

AXIONS AND AXION-LIKE PARTICLES

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The physics case for axions and axion-like particles is reviewed and an overview of ongoing and near-future laboratory searches is presented.

1 Introduction

The Standard Model (SM) of particle physics describes the properties of known matter and forces to a fantastic accuracy. However, it can not be considered to be a complete and fundamental theory. Most notably, it does not give a satisfactory explanation for the values of its many parameters, it does not provide a predictive quantum theory of gravity, and it does not explain the origin of the dark sector of the Universe. In fact, the observation that nearly thirty percent of the energy content of the universe consists of dark matter provides a very strong case for the existence of particles beyond the SM.

Intriguingly, dark matter candidates occur automatically in a number of SM extensions which were originally motivated by completely other reasons. Prominent examples are the lightest supersymmetric (SUSY) partners, e.g. neutralinos (partners of the SM gauge or Higgs bosons) or gravitinos (partners of the graviton), in R-parity conserving SUSY extensions of the SM – the latter being motivated by grand unification and the solution of the hierarchy problem. Another widely-discussed example are the axion – occurring in extensions of the SM solving the strong CP problem – and axion-like particles (ALPs), which are often predicted by embeddings of the SM in string theory.

In this contribution we review the theoretical motivation for the axion and ALPs (Sec. 2), summarize the various hints for their existence from astrophysics and cosmology (Sec. 3), and report on ongoing and near-future laboratory experiments and observatories (Sec. 4).

2 Theoretically Favored ALP Candidates

Many extensions of the SM feature one or several spontaneously broken global $U(1)_i$ symmetries, $i = 1, \dots, n_{\text{ax}}$. At energies below their symmetry breaking scales v_i , Nambu-Goldstone boson fields a'_i arise as massless excitations of the angular part of the SM singlet complex scalar fields ϕ_i whose vacuum expectation values (vev) $\langle \phi_i \rangle = v_i/\sqrt{2}$ break the $U(1)_i$ symmetries: $\phi_i(x) = (v_i + \sigma_i(x))e^{ia'_i(x)/v_i}/\sqrt{2}$. Their interactions with SM particles, e.g. with gluons (described by the gluonic field strength $G_{\mu\nu}$), photons (described by the electromagnetic field strength $F_{\mu\nu}$), and electrons (described by the spinor e), are suppressed by inverse powers of the supposedly large

symmetry breaking scales, $f_{a'_i} = v_i \gg v$, where $v = 246$ GeV is the electroweak Higgs vev,

$$\begin{aligned} \mathcal{L} = & \frac{1}{2} \partial_\mu a'_i \partial^\mu a'_i - \frac{\alpha_s}{8\pi} \left(\sum_{i=1}^{n_{\text{ax}}} C'_{ig} \frac{a'_i}{f_{a'_i}} \right) G_{\mu\nu}^b \tilde{G}^{b,\mu\nu} \\ & - \frac{\alpha}{8\pi} \left(\sum_{i=1}^{n_{\text{ax}}} C'_{i\gamma} \frac{a'_i}{f_{a'_i}} \right) F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{2} \left(\sum_{i=1}^{n_{\text{ax}}} C'_{ie} \frac{\partial_\mu a'_i}{f_{a'_i}} \right) \bar{e} \gamma^\mu \gamma_5 e + \dots \end{aligned} \quad (1)$$

Here, the couplings to the gluons, C'_{ig} , and to the photons, $C'_{i\gamma}$, arise from integrating out fermions with chirally anomalous $U(1)_i$ charge assignments. Particularly well-motivated examples for such Nambu-Goldstone bosons are:

- The *axion* A – the particle excitation of the superposition of all Nambu-Goldstone fields a'_i coupling to the topological charge density in QCD, $q \equiv \frac{\alpha_s}{8\pi} G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$, in Eq. (1),

$$\frac{A}{f_A} \equiv \sum_{i=1}^{n_{\text{ax}}} C'_{ig} \frac{a'_i}{f_{a'_i}}. \quad (2)$$

