

## 2.24 Simulation Study of $K_L$ Beam: $K_L$ Rates and Background

Ilya Larin

*Department of Physics*

*Old Dominion University*

*Norfolk, VA 23529, U.S.A.*

### Abstract

We report our simulation results for  $K_L$ -beam and neutron background production, estimated rates for certain run conditions and resolution for  $K_L$ -beam momentum.

#### 1. $K_L$ Beam Line

Our calculations have been performed for Jefferson Lab Hall D setup geometry. Primary  $K_L$ -production target has been placed in Hall D collimator cave. For the target material, we selected beryllium as for thick targets  $K_L$ -yield roughly proportional to the radiation length and density, which gives beryllium as the best candidate. Beam plug and sweeping magnet are placed right after the target. For our calculations we took a simple beam plug: 15 cm thick piece of lead. Sweeping magnet is cleaning up charged component and has a field integral 2 Tesla . meter, which is enough to remove all charged background coming out of the beam plug. Vacuum beam pipe has 7 cm diameter and preventing neutron rescattering in air. There are two collimators: one placed before the wall between collimator cave and experimental hall, another - in front of the Hall D detector. Distance between primary Be target and liquid hydrogen ( $LH_2$ ) target (located inside Hall D detector) has been taken 16 m in our calculations, it can be increased up to 20 m.

#### 2. $K_L$ Production

We simulated  $K_L$ -production in photon bremsstrahlung beam produced by 12 GeV electron beam in Hall D tagger amorphous radiator. We analyzed  $K_L$ -production via  $\phi$ -meson photoproduction in detail. This is one of the main mechanisms of  $K_L$ -production at our energy range. It gives the same number of  $K^0$  and  $\bar{K}^0$ . Another mechanism is hyperon photoproduction (which gives only  $K^0$ ) was not studied in our simulations separately. Instead, we have taken as an alternative model Pythia generator [1], which includes hyperon production.  $\phi$ -meson photoproduction total and differential cross sections on proton and complex nuclei (coherent and incoherent) were taken from Refs. [2,3]. Angular distributions for  $\phi \rightarrow K_L K_S$  decay, we used are from Ref. [2,4,5]. These calculations show that  $\phi$  decay in its rest frame is going mostly perpendicular to the axis of  $\phi$  momentum. Since  $K_L$ s need to stay on original photon beam direction to get  $LH_2$  target, this condition requires that  $\phi$  production and decay angles in laboratory frame should be about the same. That means we will have in the  $LH_2$  only  $K_L$ s from  $\phi$ -mesons produced at relatively high  $t$ . It suppresses the number of "useful"  $K_L$ s by factor of  $\sim 3$  or more (in comparison with the case if  $K_L$  and  $K_S$  momenta are parallel to  $\phi$  momentum).  $K_L$  absorption used in our calculations was studied in Ref. [6] very well. About 80% of produced  $K_L$ s will be absorbed in Be target itself and beam plug. This value of absorbed  $K_L$ s can be reduced by optimizing beam plug setup.

### 3. $K_L$ Beam Parameters

One of the main  $K_L$ -beam parameters is momentum distribution. Momentum spectrum is a function of the distance and angle. We are giving here resulting spectra for  $K_L$  reaching  $\text{LH}_2$  target. Results of our simulations for  $K_L$  momentum spectrum is shown on Fig. 1. The spectrum first has increasing shape since  $\phi$  decay cone angle decreasing at higher  $\gamma$ -beam and  $K_L$  momentum. This selecting lower  $\phi$  production  $t$  values, which are more favorable according to  $\phi$  differential cross section. At certain point highest possible  $\gamma$ -beam momentum is reached and  $K_L$  momentum spectrum is dying out pretty fast. For comparison, we selected part of  $K_L$  spectrum from Pythia generator originated only from  $\phi$  decays and showed it on the same plot (red histogram).

