

# DECOUPLING OF NITROGEN AND OXYGEN IMPURITIES IN NITROGEN DOPED SRF CAVITIES\*

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## Abstract

The performance of superconducting radiofrequency (SRF) cavities is critical for enabling the next generation of efficient high-energy particle accelerators. Recent developments have focused on altering the surface impurity profile through *in-situ* baking, furnace baking, and doping to introduce and diffuse beneficial impurities such as nitrogen, oxygen, and carbon. However, the precise role and properties of each impurity are not well understood. In this work, we attempt to disentangle the role of nitrogen and oxygen impurities through time-of-flight secondary ion mass spectrometry using niobium cavity cutouts baked at temperatures varying from 120-800 °C with and without nitrogen injection. Understanding the underlying mechanisms behind these impurities will drive further optimization in the tailoring of nitrogen and oxygen profiles for high-performance SRF cavities.

## INTRODUCTION

Improving the performance of superconducting radiofrequency (SRF) cavities is the key to enabling the next generation of particle accelerators such as the International Linear Collider (ILC) or the 8 GeV Booster Replacement Linac at Fermilab [1, 2]. Two key performance metrics are quality factor ( $Q_0$ ) and accelerating gradient ( $E_{acc}$ ).  $Q_0$  is inversely proportional to power loss in a cavity, so higher  $Q_0$  cavities will lower operating costs and helium expenses [3]. High  $E_{acc}$  in cavities increases the magnitude of the kick provided to a particle, making an accelerator more efficient per unit length and lowering the cost of building a future accelerator. The performance of the SRF cavity depends primarily on the composition of the rf layer, which includes the first 100 nm of the surface, through which superconducting currents can penetrate [4]. One of the main culprits which limits SRF cavity performance has been hydrogen; the presence of hydrogen impurities in the rf layer leads to the precipitation of highly lossy niobium hydrides and nanohydrides on the surface [5–7]. Many studies have been conducted on the development of surface treatments to minimize these prohibitive effects through the introduction and diffusion of impurities such as oxygen and nitrogen [8–11].

One such treatment is nitrogen doping: a state-of-the-art high  $Q_0$  treatment that enabled the construction of LCLS-II [12, 13]. Nitrogen doping introduces uniform and dilute

concentrations of nitrogen into the rf layer [8]. A key performance feature of a N doped cavity is the anti Q-slope, an increase in  $Q_0$  with increasing  $E_{acc}$ , which typically occurs at  $E_{acc}$  below 20 MV/m [8]. While there are some theories which propose the cause of anti-Q slope, there is no general consensus on which underlying mechanisms are responsible for it [14–17]. One theory is that nitrogen impurities act as a trap for interstitial hydrogen, preventing the precipitation of niobium hydrides [18].

The presence of interstitial oxygen in the Nb lattice has also been shown to reduce the precipitation of lossy niobium hydrides [18]. Recent work from E. Lechner and D. Bafia have highlighted the role of diffused oxygen in improved SRF cavity performance [19–21]. A new surface treatment termed oxygen doping (200°C *in-situ* baking) displays high Q performance by yielding an approximately uniform concentration of oxygen in the surface without completely removing the oxide [19]. The performance of the O doped cavities exhibit the anti-Q slope, a characteristic feature of N doped cavities [23, 25]. Given the similarities in performance, it has been proposed by D. Ford and P. Koufalis, among others, that impurities in the surface layer of N and O doped cavities perform similar functions in mitigating sources of loss [18, 24]. D. Ford has shown through first principle calculations that Nb will preferentially bond to O and N over H, with binding energies of -7.02 eV, -7.39 eV, and -2.41 eV, respectively [18]. In this work, we correlate the concentration of impurities, specifically N and O, to SRF cavity performance.

## EXPERIMENTAL METHOD

Nb cavity cutouts from cavity TE1AES008 and an array of single-cell TESLA shaped 1.3 GHz Nb cavities were treated with the following treatments: Electropolishing (EP), 2/0+5 N-doped, Low Temperature Baking (LTB) at 120°C, and O doping at 200°C. Each cutout and cavity was subjected to the same baseline treatment of 800°C degas and 40 um EP. Between treatments, cutouts would receive between 10 um and 30 um EP to reset the surface. The specifics of each treatment are presented in Table 1.