This field replaces the theta parameter in QCD by a dynamical quantity, $\theta_A(x) = A(x)/f_A$, spontaneously relaxing to zero, $\langle \theta_A \rangle = 0$ – thereby explaining the non-observation of strong CP violation^{1,2,3}. In fact, topological non-trivial gluonic fluctuations result in an effective potential for θ_A ,

$$V(\theta_A) = \frac{\chi(0)}{2} \theta_A^2 + \mathcal{O}(\theta_A^4) \simeq \frac{m_\pi^2 f_\pi^2}{2} \frac{m_u m_d}{(m_u + m_d)^2} \theta_A^2 + \mathcal{O}(\theta_A^4), \quad (3)$$

which has a localized minimum at $\theta_A = 0$. Here, $\chi(0) = \langle Q^2 \rangle |_{\theta=0} / \int d^4x$, with $Q = \int d^4x q(x)$, is the topological susceptibility, m_π and f_π are the mass and the decay constant of the pion, and m_u and m_d are the quark masses, respectively. Moreover, the topological fluctuations give the axion a small mass, which can be read off from the quadratic part in Eq. (3),

$$m_A = \frac{m_\pi f_\pi}{f_A} \frac{\sqrt{m_u m_d}}{m_u + m_d} \simeq 0.6 \text{ meV} \times \left(\frac{10^{10} \text{ GeV}}{f_A} \right), \quad (4)$$

rendering the axion a pseudo-Nambu-Goldstone boson. Due to the mixing with the neutral pion, the axion has a universal coupling to photons,

$$\mathcal{L} \supset -\frac{g_{A\gamma}}{4} A F_{\mu\nu} \tilde{F}^{\mu\nu}; \quad |g_{A\gamma}| \sim \frac{\alpha}{2\pi f_A} \sim 10^{-12} \text{ GeV}^{-1} \left(\frac{10^9 \text{ GeV}}{f_A} \right), \quad (5)$$

which allows for various experimental probes.

- The *majoron* – the Nambu-Goldstone boson arising from the breaking of global lepton number symmetry^{4,5} at a high energy scale $f_L = v_L$ – thereby explaining the smallness of the masses of the left-handed SM active neutrinos by a see saw relation involving the electroweak scale Dirac-type mass $M_D = F v$ and the large Majorana-type mass $M_M = y f_L$ of extra right-handed SM singlet neutrinos,

$$m_\nu = -M_D M_M^{-1} M_D^T = -F y^{-1} F^T \frac{v^2}{f_L} = 0.6 \text{ eV} \left(\frac{10^{12} \text{ GeV}}{f_L} \right) \left(\frac{-F y^{-1} F^T}{10^{-2}} \right). \quad (6)$$

- *Familons* arising from the breaking of global family symmetries^{6,7,8}.
- *Closed string axions*^{9,10,11} – Kaluza-Klein zero modes of antisymmetric tensor fields – the latter belonging to the massless spectrum of the bosonic string propagating in ten dimensions. Their number n_{ax} is determined by the topology of the compactified manifold.

- *Accidental axion and ALPs* arising from the breaking of accidental global $U(1)$ symmetries that appear as low energy remnants of exact discrete symmetries – the latter being postulated in purely field theoretic set ups^{12,13} or occurring automatically in orbifold compactifications of the heterotic string¹⁴.

Therefore, searches for the axion A – the linear combination coupling to the topological charge density in QCD, Eq. (2) – and further ALPs a_j – the $n_{\text{ax}-1}$ Nambu-Goldstone bosons perpendicular to the axion in field space – are theoretically very well motivated. Moreover, as we will review next, their existence is also suggested on cosmological and astrophysical grounds.