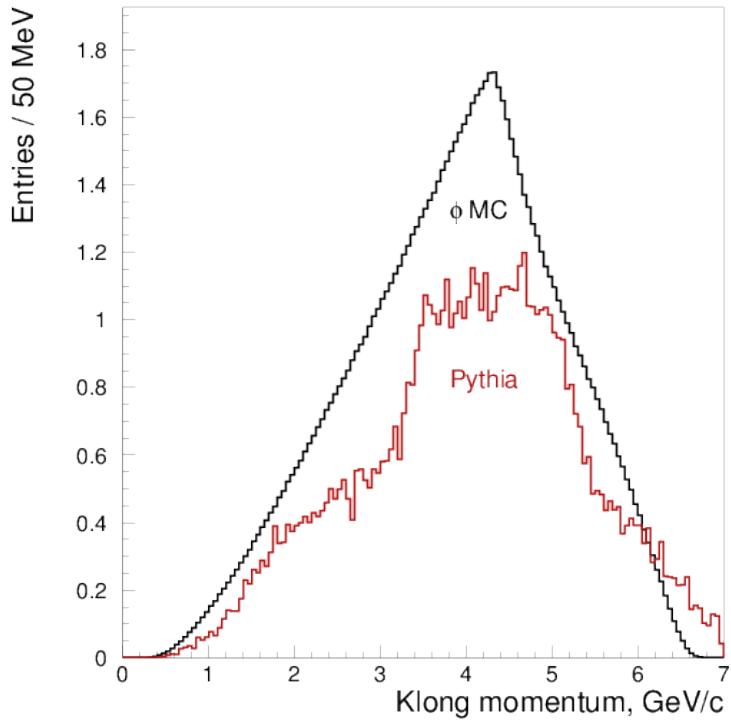


Figure 1:  $K_L$  momentum spectra originated from  $\phi$  decays: black histogram - our simulation using GEANT [7], red - Pythia generator result [1].

Pythia shows, that  $\phi$  decays give about 30% of  $K_L$ s. Number of  $K^0$  exceeds number of  $\bar{K}^0$  by 30% according to this generator for our conditions. Their momentum spectra are shown on Fig. 2 separately.

To estimate expected rate of  $K_L$  at  $\text{LH}_2$  target we used the following conditions: electron beam current is  $3.2 \mu\text{A}$ , tagger radiator thickness is 1% of radiation length, Be target thickness is 40 cm, distance Be to  $\text{LH}_2$  target is 16 m, radius of  $\text{LH}_2$  target is 2 cm. Our calculations are related to the  $K_L$  flux at that distance and solid angle. For  $K_L$ -beam intensity under the above condition, we got 100  $K_L$ s per second for our  $\phi$  photoproduction simulation and

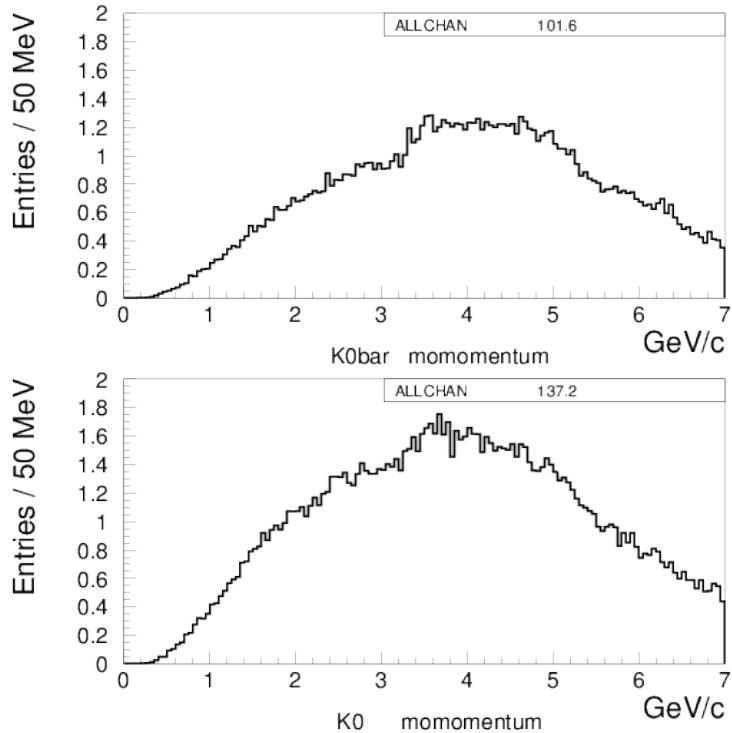


Figure 2:  $\bar{K}^0$  (top plot) and  $K^0$  (bottom plot) momentum spectra from Pythia generator.

