

## Exploring strangeness enhancement in proton-proton collisions at the LHC with event shapes

Suraj Prasad<sup>1,\*</sup>, Sushanta Tripathy<sup>2</sup>, Bhagyarathi Sahoo<sup>1</sup>, Neelkamal Mallick<sup>1</sup>, and Raghunath Sahoo<sup>1</sup>

<sup>1</sup>Department of Physics, Indian Institute of Technology Indore, Simrol, Indore 453552, India and

<sup>2</sup>CERN, 1211 Geneva, Switzerland

### Introduction

Strangeness enhancement is one of the important signatures to probe the formation of a thermalised and deconfined medium of partons in high-energy heavy-ion collisions, where proton-proton ( $pp$ ) collisions are taken as the baseline. Recently, an enhancement of strange hadron production with respect to pions is observed in high multiplicity  $pp$  collisions [1]. Several theoretical and phenomenological models have been proposed to explain the observed signatures of QGP in small systems. In fact, p-QCD based models such as PYTHIA8, with a multi-partonic interaction (MPI) based picture of colour reconnection (CR) and rope hadronisation (RH), can reproduce the experimental features of strangeness production [1]. Figure 1 shows the self-normalised yield ratios to pions as a function of average charged-particle multiplicity ( $\langle dN_{ch}/d\eta \rangle$ ) in  $pp$  collisions at  $\sqrt{s} = 13$  TeV using PYTHIA8 with CR and RH. Here, a large value of  $\langle dN_{ch}/d\eta \rangle$  corresponds to a large value of the number of multi-partonic interactions ( $N_{mpi}$ ). As one approaches high  $N_{mpi}$  events, the self-normalised yield ratios of strange hadrons to pions increase. The increment is higher for hadrons having larger valence strange quarks and is higher for baryons than for mesons with a similar number of valence strange quarks. However, MPI is a phenomenological mechanism that explains different experimental measurements and direct measurement of  $N_{mpi}$  is not viable in the experiments. Thus, we define event shapes, such

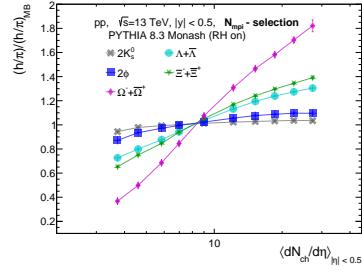


FIG. 1: Strange hadron production ratio to pions normalised to minimum bias events as a function of average charged particle multiplicity at midrapidity in  $pp$  collisions at  $\sqrt{s} = 13$  TeV using PYTHIA 8 [6].

as charge particle multiplicity in the midrapidity ( $N_{ch}^{\text{mid}}$ ) and forward rapidity ( $N_{ch}^{\text{fwd}}$ ), unweighted transverse spherocity ( $S_0^{\text{PT}=1}$ ) [2], transverse sphericity ( $S_T$ ) [3], relative transverse activity classifier ( $R_T$ ) [4] and charged particle flattening ( $\rho_{ch}$ ) [5], that are capable of selecting the soft-QCD-dominated isotropic events from the jetty events and possess large correlations with  $N_{mpi}$ . The details for event shape definition, event generation, and event and track selection cuts can be found in Ref. [6]. In this study, we show the feasibility of probing the strange hadron production with different event shapes in  $pp$  collisions at  $\sqrt{s} = 13$  TeV using PYTHIA8.

### Results

Figure 2 shows self-normalised yield ratios to pions as a function of  $\langle dN_{ch}/d\eta \rangle$  measured in different percentiles of  $N_{ch}^{\text{mid}}$  (upper left),  $N_{ch}^{\text{fwd}}$  (upper middle),  $S_0^{\text{PT}=1}$  (upper right),  $S_T$  (lower left),  $R_T$  (lower middle), and  $1 - \rho_{ch}$  (lower right) in  $pp$  collisions at

