

ELASTIC K^+ -p SCATTERING AT 0.97, 1.17, AND 1.97 $\text{GeV}/c^{(*)}$

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I. INTRODUCTION

Previous experiments on elastic K^+ -p scattering have been carried out at momenta up to 810 MeV/c by using bubble chambers ¹⁾, emulsions ²⁾, and counters ³⁾. In this paper we report an experiment on K^+ -p elastic scattering at higher momenta, using spark chambers to measure $d\sigma/d\Omega$ at angles greater than about 10° (lab) and counters for the small-angle scattering. The characteristics of spark chambers that provide a crucial advantage over other devices for this type of experiment are: (a) time resolution (about 500 nsec) which makes possible the use of electronically separated beams, rather than the spatially separated beams needed for bubble chambers; and (b) spatial resolution enabling the measurement of angles with an accuracy of about $1/2^\circ$, and a fairly large accepted solid angle—a combination which is very difficult to achieve with counters alone.

II. EXPERIMENTAL METHOD AND RESULTS

The variable-momentum K -particle beam was selected by using two high-pressure gas Čerenkov counters and time-of-flight, in a secondary beam from the Bevatron.

Scattered particles were detected either by the hodoscope or the spark-chamber system. The arrangement of the scattering detector systems is shown in Fig. 1.

The hodoscope system consisted of a bending magnet and gas Čerenkov counter to measure simultaneously the momentum and velocity of the scattered K meson. For scattering angles greater than $\theta_K^{\text{lab}} \approx 10^\circ$

the recoil proton can escape from the liquid-hydrogen target and be detected. These events were recorded by triggering three spark chambers placed around the target. The chambers had rectangular thin foil plates and were filled with argon.

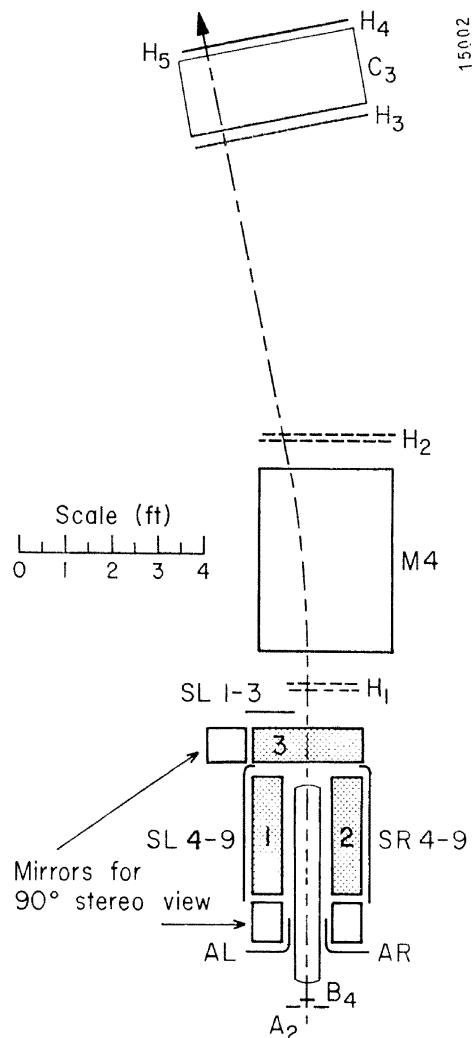


Fig. 1 Arrangement of scattering detector apparatus.

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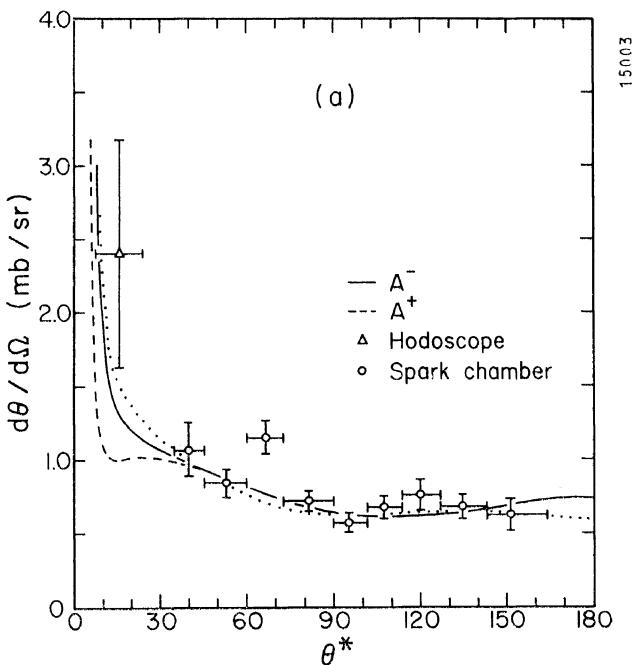