The cutouts were analyzed with time of flight secondary ion mas spectrometry (SIMS) to obtain depth profiles of the impurities present in the Nb lattice. Each reported impurity depth profile is taken as the average of three 200  $\mu$ m x 200  $\mu$ m locations on each sample. The sputter crater for each spot is 600  $\mu$ m x 600  $\mu$ m to avoid edge effects, and the sputtering energy for the Cs ion beam is 2 keV. Vacuum was maintained at  $< 4 \times 10^{-10}$  mbar in the analysis chamber and  $10^{-7}$  mbar in the heating chamber. Cutouts

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Table 1: Treatment process history for cavity cutouts and 1.3 GHz single-cell SRF cavities.

Cavity	Steps	
TE1AES010	EP baseline	
TE1AES010	120°C × 3 hrs <i>in-situ</i>	+120°C × 3 hrs (total 6 hrs) <i>in-situ</i>
TE1PAV009	120°C × 48 hrs <i>in-situ</i>	
TE1AES017	200°C × 1 hr <i>in-situ</i>	+200°C × 10 hrs (total 11 hrs) <i>in-situ</i>
TE1AES021	200°C × 20 hr <i>in-situ</i>	
TE1AES024	800°C × 3 hrs UHV	2 min N inj at 25 mTorr
		5 μm cold EP

were analyzed alongside implanted N and O standards to calibrate relative SIMS intensity to concentration [26]. O<sup>-</sup> and NbN<sup>-</sup> signals were used to determine the relative sensitivity factor (RSF) for O and N respectively. Experimentally determined RSF values relative to the Nb were found to be  $3.66 \times 10^{21}$  ions/cm<sup>3</sup> for N and  $2.11 \times 10^{20}$  ions/cm<sup>3</sup> for O.

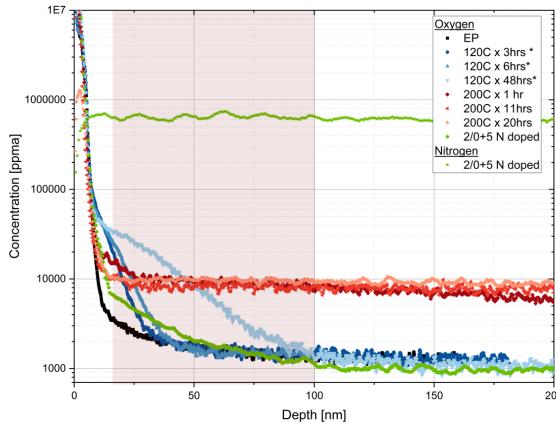


Figure 1: SIMS depth profile for absolute concentrations of N and O in Nb cavity cutouts. Concentration for starred (\*) measurements were not measured directly and are scaled to the background EP baseline measurement.

Figure 1 contains the absolute concentration for each of the treatments on cavity cutouts. When available, RSF was acquired from standards during the same measurement session. If not, concentrations were scaled to background signal of the EP baseline. Baking at 120°C gradually dissolves oxygen into the bulk with a diffusion length of 40—100 nm [19]. *In-situ* baking at 200°C introduces a uniform concentration of O in the rf layer. No N was measured in any of the EP 120°C, or 200°C baked cavity cutouts. Nitrogen doping yields a uniform concentration of N in the rf layer, and its O concentration is comparable to that of the baseline EP O profile. This suggests that the presence of O in N doping is negligible, and we can effectively decouple N and O.

Cavity cutout results are then compared to SRF cavity performance of single-cell TESLA shaped Nb cavities of resonant frequency 1.3 GHz subjected to the same treatments in Table 1. For *in-situ* treatments, the cavities were fully assembled after EP and prior to baking. Vacuum was maintained after during testing and baking to prevent the surface oxide from regrowing. Cavities were tested at the

Fermilab Vertical Test Stand (VTS) to find  $Q_0$  vs.  $E_{acc}$  at 2 K and <1.5 K in continuous wave (CW) operation. Cooling followed the fast cool down protocol to minimize the possibility of trapping magnetic flux [19].