240  $K_L$ s per second from all sources from Pythia. There are ways to increase the  $K_L$ -beam intensity by increasing tagger radiator thickness, electron beam current and other parameters. Increasing  $\text{LH}_2$  target radius will increase number of  $K_L$ s reaching it proportionally to the solid angle. For example for  $\text{LH}_2$  target radius 4 cm, electron beam current 5  $\mu\text{A}$ , 5% rad. length radiator and increased Be target sizes we shall be able to obtain beam rate about 7,000  $K_L$ s per second from all production mechanisms at  $\text{LH}_2$  target face. For comparison this value corresponds to  $\sim$ 10 million of produced  $K_L$ s in Be target per second.

#### 4. $K_L$ Beam Resolution

$K_L$ -beam momentum can be measured using TOF - time between accelerator bunch and reaction in  $\text{LH}_2$  target detected by start counter. Hall D tagger timing can not be used at such high intensity conditions. Thus TOF resolution is a quadratic sum of accelerator time and start counter time resolutions. Since accelerator signal has very good time resolution ( $\sim$ 0.1 ns or better), TOF resolution will be defined by start counter. Hall D start counter has resolution  $\sim$ 0.35 ns. This value can be hopefully improved with upgrading counter design and parameters. In our calculations we used an optimistic value of 0.25 ns start counter time resolution. Of course to get TOF information electron beam needs to have narrow bunch time structure with the distance between bunches at least 30 ns. At low ( $< 1 \text{ GeV}/c$ )  $K_L$  momenta uncertainty in  $K_L$  production point position within Be target will also affect TOF

calculation precision. Fig. 3 shows TOF and beam momentum resolution as a function of  $K_L$  beam momentum.

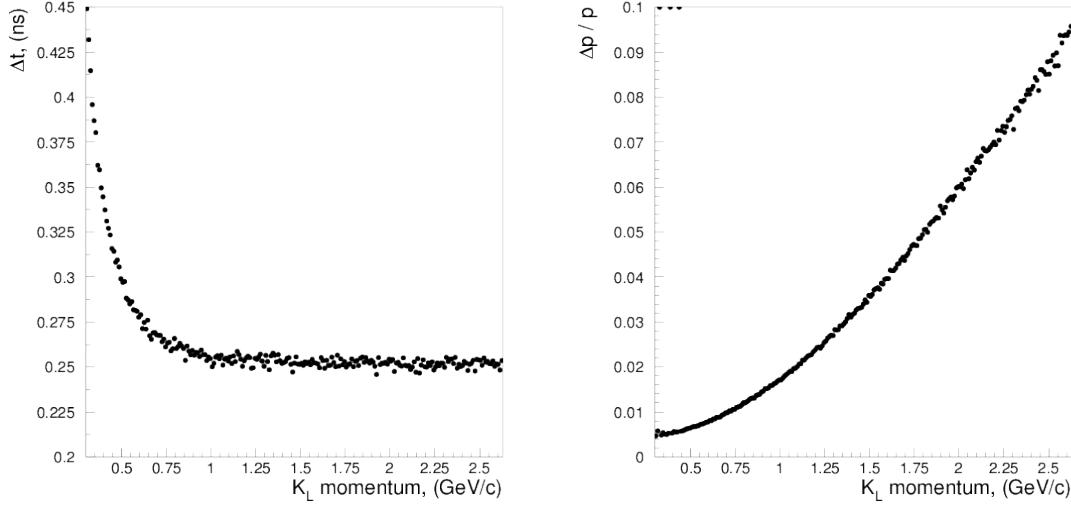


Figure 3:  $K_L$ -beam TOF (left plot) and momentum (right plot) resolution as a function of momentum.

TOF resolution is flat for momenta higher than 1  $\text{GeV}/c$ . Momentum resolution is growing with momentum value, for 1  $\text{GeV}/c$  it is  $\sim 1.7\%$ , for 2  $\text{GeV}/c \sim 6\%$ .

