

Low Energy Particle Physics and Cosmology of Nonlinear Supersymmetric General Relativity

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Abstract

Basic ideas of nonlinear supersymmetric general relativity(NLSUSY GR) are explained and some cosmological and the low emergy particle physics consequences are discussed.

1 NLSUSY GR

By extending the geometric arguments of Einstein general relativity (EGR) on Riemann space-time to new space-time inspired by NLSUSY, where tangent space-time is specified not only by x_a for $SO(1, 3)$ but also by the Grassmanian ψ_a for isomorphic $SL(2C)$ of NLSUSY, the fundamental action (called nonlinear supersymmetric general relativity) has been constructed[1]:

$$L(w) = \frac{c^4}{16\pi G} |w| (\Omega(w) - \Lambda), \quad (1)$$

$$|w| = \det w^a_\mu = \det(e^a_\mu + t^a_\mu(\psi)), \quad t^a_\mu(\psi) = \frac{\kappa^2}{2i} (\bar{\psi}^i \gamma^a \partial_\mu \psi^i - \partial_\mu \bar{\psi}^i \gamma^a \psi^i), \quad (N = 1, 2, \dots, N), \quad (2)$$

where $w^a_\mu(x) = e^a_\mu + t^a_\mu(\psi)$, e^a_μ , $t^a_\mu(\psi)$ and $\Omega(w)$ are the invertible unified vierbein of new spacetime, the ordinary vierbein of EGR, the stress-energy-momentum of N NG fermion $\psi(x)$ (called *superons* as hypothetical spin 1/2 objects constituting all observed particles) and the the unified scalar curvature of new(SGM) spacetime, respectively. $s_{\mu\nu} \equiv w^a_\mu \eta_{ab} w^b_\nu$ and $s^{\mu\nu}(x) \equiv w^\mu_a(x) w^{\nu a}(x)$ are unified metric tensors of SGM spacetime. New space-time is the generalization of the compact isomorphic groups $SU(2)$ and $SO(3)$ for the gauge symmetry of 't Hooft-Polyakov monopole into the noncompact isomorphic groups $SO(1,3)$ and $SL(2C)$ for space-time symmetry. NLSUSY GR action possesses promissing large symmetries isomorphic to $SO(10)$ SP[2]. Note that the so called no-go theorem is overcome (circumvented) in a sense that the non-tivial N -extended SUSY gravity theory with $N > 8$ has been constructed in a NLSUSY invariant way.

2 Cosmology and Low Energy Physics

NLSUSY GR $L(w)$ on new *empty* space-time written in the form of the *vacuum* EH type is unstable due to NLSUSY structure of tangent space-time and decays (called *Big Decay* [3]) spontaneously into ordinary Einstein-Hilbert(EH) action with the cosmological constant Λ , NLSUSY action for N superons and their gravitational interactions on ordinary Riemann space-time written formally as follows, which ignites Big Bang of the present universe.

$$L(e, \psi) = \frac{c^4}{16\pi G} |e| \{R(e) - \Lambda + \tilde{T}(e, \psi)\}, \quad (3)$$

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which is called superon-graviton model(SGM) from composite viewpoint.

Considering SGM action reduces to N -extended NLSUSY action in asymptotic Riemann-flat ($e^a_\mu \rightarrow \delta^a_\mu$) space-time, we can study the low energy contents of SGM by constructing the N -extended LSUSY theory equivalent to N -extended NLSUSY model. We have shown explicitly by the heuristic arguments for simplicity in two space-time dimensions ($d = 2$) [4] that $N = 2$ LSUSY interacting QED is equivalent to $N = 2$ NLSUSY model. (Note that the minimal realistic SUSY QED in SGM composite scenario is described by $N = 2$ SUSY.)

