

## Transfer reaction for the study of $\alpha$ -cluster states in $^{54}\text{Cr}$

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## Introduction

Emergence of  $\alpha$ -clustering in nuclei and its impact on their various properties, particularly starting from light nuclei within the nuclear table [1] remains a topic of contemporary interest. The concept of clustering has been observed since the early days of nuclear physics, where an  $\alpha$ -particle can serve as the fundamental building block for certain nuclei. The  $\alpha$ -cluster model serves as a valuable tool to elucidate significant aspects of nuclear behavior, including electromagnetic transition strength,  $\alpha$ -decay width, and  $\alpha$ -particle scattering data. Moreover, it simplifies the computationally challenging task by reducing the degrees of freedom. In lighter nuclei with mass numbers less than  $A = 40$ ,  $\alpha$ -clustering is a prominent characteristic in numerous states [2–4]. In the  $fp$ -shell region,  $^{44}\text{Ti}$  has been extensively studied with the interpretation of  $\alpha + \text{core}$  ( $^{40}\text{Ca}$ ) configurations. Recently, M. A. Souza and H. Miyake proposed that  $^{46}\text{Cr}$  and  $^{54}\text{Cr}$  are the most favorable even-even Cr isotopes for the  $\alpha + \text{core}$  configuration [5]. Hence the cluster structure study of  $^{54}\text{Cr}$  is important with available stable projectile-target combinations.

In this present work, we have populated the excited states in  $^{54}\text{Cr}$  via the  $\alpha$ -transfer to  $^{50}\text{Ti}$  nucleus. The reactions were studied through particle- $\gamma$  coincidence measure-

ments. Different excited states of  $^{54}\text{Cr}$  have been identified with triton gate of  $E_t > 17.5$  MeV. Along with the yrast states, several low-lying non-yrast states were found through this reaction. The  $\alpha + \text{core}$  spectroscopic factors for the observed excited states have been obtained.

## Experimental details

The excited states of  $^{54}\text{Cr}$  were populated via an  $\alpha$ -transfer reaction, specifically  $^{50}\text{Ti}(^7\text{Li}, t)^{54}\text{Cr}$ , using 20 MeV beam provided by the 14UD Pelletron Linac accelerator Facility (PLF) at TIFR, Mumbai, India. For this experiment, a self-supporting target of  $^{50}\text{Ti}$  with a thickness of approximately 1.48 mg/cm<sup>2</sup> was employed. The de-exciting  $\gamma$ -rays of residual nuclei were detected by the Indian National Gamma Array (INGA) at TIFR, consisting of 17 Compton-suppressed clover HPGe detectors at different angles with respect to the beamline. To detect outgoing charged particles, we utilized an array of thirteen CsI(Tl) detectors. Additionally, a Si surface barrier detector (monitor) was incorporated to detect scattered beam particles for cross-section calculations. The signals from the individual CsI(Tl) detectors and each crystal of the clover detectors were collected using a total of six 12-bit 100 MHz PIXIE-16 modules. The CsI(Tl) detectors were calibrated using the triton energy spectrum from the reaction  $^{12}\text{C}(^7\text{Li}, t)^{16}\text{O}$  at a beam energy of 20 MeV. For the calibration of the clover HPGe detectors, we employed a mixed  $^{133}\text{Ba}$

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$^{152}\text{Eu}$  source.

## Results and discussion

The collected experimental data have been sorted by a MultipARameter time-stamped based COincidence Search (MARCOS) code, developed at TIFR. The particle- $\gamma$  matrix files and ROOT NTuple with two- and higher-fold coincidence events have been generated. From each raw pulse of the CsI(Tl), the QDC-short and QDC-long have been calculated. The 2D particle identification spectrum has been generated in coincidence with any of the clover detectors. A two-dimensional spectrum of QDC-short vs QDC-long for  $34^\circ$  CsI(Tl) detectors has been shown in Fig. 1. Proton-, triton-, and alpha-band have been identified in the spectrum. The  $\gamma$  spectrum of the clover detectors has been obtained and cleaned by applying the 2d gate to the particle spectrum. In the present heavy-ion collisions, various re-

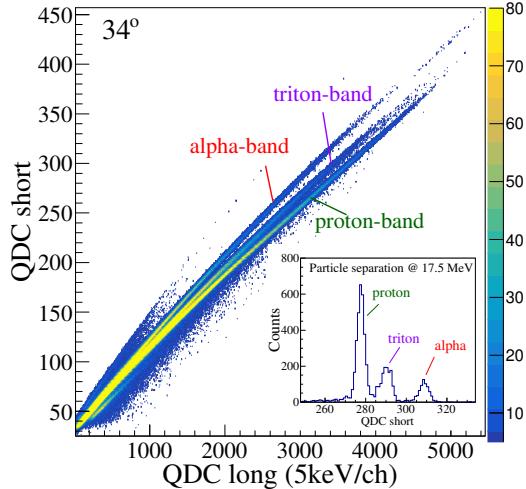


FIG. 1: The calibrated particle spectrum from a CsI(Tl) detector positioned at  $\theta_{\text{lab}} = 34^\circ$  reveals distinct bands of light-charged particles, including protons, tritons, and  $\alpha$ -particles, as observed in the current measurement. Insets demonstrate a sharp differentiation between these particle types at an energy of 17.5 MeV.

action channels become energetically favored. The complete fusion evaporation reaction has a dominant cross-section. Specifically, nu-

clei such as  $^{54,55}\text{Mn}$  and  $^{53-55}\text{Cr}$  are substantially populated through different pathways, including  $(3n)$ ,  $(2n)$ ,  $(p3n)$ ,  $(p2n)$ , and  $(pn)$  complete fusion evaporation channels. Moreover, the nucleus  $^{54}\text{Cr}$  can also be generated through the  $\alpha$ -cluster transfer to  $^{50}\text{Ti}$ . In the incomplete fusion channel, the  $^{54}\text{Cr}$  compound nucleus will be formed which subsequently undergoes de-excitation processes, leading to the production of  $^{51-53}\text{Cr}$  through multi-neutron ( $xn$ )-evaporation and  $^{51-53}\text{V}$  via  $(pxn)$  channels. In this work, along with yrast state (up to  $8^+$ ), multiple low-lying non-yrast  $2^+$  and  $4^+$  states of  $^{54}\text{Cr}$  have been populated which are typically challenging to access through fusion evaporation reactions. The integrated transfer cross-sections for both the observed yrast and non-yrast states have been measured. We have then compared these measurements with calculations performed using Coupled Reaction Channels in the FRESCO framework, allowing us to extract  $\alpha$ -core spectroscopic factors [6]. Our findings suggest moderate  $\alpha$ -cluster structure for the yrast states of  $^{54}\text{Cr}$  and the clustering effect reduces for the non-yrast states.

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