## 5. $K_L$ Beam Background

Background conditions is one of the most important parameter of the beam. After passing through 30 radiation length beam plug and swiping out charged background component, we will have some residual  $\gamma$  background and neutrons produced by electromagnetic showers. Momentum spectrum of residual  $\gamma$ s shown on Fig. 4.

It decreases exponentially with increasing  $\gamma$  energy. For the rates we obtained  $\sim 100,000$  per second for  $\gamma$ s with energy above 50 MeV and  $\sim 1,000$  per second for  $\gamma$ s above 500 MeV.

The most important and unpleasant background for  $K_L$ -beam is neutron background. Special care needs to be taken to estimate and if possible to eliminate this kind of background. In our calculations to estimate neutron background we used two independent program packages: Pythia [1] and DINREG [8]. Both packages give the same order of magnitude neutron background level. At our condition it is  $\sim 140$  neutrons per second at  $\text{LH}_2$  target for neutrons with momenta higher than 500  $\text{MeV}/c$ . These spectra along with  $K_L$  momentum spectrum are shown on Fig. 5.

Additionally we calculated muon production level. Muon will be swiped out of the beam line thus they are not our background. But since their high penetration ability it might be important the for purposes of the shielding. We taken into account only Bethe-Heitler muon production process. Muons from pion decays and other production mechanisms will increase total muon yield only slightly. They were not included in our model. Number of produced

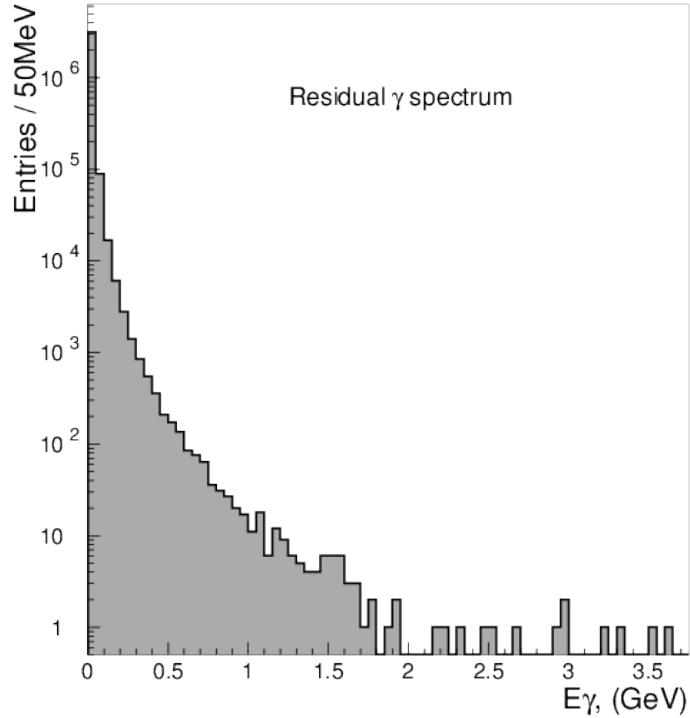


Figure 4: Momentum spectrum of residual  $\gamma$ s.

muon in Be target and lead beam plug is about the same, lead originating muons have much softer momentum spectrum. Estimated number of produced muons  $\sim$ 6 million per second. Their momentum spectrum is shown on Fig. 6.

Half of muons have momentum higher than  $2 \text{ GeV}/c$ ,  $\sim 10\%$  of muons have momentum higher than  $6 \text{ GeV}/c$  and  $\sim 1\%$  of muons - momentum above  $10 \text{ GeV}/c$ .

## 6. Summary

In the summary part, we want to emphasize that  $K_L$ -beam facility opens horizons for new rich physics. Jefferson Lab GlueX spectrometer has very good acceptance and resolution parameters [9], which perfectly fit  $K_L$ -beam facility requirements. Expected rates for  $K_L$ -beam with increased  $\gamma$ -beam luminosity will allow to collect statistics order of magnitude higher than other facilities can provide. One of the main advantage of such facility is that  $K_L$ -beam is produced by  $\gamma$ -beam which provides low neutron background level comparing with hadron produced  $K_L$ -beam. To get more precise  $K_L$ -beam rates and neutron background estimations as well as radiation levels induced, it is important to conduct a few days measurements on low intensity test beam.

