

## Study of hot giant dipole resonance in medium mass nuclei

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### Introduction

The isovector giant dipole resonance (IVGDR) is a fundamental high-frequency mode of nuclear collective excitation that decays by emitting  $\gamma$ -rays [1]. Because these  $\gamma$ -rays are minimally affected by nuclear surroundings, it serves as a clean probe for studying hot nuclear systems. The IVGDR can be characterized by three key parameters: centroid energy ( $E_G$ ), resonance width ( $\Gamma_G$ ), and strength ( $S_G$ ). Among these, the IVGDR width ( $\Gamma_G$ ) is particularly significant because it directly relates to various damping mechanisms inside nucleus. Over the past decades, experiments have shown that  $\Gamma_G$  increases with temperature in the range of approximately  $1 \text{ MeV} \lesssim T \lesssim 3 \text{ MeV}$ , with some debate over whether it saturates at higher temperatures beyond  $T > 3 \text{ MeV}$ . At lower temperatures ( $T \lesssim 1 \text{ MeV}$ ), the situation is more complex due to strong microscopic effects, such as pairing correlations and shell effects, which obscure the thermal broadening of  $\Gamma_G$ . Alongside experimental studies, various theoretical models have been developed to explain how  $\Gamma_G$  behaves with temperature. Among these, the classical thermal shape fluctuation model (TSFM) is most popular.

However, validating these models requires data over a broad mass range. Most existing experimental data focus on nuclei with  $A > 90$  and cover only limited temperature ranges. To address this gap, we conducted an experimental study in mid-mass nuclei. We analyzed high-energy  $\gamma$ -ray spectra from the excited  $^{62}\text{Zn}$  compound nucleus, which was populated through light-ion induced fusion reaction, and measured the IVGDR width at different nuclear temperatures. We then compared our findings with TSFM prediction.

### Experimental details and data analysis

The experiment was conducted at VECC in Kolkata [3]. Nuclei was populated by bombarding  $^{58}\text{Ni}$  target with accelerated alpha beams of energy 28 MeV and 40 MeV from the K-130 RTC. The emitted high-energy  $\gamma$  rays (with energies  $E_\gamma > 4 \text{ MeV}$ ) were detected using the LAMBDA high-energy  $\gamma$ -ray spectrometer.

To determine the total energy deposited by  $\gamma$ -rays within the detector volume, nearest-neighbor cluster summing algorithm was employed [2]. Before summing, various unwanted events had to be identified and excluded to isolate the IVGDR bump in 10-25 MeV energy region, which is superimposed on the statistical  $\gamma$ -ray distribution. Major contamination in the high-energy  $\gamma$ -ray spectra, primarily

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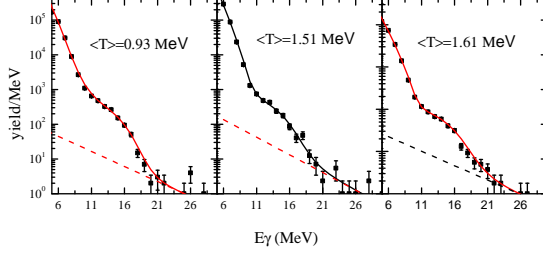


FIG. 1: Measured  $\gamma$ -ray spectra (symbols) from the decay of  $^{62}\text{Zn}$  is compared with the statistical model results (red solid lines). The Bremsstrahlung contributions are shown by the red dashed lines.

from neutron background, was identified using time-of-flight technique. Pile-up events were rejected through pulse shape discrimination, and cosmic muon events, which have distinct hit patterns compared to high-energy gamma particles, were effectively identified and eliminated. Events that passed all these checks were considered to be valid.

To extract the IVGDR parameters, measured  $\gamma$ -ray spectra were compared with statistical model calculations in the Hauser-Feshbach formalism using TALYS-1.95. This model predicts decay through various channels based on excitation energy, with probabilities determined by transmission coefficients ( $T$ ) and final state level densities. For the decay through  $\gamma$ -ray emission,  $T(E_\gamma)$  is related to the energy dependent photon strength function, represented by a standard Lorentzian function:

$$f_{E1}(E_\gamma) = \frac{2\sigma_{TRK}S_G}{3\pi^3(\hbar c)^2} \frac{E_\gamma\Gamma_G}{(E_\gamma^2 - E_G^2)^2 + E_\gamma^2\Gamma_G^2}, \quad (1)$$

where  $\Gamma_G$  is the resonance width,  $E_G$  is the centroid energy, and  $S_G$  is the fraction of the TRK dipole sum rule strength exhausted by the IVGDR. The best-fit values of these parameters were obtained through visual inspection.

## Results and discussion

In Fig.1, we compare measured spectra with the best-fit statistical model calculations.

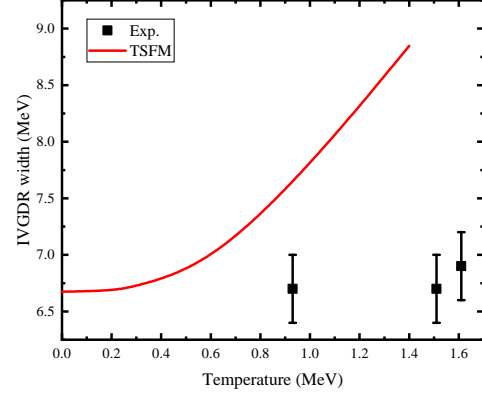


FIG. 2: Variation of the measured IVGDR width with temperature. The solid red line represents the TSFM calculation.

Fig.2 shows the extracted width as a function of temperature, indicating a flat width-temperature behavior, which suggests compactness of the system. The line in Fig.2 represents the TSFM calculation, performed with macroscopic-microscopic potential energy surface, but it fails to account for the low-temperature suppression of  $\Gamma_G$ . This indicates the need for improved models with additional microscopic contributions. Our current results are preliminary, relying on visual fitting of independently varied parameters. A more rigorous approach, such as Bayesian inference, is needed for final results, and this work is ongoing.

## References

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