$N = 2$ NLSUSY action for two superons ψ^i ($i = 1, 2$) in $d = 2$ is written as follows,

$$L_{N=2\text{NLSUSY}} = -\frac{1}{2\kappa^2} |w| = -\frac{1}{2\kappa^2} \left\{ 1 - i\kappa^2 \bar{\psi}^i \not{\partial} \psi^i - \frac{1}{2} \kappa^4 (\bar{\psi}^i \not{\partial} \psi^i \bar{\psi}^j \not{\partial} \psi^j - \bar{\psi}^i \gamma^a \partial_b \psi^i \bar{\psi}^j \gamma^b \partial_a \psi^j) \right\} \quad (4)$$

where $\kappa^2 = (\frac{c^4 \Lambda}{8\pi G})^{-1}$ in the SGM scenario and $|w| = \det(w^a_b) = \det(\delta^a_b + t^a_b)$, $t^a_b = -i\kappa^2 \bar{\psi}^i \gamma^a \partial_b \psi^i$. The most general massless $N = 2$ LSUSY QED action in $d = 2$ is written as follows,

$$\begin{aligned} L_{N=2\text{SUSYQED}} &= -\frac{1}{4} (F_{ab})^2 + \frac{i}{2} \bar{\lambda}^i \not{\partial} \lambda^i + \frac{1}{2} (\partial_a A)^2 + \frac{1}{2} (\partial_a \phi)^2 + \frac{1}{2} D^2 - \frac{1}{\kappa} \xi D \\ &+ \frac{i}{2} \bar{\chi} \not{\partial} \chi + \frac{1}{2} (\partial_a B^i)^2 + \frac{i}{2} \bar{\nu} \not{\partial} \nu + \frac{1}{2} (F^i)^2 + f(A \bar{\lambda}^i \lambda^i + \epsilon^{ij} \phi \bar{\lambda}^i \gamma_5 \lambda^j + A^2 D - \phi^2 D - \epsilon^{ab} A \phi F_{ab}) \\ &+ e \left\{ i v_a \bar{\chi} \gamma^a \nu - \epsilon^{ij} v^a B^i \partial_a B^j + \bar{\lambda}^i \chi B^i + \epsilon^{ij} \bar{\lambda}^i \nu B^j - \frac{1}{2} D (B^i)^2 + \frac{1}{2} (\bar{\chi} \chi + \bar{\nu} \nu) A - \bar{\chi} \gamma_5 \nu \phi \right\} \\ &+ \frac{1}{2} e^2 (v_a^2 - A^2 - \phi^2) (B^i)^2. \end{aligned} \quad (5)$$

We have shown explicitly[4]: $L_{N=2\text{NLSUSY}} + [\text{surface terms}] = L_{N=2\text{SUSYQED}}$. Therefore the low energy particle physics contents in asymptotic flat space-time of $N = 2$ SGM can be read from $N = 2$ SUSY QED action equivalent to $N = 2$ NLSUSY action.

Now we study the vacuum structure of $N = 2$ SUSY QED action (5). The vacuum is determined by the minimum of the potential $V(A, \phi, B^i, D)$,

$$V(A, \phi, B^i, D) = -\frac{1}{2} D^2 + \left\{ \frac{\xi}{\kappa} - f(A^2 - \phi^2) + \frac{1}{2} e (B^i)^2 \right\} D. \quad (6)$$

Substituting the solution of the equation of motion for the auxiliary field D we obtain

$$V(A, \phi, B^i) = \frac{1}{2} f^2 \left\{ A^2 - \phi^2 - \frac{e}{2f} (B^i)^2 - \frac{\xi}{f\kappa} \right\}^2 \geq 0. \quad (7)$$

The vacuum field configurations possess $\text{SO}(1,3)$ or $\text{SO}(3,1)$ isometries in (A, ϕ, B^i) -space depending upon the signatures of the parameters. One of the vacua in $\text{SO}(3,1)$ isometry for

$$ef < 0, \quad \frac{\xi}{f\kappa} > 0, \quad A^2 - \phi^2 + (\tilde{B}^i)^2 = k^2, \quad \left(\tilde{B}^i = \sqrt{-\frac{e}{2f}} B^i, \quad k^2 = \frac{\xi}{f\kappa} \right) \quad (8)$$