3 Physics Case for Axions and ALPs

3.1 Cold Dark Matter

Axions and/or ALPs – if they are pseudo Nambu-Goldstone bosons, i.e. if they have a (small) mass due to non-perturbative effects or explicit symmetry breaking – are excellent cold dark matter candidates. In fact, for large symmetry breaking scales, axions and ALPs are very long lived. They are produced in the early universe via the vacuum realignment mechanism as a coherent state of many, extremely non-relativistic particles in the form of a classical, spatially coherent oscillating field^{5,16,17}. Neglecting anharmonic effects, today’s (time t_0) fraction of axion or ALP dark matter produced via the vacuum realignment mechanism is proportional to the spatially averaged field amplitude squared, $\langle a^2 \rangle \equiv f_a^2 \langle \theta_a^2 \rangle$, at the time when the oscillations started, $t_{\text{osc}} \simeq (3/2) m_a^{-1}(t_{\text{osc}})$,¹⁸

$$R_a = \frac{\rho_a}{\rho_{\text{DM}}}(t_0) \simeq 0.2 \sqrt{\frac{m_a(t_0)}{\text{eV}}} \sqrt{\frac{m_a(t_0)}{m_a(t_{\text{osc}})}} \left(\frac{f_a}{10^{11} \text{ GeV}} \right)^2 \langle \theta_a^2 \rangle. \quad (7)$$

Here, the indicated time-dependence of the mass arises from its temperature dependence, $m_a(t) \equiv m_a(T(t))$, taking into account possible plasma effects.

The values of $\langle \theta_a^2 \rangle$ depend crucially on whether the global symmetry breaking occurred before inflation ends and there was no symmetry restoration after inflation (“pre-inflationary SSB”), or the opposite (“post-inflationary SSB”),

$$\langle \theta_a^2 \rangle = \begin{cases} \theta_i^2 + \left(\frac{H_I}{2\pi f_a} \right)^2, & \text{if } f_a > \max\left(\frac{H_I}{2\pi}, \epsilon_{\text{eff}} E_I \right), \\ \frac{\pi^2}{3}, & \text{if } f_a < \max\left(\frac{H_I}{2\pi}, \epsilon_{\text{eff}} E_I \right). \end{cases} \quad (8)$$

Here, H_I is the Hubble expansion rate during inflation, $E_I = \sqrt{\frac{3}{8\pi}} M_{\text{Pl}} H_I$ is the energy scale of inflation, and $\epsilon_{\text{eff}} \in (0, 1)$ is a reheating efficiency parameter.

In the first, “pre-inflationary SSB” case, spatial variations in $\theta_a \equiv a/f_a$ are smoothed out over our Hubble volume. Then $\theta_i \equiv \langle \theta_a \rangle$ is an essentially environmental parameter drawn from a uniform distribution, $\theta_i \in [0, \pi]$, with a small variance, $\sigma_{\theta_a} = H_I/(2\pi f_a)$, arising from quantum fluctuations during inflation. These fluctuations provide a lower bound on R_a . The non-observation of the associated isocurvature fluctuations in the anisotropies of the temperature of the cosmic microwave background (CMB) severely constrain axions and ALPs, disfavoring a sizeable region in f_a - H_I parameter space^{19,20,21,22,23}, as is exemplified in Fig. 1. In fact, an H_I of order 10^{14} GeV, as indicated by the recent detection of B-mode polarization in the CMB by BICEP2²⁴ – if the latter is interpreted as originating from primordial gravitational waves rather than from foreground dust^{25,26} – would strongly disfavor pre-inflationary SSB scenarios^{27,28} (possible ways out of this conclusion have been put forward in Refs.^{29,30,31,32}).