Fig. 2 a

Elastic scatters were selected on the basis of coplanarity and the polar scattering angles of the two particles. About 1500 events were found at each of the three momenta. The differential cross sections are shown in Fig. 2.

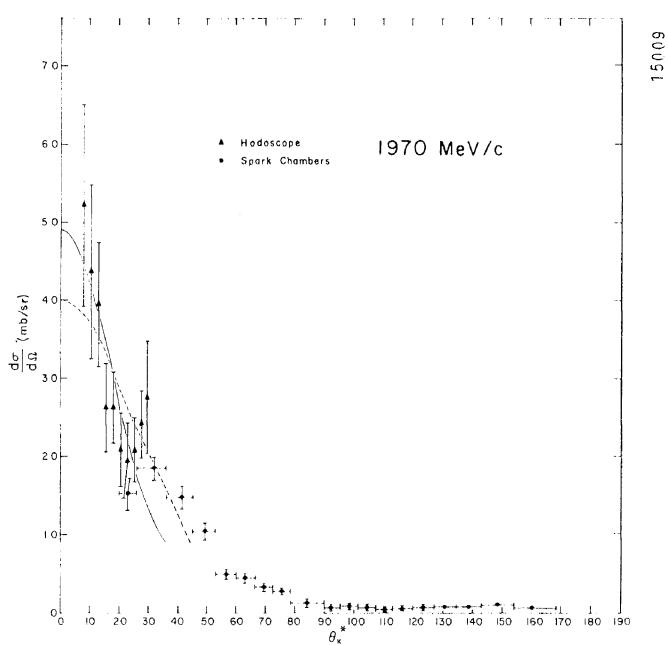


Fig. 2 c

Fig. 2 Measured angular distributions. The curves on the 970- and 1170-MeV/c data were calculated for phase-shift solutions set (a) given in Table I. The A^+ solutions are approximately the same as the A^- sets with the signs of all the phase-shifts reversed. The dotted curve on the 970 MeV/c plot was derived from the A^- solution for phase shift set (b).