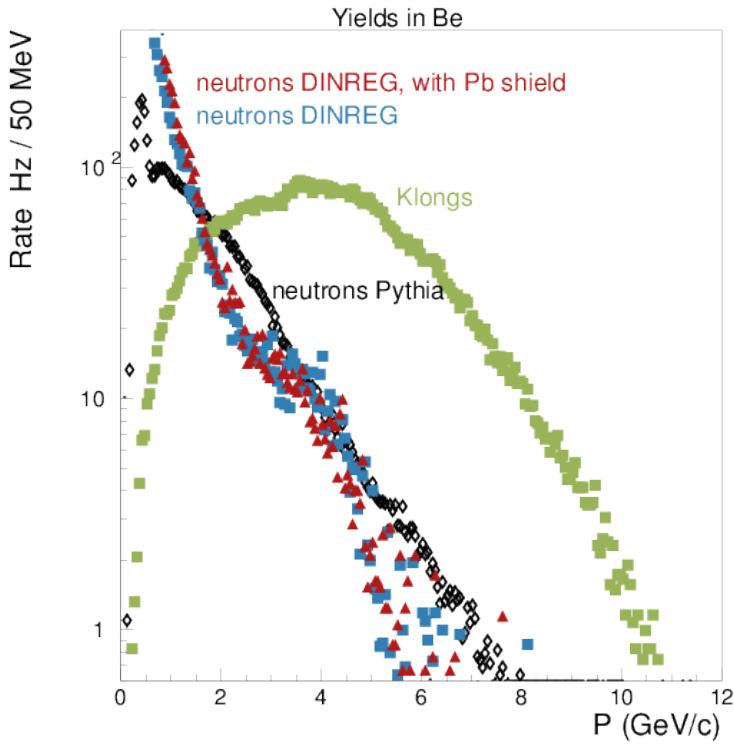


Figure 5:  $K_L$  and neutron momentum spectra obtained with different packages.

## 7. Acknowledgments

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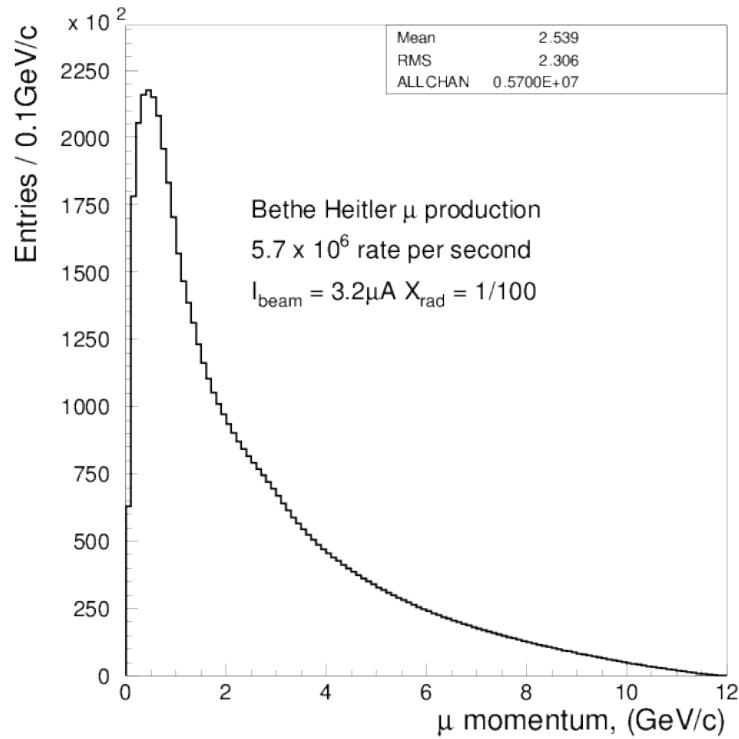


Figure 6: Muon momentum spectrum for Bethe-Heitler production.

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