This would favor the second, “post-inflationary SSB” case, in which the initial misalignment angle takes on different values in small (linear dimension $\sim m_a^{-1}(t_{\text{osc}})$) patches in the universe,

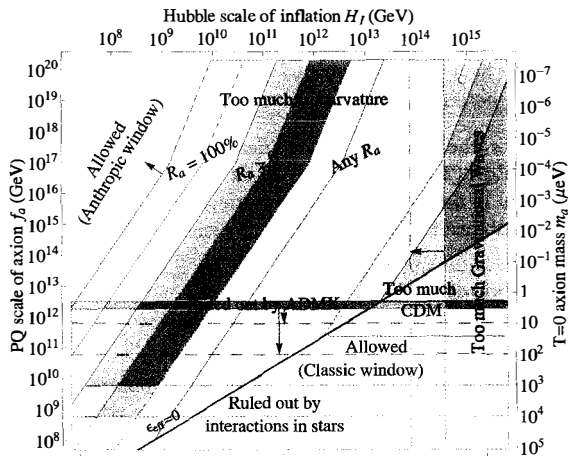


Figure 1 – Disfavored (colored) and favored (white) regions of axion – inflation parameter space²².

resulting in a vanishing mean value, $\langle \theta_a \rangle = 0$, and a variance $\sigma_{\theta_a} = \pi/\sqrt{3}$, due to small-scale variations. These variations again provide a lower bound on R_a and, correspondingly, an upper bound on f_a . For the axion, this bound is of order^{23,33,34} $f_A \lesssim \text{few} \times 10^{11}$ GeV, see Fig. 1. Moreover, an additional cold dark matter population arises in the post-inflationary SSB scenario from the decay of topological defects, which radiates non-relativistic axions or ALPs. This additional contribution narrows down the parameter space of axions or ALPs even more, however with substantial uncertainties^{23,35,36,37}.

In R-parity conserving supersymmetric models, more possibilities arise: cold dark matter might be a mixture of axions along with the lightest SUSY particle (LSP)^{38,39}. Candidates for the LSP include then the lightest neutralino, the gravitino, the axino, or a sneutrino. In the case of a neutralino LSP, saxion and axino production in the early universe have a strong impact on the neutralino and axion abundance. For large values of f_A , saxions from the vacuum re-alignment mechanism may produce large relic dilution via entropy dumping.

Recently, two groups reported the observation of an unidentified 3.55 keV line from galaxy clusters and from the Andromeda galaxy^{40,41}. It is tempting to identify this line with the expected signal from two photon decay of 7.1 keV mass ALP dark matter^{42,43}. To match the observed X-ray flux, but allowing for the likely possibility, that the ALP dark matter makes only a fraction $R_a \equiv \rho_a/\rho_{\text{DM}}$ of the total density of dark matter, the required lifetime and thus decay constant of the ALP is (cf. Fig. 2)

$$g_{a\gamma} \sim R_a^{-1/2} 10^{-(17 \div 18)} \text{ GeV}^{-1}; \quad f_a \sim C_{a\gamma} R_a^{1/2} 10^{14 \div 15} \text{ GeV}. \quad (9)$$

However, some observations such as the anomalous line strength in the Perseus cluster and the enhanced strength of the line emission in the cool cores of the Perseus, Ophiuchus and Centaurus clusters seem to be better fitted in models in which the unidentified line arises from the decay of a 7.1 keV (scalar or fermionic) dark matter species into a very light ($m_a \lesssim 10^{-10}$ eV) ALP, that subsequently converts to photons in astrophysical magnetic fields^{48,49}.

3.2 Gamma Transparency of the Universe

Gamma-ray spectra from distant active galactic nuclei (AGN) should show an energy and redshift-dependent exponential attenuation, $\exp(-\tau(E, z))$, due to e^+e^- pair production off the extra-

galactic background light (EBL) – the stellar and dust-reprocessed light accumulated during the cosmological evolution following the era of re-ionization. In fact, Fermi⁵⁰ and H.E.S.S.⁵¹ have put sensible constraints on the EBL using their recent first detection of this effect. However, there is growing evidence^{52,53,54,55} for an anomalous transparency of the universe for gamma-rays at large optical depth, $\tau \gtrsim 2$. This may be explained by photon \leftrightarrow ALP oscillations: the conversion of gamma rays into ALPs in the magnetic fields around AGNs or in the intergalactic medium, followed by their unimpeded travel towards our galaxy and the consequent reconversion into photons in the galactic/intergalactic magnetic fields^{53,56,57,58,59}. This explanation requires a very light ALP, which couples to two photons with strength⁵⁸ (cf. Fig. 2),

$$g_{a\gamma} \gtrsim 10^{-12} \text{ GeV}^{-1}; \quad m_a \lesssim 10^{-7} \text{ eV}. \quad (10)$$