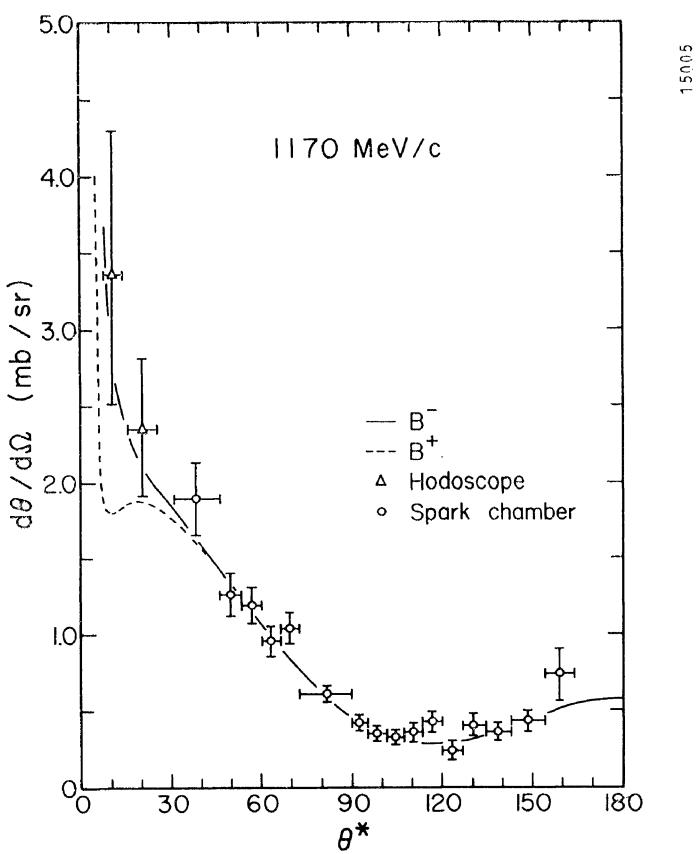


Fig. 2 b

III. DISCUSSION

(a) Phase-Shift Analysis

Elastic $K^+ - p$ angular distributions have been measured previously up to 810 MeV/c, and all are consistent with pure S -wave scattering. In addition to an S -wave solution, Stubbs *et al.*¹⁾ also found satisfactory fits at 810 MeV/c with dominant $P_{1/2}$ and with a mixture of $P_{1/2}$ and $P_{3/2}$ amplitudes.

Our angular distributions show a deviation from isotropy that increases with energy. At 970 MeV/c, the inelastic cross-section is one-third of the total, as compared with one-tenth of the total at 810 MeV/c (see Fig. 3).

Phase-shift analyses including (a) complex $S_{1/2}$, $P_{1/2}$, and $P_{3/2}$ phase shifts, and (b) real $S_{1/2}$, $P_{1/2}$ and $P_{3/2}$ and complex $D_{3/2}$ phase shifts were made for the 970- and 1170-MeV/c data. Set (b) would be expected to provide a good fit to the data if $N^*(3/2, 3/2)$ production was the dominant inelastic process, and the final orbital angular momentum state was S wave. The 1170 MeV/c data were also fitted with (c) complex $S_{1/2}$, real $P_{1/2}$ and $P_{3/2}$, and complex $D_{3/2}$ phase shifts. Set (c) would be expected to pro-

vide a good fit to the data if K^* production was the dominant inelastic process and the K^*-N system were in an S state. The situation is more complicated than at 810 MeV/c because of the large inelastic

cross-section (see Fig. 3), and the number of satisfactory solutions is accordingly larger. Some representative solutions are given in Table I.

In spite of the multiplicity of phase-shift solutions that were found, we can make the following general observations:

(i) The S -wave contribution to be elastic scattering is decreasing rapidly at 1 GeV/c.

(ii) Although the various solutions predict quite different values for the phase shifts, they lead to a more or less stable estimate of the magnitude of the real part of the forward scattering amplitude.

(iii) The SPD solutions for the 970 MeV/c-data are somewhat better than the SP solutions. This supports the N^* ($3/2, 3/2$) production hypothesis discussed above.

At 1970 MeV/c the angular distribution shows a pronounced diffraction peak upon a small almost isotropic distribution of about 0.05 mb/sr.

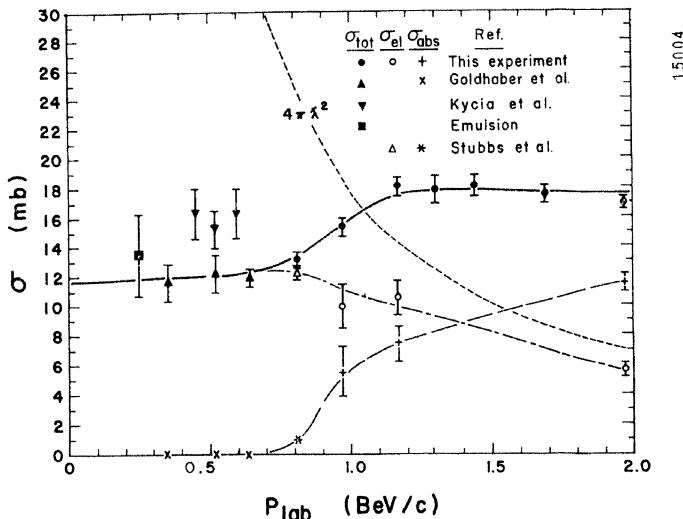


Fig. 3 The $K^+ - p$ total, elastic, and inelastic cross-sections in the momentum range $p_{\text{lab}} = 0$ to 2 GeV/c. The curves through the measured cross-sections have no theoretical significance.

