

# Studying the dependence of transverse spherocity on the initial and final state anisotropies in heavy-ion collisions

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## Introduction

In non-central heavy-ion collisions, the overlap region of the colliding nuclei is oval-shaped and possesses event-by-event fluctuation in the nucleon density distribution. These initial profiles for initial nucleon distribution are quantified by eccentricity ( $\varepsilon_2$ ) and triangularity ( $\varepsilon_3$ ). For a hydrodynamically evolving system, if the pressure gradient formed during the collision of the nuclei is large enough, initial spatial eccentricities can be transformed into final state momentum space azimuthal anisotropies, quantified by the coefficients of the Fourier expansion of the invariant yield, *viz.*, elliptic flow ( $v_2$ ) and triangular flow ( $v_3$ ).

On the other hand, transverse spherocity ( $S_0$ ), an event shape observable, is known for its excellent capability for separating the jetty events dominated by the hard interactions from the isotropic events dominated by the soft-QCD interactions. Recent studies show that the event selection based on  $S_0$  can achieve higher radial flow velocity and lower kinetic freezeout temperatures compared to when no  $S_0$  selection is made [1]. Here, we use transverse spherocity as an event classifier to see its dependence on  $\varepsilon_2$ ,  $\varepsilon_2$ ,  $v_2$  and  $v_3$  in Pb–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV using a multi-phase transport model (AMPT). In this study, the centrality selection is performed using geometrical impact parameter slicing. For each centrality class, 20% events having the highest and lowest  $S_0$  values are termed as high- $S_0$  and low- $S_0$ , respectively [2].

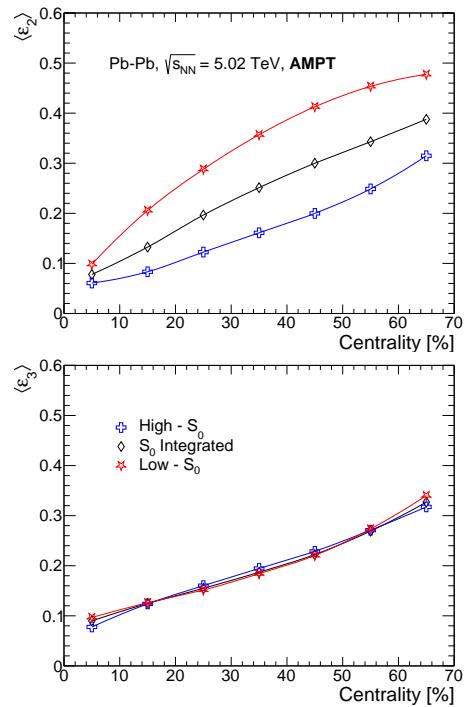


FIG. 1: Eccentricity ( $\varepsilon_2$ ) and triangularity ( $\varepsilon_3$ ) versus centrality for different transverse spherocity selections in Pb–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV using AMPT [2].

## Results and Discussions

Figure 1 shows  $\varepsilon_2$  and  $\varepsilon_3$  as a function of collision centrality and transverse spherocity in Pb–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV using AMPT. Both  $\varepsilon_2$  and  $\varepsilon_3$  gradually increase as one moves from the central to the peripheral collisions as the collision overlap region gets more anisotropic. Additionally,  $\varepsilon_2$  shows reasonable dependence on  $S_0$ , while  $\varepsilon_3$  is independent of  $S_0$  selection as it is caused by event-by-event density fluctuations and, in principle,

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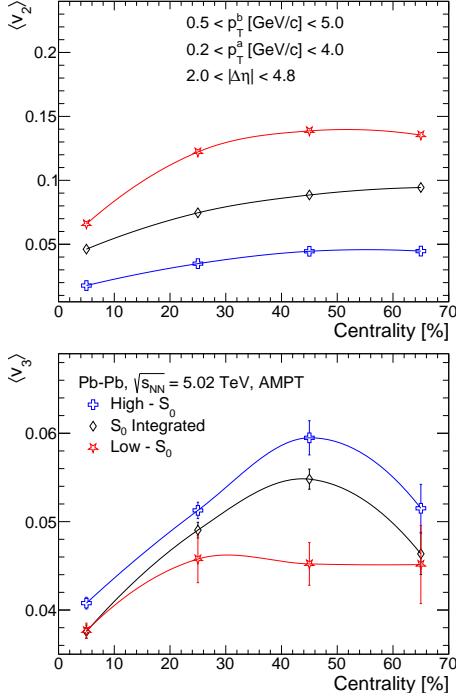


FIG. 2: Elliptic flow ( $v_2$ ) and triangular flow ( $v_3$ ) versus centrality for different transverse spherocity selections in Pb–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV using AMPT [2].

event-by-event fluctuations should not be affected by selections based on transverse spherocity. As high- $S_0$  events represent isotropic events, the corresponding value of  $\varepsilon_2$  is expected to be lower, as shown in Fig. 1.

Figure 2 shows  $v_2$  and  $v_3$  versus collision centrality and transverse spherocity in Pb–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV using AMPT. Both  $v_2$  and  $v_3$  are estimated using the two-particle correlation method, the methodology and corresponding track selection cuts can be found in Ref. [2]. Here, both  $v_2$  and  $v_3$  increase from the most central to the mid-central collisions, where the rate of increase gradually decreases, and the  $v_3$  value starts decreasing again towards the peripheral collisions. This is because, as one moves from most-central to mid-central collisions, the anisotropy increases, and as a result, both  $v_2$  and  $v_3$  increase. However, towards the peripheral colli-

sions,  $\varepsilon_2$  and  $\varepsilon_3$  cannot transform into  $v_2$  and  $v_3$  due to fewer participants to carry the information [2]. Additionally, both  $v_2$  and  $v_3$  show a significant dependence on event selections based on  $S_0$ , low- $S_0$  events have higher  $v_2$  and smaller  $v_3$  compared to the  $S_0$  integrated case. While  $v_2$  shows a similar  $S_0$  dependence as  $\varepsilon_2$ , the manifestation of finite transverse spherocity dependence on  $v_3$  when  $\varepsilon_2$  is independent of  $S_0$  is a matter of importance. Another crucial point to note is that while  $v_2$  is found to be anti-correlated with  $S_0$  selection,  $v_3$  is observed to be positively correlated. This behaviour between  $v_2$  and  $v_3$  is also observed in experiments with respect to a different event shape classifier, the reduced flow vector [3].

## Summary

In summary, we have studied the initial azimuthal anisotropies, such as eccentricity ( $\varepsilon_2$ ) and triangularity ( $\varepsilon_3$ ), and final state azimuthal anisotropies, *viz.*, elliptic flow ( $v_2$ ) and triangular flow ( $v_3$ ) as a function of centrality and transverse spherocity.  $\varepsilon_2$ ,  $v_2$  and  $v_3$  show appreciable spherocity dependence, whereas  $\varepsilon_2$  do not show any spherocity dependence. Both  $\varepsilon_2$  and  $v_2$  show anticorrelation with  $S_0$  while  $v_3$  shows a positive correlation.

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## References

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