3.3 Cosmic ALP Background Radiation

There are observational hints on extra dark radiation in the primordial plasma during big bang nucleosynthesis and before photon decoupling – beyond the one from the three active neutrino species – at the one to three sigma level^{60,61}. Intriguingly, a cosmic ALP background (CAB) radiation corresponding to an effective number $\Delta N_{\text{eff}} \sim 0.5$ of extra neutrinos species can be naturally produced by the decay of a heavy ($\sim 10^6 \text{ GeV}$) modulus with Planck mass suppressed couplings^{62,63}. In fact, an observed soft X-ray excess in the Coma cluster may be explained by the conversion of such a CAB into photons in the cluster magnetic field^{64,65}. This explanation requires that the CAB spectrum is peaked in the soft X-ray region and that the ALP coupling and mass satisfy

$$g_{a\gamma} \gtrsim 10^{-13} \text{ GeV}^{-1} \sqrt{0.5/\Delta N_{\text{eff}}}; \quad m_a \lesssim 10^{-12} \text{ eV}, \quad (11)$$

respectively, overlapping with the parameter range (10) preferred by the ALP solution of the gamma-ray transparency puzzle, as is apparent in Fig. 2.

4 Axion and ALP Experiments and Observatories

We have seen in the last section that there are strong theoretical, cosmological and astrophysical hints suggesting the existence of the axion plus additional two to three ALPs (cf. red regions in Fig. 2). Fortunately, a sizeable part of the favored regions in axion and ALP parameter space can be explored in the foreseeable future with experimental searches based on axion or ALP photon oscillations in magnetic fields (cf. green regions in Fig. 2), as we review in this section.

4.1 Haloscope Searches

Haloscopes directly search for galactic halo dark matter axions and ALPs in the laboratory. Currently, the most sensitive ones exploit electromagnetic cavities placed in a strong magnet⁶⁶. They aim for the detection of electromagnetic power arising from the conversion of dark matter axions or ALPs into real photons, with frequency $\nu = m_a/(2\pi) = 0.24 \text{ GHz} \times (m_a/\mu\text{eV})$. The best sensitivity is reached on resonance, the power output then being proportional to the quality factor of the cavity. The Axion Dark Matter eXperiment (ADMX) has indeed reached recently the sensitivity to probe axion dark matter⁶⁷ (see Fig. 2). The ongoing experiments ADMX-II and ADMX-High Frequency (HF) aim to explore the green regions labelled “Haloscopes” in the same figure. Further haloscope opportunities in complementary mass ranges may arise from recycling available microwave cavities and magnets at accelerator laboratories^{68,69}.

Other new haloscope concepts are also being investigated. A microwave Fabry-Perot resonator in a spatially varying magnetic field may be exploited⁷⁰ to search for axion/ALP dark matter with masses above $40 \mu\text{eV}$. Converted photons from axion/ALP dark matter could be focused in a manner similar to a dish antenna, allowing for broad-band searches⁷¹. Precision

a dedicated superconducting toroidal magnet with much bigger aperture than CAST, a detection system consisting of large X-ray telescopes coupled to ultra-low background X-ray detectors, and a large, robust tracking system⁹⁰, and aims at the sensitivity shown in Fig. 2.

5 Summary

There is a strong physics case for the axion and ALPs. They occur naturally in many theoretically appealing ultraviolet completions of the SM. They are dark matter candidates and can explain the anomalous cosmic gamma ray transparency, soft X-ray excesses from galaxy clusters, and the unidentified 3.55 keV line from Andromeda and galaxy clusters. A significant portion of their parameter space will be tackled in this decade by experiments. Stay tuned!

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