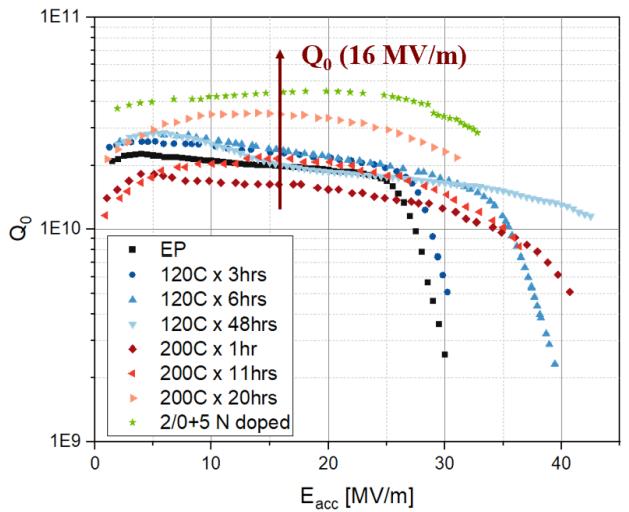


Figure 2:  $Q_0$  vs  $E_{acc}$  performance for single-cell 1.3 GHz SRF cavities of various treatment recipes. Colors transition from darker to lighter with increased O diffusion.

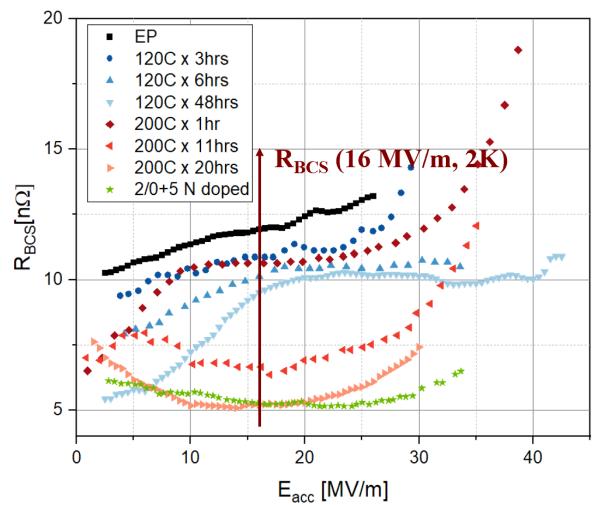


Figure 3:  $R_{BCS}$  vs  $E_{acc}$  performance for single-cell 1.3 GHz SRF cavities of various treatment recipes. Colors transition from darker to lighter with increased O diffusion.

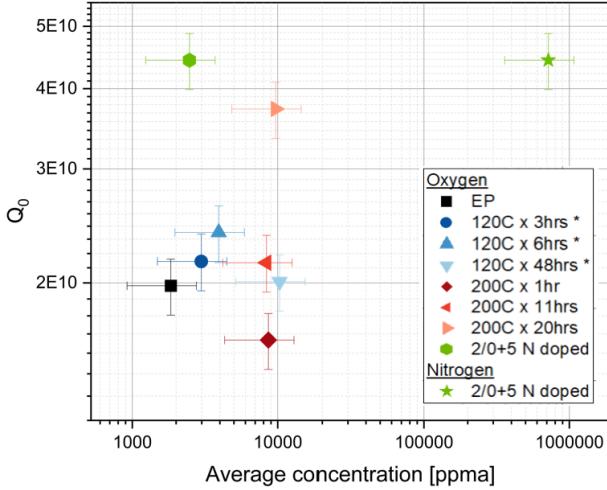


Figure 4:  $Q_0(16 \text{ MV/m})$  vs average absolute concentration of N and O.