\*Electronic address: suraj.prasad@cern.ch

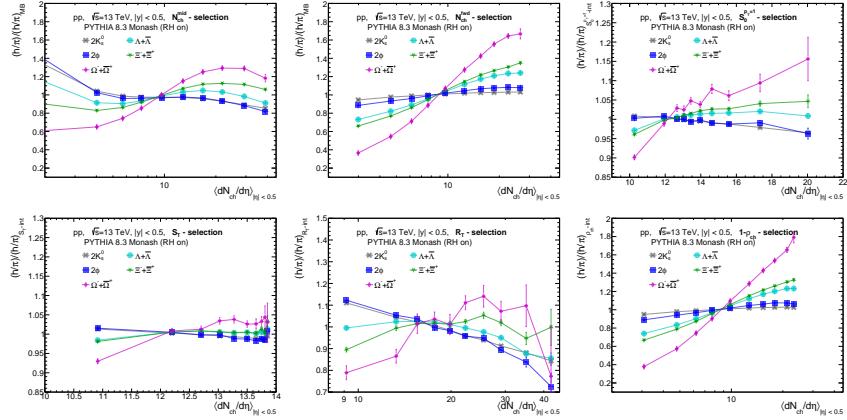


FIG. 2: Self-normalised yield ratios to pions as a function of  $\langle dN_{ch}/d\eta \rangle$  for different percentiles of  $N_{ch}^{mid}$  (upper left),  $N_{ch}^{fwd}$  (upper middle),  $S_0^{pT=1}$  (upper right),  $S_T$  (lower left),  $R_T$  (lower middle), and  $1 - \rho_{ch}$  (lower right) in  $pp$  collisions at  $\sqrt{s} = 13$  TeV using PYTHIA8 [6].

$\sqrt{s} = 13$  TeV using PYTHIA8. Similar to Fig. 1, in Fig. 2, we find a distinction in the self-normalised yield ratios to pions for the strange baryons having different numbers of valence strange quarks. Here, the increment of self-normalised yield ratios with  $\langle dN_{ch}/d\eta \rangle$  for all event shapes is higher for  $\Omega$  having three valence strange quarks as compared to  $\Lambda$ , having only one valence strange quark. Further, one finds a close resemblance between the self-normalised yield ratios to pions for the hidden strange meson,  $\phi$  and  $K_S^0$  possessing one valence strange quark. The self-normalised yield ratios of  $\phi$  and  $K_S^0$  mesons to pions are observed to decrease with an increase in  $\langle dN_{ch}/d\eta \rangle$  for the event shapes measured in the midrapidity which includes  $N_{ch}^{mid}$ ,  $S_0^{pT=1}$ ,  $S_T$  and  $R_T$ . This is an autocorrelation bias caused by defining the event shapes in a rapidity region where one measures the identified particles. In contrast, this selection bias is absent when event selection is performed using the event shapes measured in a different rapidity region, which includes  $N_{ch}^{fwd}$  and  $1 - \rho_{ch}$ . Here, the self-normalised yield ratios to pions for the strange hadrons are all distinct for events selected with  $N_{ch}^{fwd}$  and  $1 - \rho_{ch}$ , representing their sensitivity to probe the strangeness production. However,

one finds  $1 - \rho_{ch}$  very closely mimics the event selection with  $N_{mpi}$  shown in Fig. 1, making  $1 - \rho_{ch}$  an ideal choice among the currently available event shapes to study strangeness production at the LHC.

## Summary

In summary, we show the self-normalised yield ratios of different strange hadrons to pions with event selection based on different event shapes in  $pp$  collisions at  $\sqrt{s} = 13$  TeV using PYTHIA 8. We find that charged particle flattening is one of the best among the currently available event shapes for studying strangeness production at the LHC energies.

## References

- [1] S. Acharya *et al.* [ALICE], Eur. Phys. J. C **80**, 693 (2020).
- [2] S. Acharya *et al.* [ALICE], JHEP **05**, 184 (2024).
- [3] B. Abelev *et al.* [ALICE], Eur. Phys. J. C **72**, 2124 (2012).
- [4] T. Martin, P. Skands and S. Farrington, Eur. Phys. J. C **76**, 299 (2016).
- [5] A. Ortiz *et al.* Phys. Rev. D **107**, 076012 (2023).
- [6] S. Prasad, B. Sahoo, S. Tripathy, N. Mallick and R. Sahoo, [arXiv:2409.05454 [hep-ph]].