TABLE I

Phase-shift solutions found at 970 MeV/c and 1170 MeV/c for phase-shift sets (a), (b), and (c)

970 MeV/c								
Set (a)	δ_0	η_0	δ_1^-	η_1^-	δ_1^+	η_1^+	$D(f)$	$P(\chi^2)$
A^-	-38 ± 1	1.0	-10 ± 1	1.0	3 ± 4	0.6	-0.10 ± 0.03	0.14
B^-	9 ± 1	0.7	15 ± 7	1.0	-29 ± 2	0.8	-0.12 ± 0.04	0.13
Set (b)								
	δ_0	δ_1^-	δ_1^+	δ_2^-	η_2^-		$D(f)$	$P(\chi^2)$
A^-	-29 ± 3	-20 ± 2	9 ± 2	4 ± 1	0.6		-0.13 ± 0.05	0.35
B^-	-2 ± 4	13 ± 2	-25 ± 2	6 ± 4	0.6		-0.18 ± 0.05	0.41
1170 MeV/c								
Set (a)	δ_0	η_0	δ_1^-	η_1^-	δ_1^+	η_1^+	$D(f)$	$P(\chi^2)$
A^-	-33 ± 2	1.0	-10 ± 19	0.8	4 ± 1	0.2	-0.19 ± 0.09	0.22
B^-	-3 ± 1	0.2	-64 ± 14	0.3	-15 ± 1	0.9	-0.20 ± 0.09	0.22
Set (b)								
	δ_0	δ_1^-	δ_1^+	δ_2^-	η_2^-		$D(f)$	$P(\chi^2)$
A^-	-25 ± 7	-28 ± 1	4 ± 5	-11 ± 5	0.3		-0.26 ± 0.10	0.20
B^-	-12 ± 9	-28 ± 2	1 ± 6	-31 ± 9	0.2		-0.25 ± 0.10	0.25
Set (c)								
	δ_0	η_0	δ_1^-	δ_1^+	δ_2^-	η_2^-	$D(f)$	$P(\chi^2)$
A^-	-32 ± 10	0.6	6 ± 2	-25 ± 3	10 ± 3	0.6	-0.24 ± 0.10	0.15
B^-	-11 ± 10	0.2	-38 ± 8	-2 ± 4	1 ± 7	0.6	-0.19 ± 0.10	0.27

(b) Forward Scattering Dispersion Relations

Using the new data reported here we have examined the forward scattering K - p dispersion relations to see: (a) how well all the existing information can be described by these relations; and (b) how well the sign and magnitude of the effective pole term (including Λ , Σ , Y_1^* , and, possible, Y_0^* poles) can now be estimated.

Our approach has been to try to fit all of the available K^+ - p and K^- - p data by using a singly subtracted dispersion relation

$$D(\omega) - D(\omega_0) = \Gamma \left(\frac{1}{\bar{\omega} + \omega} - \frac{1}{\bar{\omega} + \omega_0} \right) + f(\omega) - f(\omega_0),$$

which we write in the form

$$y = \Gamma(x - x_0) + D(\omega_0),$$

where

$$x = \frac{1}{\bar{\omega} + \omega}.$$

The parameter ω_0 can be chosen arbitrarily; we have taken $\omega_0 = +1.0$. The integrals $f(\omega)$ were evaluated up to $\omega = 40 m_K$ by numerical integration, and beyond this energy by using the asymptotic expressions of Udgaonkar ⁴⁾. Thus, the slope of a linear fit to the measured values of $D(\omega)$ determines the pole residue Γ , and the intercept at $x = x_0$ determines the value of $D(\omega_0)$.

Using available values for $D(\omega)$, we find

$$\left. \begin{aligned} \Gamma &= -0.12 \pm 0.32 \\ D(\omega_0) &= -0.9 \pm 0.2_{(m_K)-1} \end{aligned} \right\} P(\chi^2) = 0.70.$$

Although $D(\omega)$ for K^+ has been reasonably well determined at several energies, the data on $D(\omega)$ for K^- is still too meagre to permit a satisfactory test of the forward scattering dispersion relations for K mesons.

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