Figures 2 and 3 shows the  $Q_0$  and  $R_{BCS}$  performance of single-cell 1.3 GHz SRF cavities for each treatment recipe. We can see the gradual evolution towards high  $E_{acc}$  with increased 120°C baking as well as towards high  $Q_0$  with increased 200°C baking. A decrease in  $R_{BCS}$  from longer baking and more diffused O drives the observed increase in  $Q_0$ . The performance for the O doping treatment of 200°C × 20 hrs is comparable to the high  $Q_0$  and low  $R_{BCS}$  of 2/0+5 N doping.

## RESULTS AND DISCUSSION

To correlate the results from sample cutouts and cavities, we extracted  $Q_0$  and  $R_{BCS}$  at 16 MV/m for each of the treatments. 16 MV/m was chosen as it avoids both high field Q-slope degradation at >25 MV/m and low field (<5 MV/m) effects from surface oxides. From Fig. 1, we averaged the concentration over the first 100 nm of the surface, determined by the penetration depth of fields in Nb. The first 15 nm were excluded to consider only the concentration of dissolved O and N within the bulk.

Figure 4 correlates  $Q_0(16 \text{ MV/m})$  with the average concentration of dissolved O and N in the surface. We do not observe any distinguishable pattern since  $Q_0$  is sensitive to both changes in the native oxide and dissolved impurities in the bulk. The concentration of N in N doped is about 2 orders of magnitudes higher than the concentrations of O from baking; the ratio of N:Nb is 0.7:1, which is much higher than expected. Further investigation is necessary to confirm this elevated concentration of N.

In Fig. 5, we correlate  $R_{BCS}$ , the quasiparticle losses from impurities within the bulk, with the average concentration of impurities in the rf layer. From a baseline  $R_{BCS}$  of 12 nΩ for EP,  $R_{BCS}$  decreases as more O is diffused into the Nb bulk. We can also confirm that the concentration of O in N doped is too low to account for the large improvement in  $R_{BCS}$  for the N doped treatment. Assuming the concentration of N is accurate, 200°C × 20 hours achieves similarly low  $R_{BCS}$

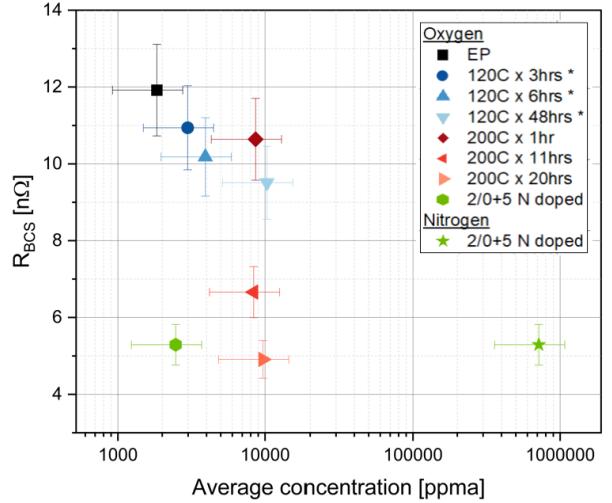


Figure 5:  $R_{BCS}(16 \text{ MV/m}, 2 \text{ K})$  vs average absolute concentration of N and O.

as N doped with significantly less O. Microscopically, this would suggest that O is much more effective at trapping H than N is. However, this disagrees with first principle calculations that demonstrate similar binding energies for NbO and NbN [18].

## CONCLUSION

We performed material and cavity performance measurements on a number of SRF treatment recipes. We were able to decouple the effect of N and O impurities as there is no N for *in-situ* baking and negligible O for N doping. We compared the results from cavity cutouts and cavities to that as concentration of dissolved O increased,  $R_{BCS}$  decreased. With sufficient O diffusion, we can achieve a similarly low  $R_{BCS}$  as with N in N doped. In addition, it seems like significantly less O is necessary to achieve the same improvement in performance as N, suggesting that O is more effective than N at trapping H. O diffusion may be a promising technology towards the development of a next generation high  $Q_0$  and high  $E_{acc}$  SRF cavity treatment. More studies are necessary to confirm these results.

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