can be studied by substituting the following expressions

$$\begin{aligned} A &= -(k + \rho) \cos \theta \cos \varphi \cosh \omega, \quad \phi = (k + \rho) \sinh \omega, \\ \tilde{B}^1 &= (k + \rho) \sin \theta \cosh \omega, \quad \tilde{B}^2 = (k + \rho) \cos \theta \sin \varphi \cosh \omega, \end{aligned}$$

into the action and expanding the results around the vacuum. We obtain

$$\begin{aligned} L_{N=2\text{SUSYQED}} &= \frac{1}{2} \{ (\partial_a \rho)^2 - 4f^2 k^2 \rho^2 \} + \frac{1}{2} \{ (\partial_a \theta)^2 + (\partial_a \varphi)^2 - e^2 k^2 (\theta^2 + \varphi^2) \} - \frac{1}{4} (F_{ab})^2 \\ &+ \frac{1}{2} (\partial_a \omega)^2 - \frac{1}{4} (F_{ab})^2 + \frac{1}{2} (i \bar{\lambda}^i \not{\partial} \lambda^i - 2fk \bar{\lambda}^i \lambda^i) + \frac{1}{2} \{ i(\bar{\chi} \not{\partial} \chi + \bar{\nu} \not{\partial} \nu) - ek(\bar{\chi} \chi + \bar{\nu} \nu) \} + \dots \quad (9) \end{aligned}$$

and the following mass spectra with mass hierarchy by the factor $\frac{e}{f}$ indicating the spontaneous SUSY breaking as anticipated;

$$m_\rho^2 = m_{\lambda^i}^2 = 4f^2k^2 = \frac{4\xi f}{\kappa}, \quad m_\theta^2 = m_\varphi^2 = m_\chi^2 = m_\nu^2 = e^2k^2 = \frac{\xi e^2}{\kappa f}, \quad m_{v_a} = m_\omega = 0. \quad (10)$$

The local $U(1)$ gauge symmetry is not broken. The massless scalar ω is a NG boson for the degeneracy of the vacuum in (A, \tilde{B}_2) -space, which is gauged away provided the gauge symmetry between the vector and the scalar multiplet is introduced. As for the cosmological significances of $N = 2$ SUSY QED in SGM scenario, the above vacuum produces the same interesting predictions as already pointed out in $N = 2$ pure SUSY QED in SGM scenario [5], which may simply explain the observed mysterious (numerical) relations

$$((\text{dark}) \text{ energy density of the universe})_{\text{obs}} \sim 10^{-12} \sim (m_\nu)_{\text{obs}}^4 \sim \frac{\Lambda}{G} \sim g_{sv}^2,$$

and give a new insight into the origin of mass provided $f\xi \sim O(1)$ and λ^i is identified with neutrino. (Λ , G and g_{sv} are the cosmological constant of NLSUSY GR (SGM) on *empty* new space-time, the Newton gravitational constant and the superon-vacuum coupling constant via the supercurrent, respectively.) Here we just mention that there is another physical (apparently pathological) vacuum in $SO(3, 1)$ isometry, which contains off-diagonal fermion mass terms[6]. The similar investigations in $d = 4$ are urgent and the extension to large N , especially to $N = 5$, is important for *superon quintet hypothesis* in SGM scenario with $N = 10 = \underline{5} + \underline{5}^*$ [7]. Also NLSUSY GR in extra space-time dimensions is an interesting problem, which can describe all observed particles as elementary *a la* Kaluza-Klein.

Our analysis shows that the vacua of N -extended NLSUSY GR action in SGM scenario possess rich structures promising for the unified description of nature.

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References

- [1] K. Shima, *Physics Letters* B501 (2001) 237.
- [2] K. Shima and M. Tsuda, *Physics Letters* B507 (2001) 260.
- [3] K. Shima and M. Tsuda, *PoS HEP2005* (2006) 011;
- [4] K. Shima and M. Tsuda, *arXiv:0706.0063 [hep-th]* (2007).
- [5] K. Shima and M. Tsuda, *Physics Letters* B645 (2007) 455.
- [6] K. Shima, M. Tsuda and W. Lang, *arXiv:0710.1680[hep-th]* (2007).
- [7] K. Shima, *European Phys. J. C* 7 (1999) 341.