a negative value of $\alpha_c$. Large deviations in the horizontal plane are visible, with the largest value of roughly 6.2 mm, shortly after the injection septum.**
The threshold from [9] (Equation 46 on p. 217) depends on the chromaticity and \( \alpha \). As the experiments with the negative vertical chromaticity \( \eta \) demonstrated. Using this threshold and the absolute value of the wake field, the threshold and behaviour are foreseen in the future. Simulations for the micro-bunching instability are planned and the applicability of Inovesa in the negative \( \alpha_c \) regime is under investigation.

**OUTLOOK**

There will be efforts to increase the current limit. Also better characterisations of the negative \( \alpha_c \) optics are necessary. In particular more investigations for the head-tail instabilities as well as measurements on the micro-bunching threshold and behaviour are foreseen in the future. Simulations for the micro-bunching instability are planned and the applicability of Inovesa in the negative \( \alpha_c \) regime is under investigation.

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