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Approximate N⁵LO Higgs Boson Decay Width $\Gamma(H \rightarrow \gamma\gamma)$

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Abstract: The precision and predictive power of perturbative QCD (pQCD) prediction depends on both a precise, convergent, fixed-order series and a reliable way of estimating the contributions of unknown higher-order (UHO) terms. It has been shown that by applying the principle of maximum conformality (PMC), which applies the renormalization group equation recursively to set the effective magnitude of α_s of the process, the remaining conformal coefficients will be well matched with the corresponding α_s at each order, leading to a scheme-and-scale invariant and more convergent perturbative series. The PMC series, being satisfied with the standard renormalization group invariance, has a rigorous foundation. Thus it not only can be widely applied to virtually all high-energy hadronic processes, but also can be a reliable platform for estimating UHO contributions. In this paper, by using the total decay width $\Gamma(H \rightarrow \gamma\gamma)$ which has been calculated up to N⁴LO QCD corrections, we first derive its PMC series by using the PMC single-scale setting approach and then estimate its unknown N⁵LO contributions by using a Bayesian analysis. The newly suggested Bayesian-based approach estimates the magnitude of the UHO contributions based on an optimized analysis of the probability density distribution, and the predicted UHO contribution becomes more accurate when more loop terms have been known to tame the probability density function. Using the top-quark pole mass $M_t = 172.69$ GeV and the Higgs mass $M_H = 125.25$ GeV as inputs, we obtain $\Gamma(H \rightarrow \gamma\gamma) = 9.56504$ keV, and the estimated N⁵LO contribution to the total decay width is $\Delta\Gamma_H = \pm 1.65 \times 10^{-4}$ keV for the smallest credible interval of 95.5% degree of belief.

Keywords: perturbative quantum chromodynamics; principle of maximum conformality; bayes analysis



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1. Introduction

The ATLAS and CMS collaborations have discovered the Higgs boson in 2012 [1,2], consistent with the elementary particle suggested by the Standard Model (SM). The Higgs boson answers some of the most profound questions in physics, such as where the masses of the elementary particles and the W^\pm/Z^0 gauge bosons come from, how the electroweak phase transition governs the evolution of the early universe, etc. It is then crucial to verify and study the Higgs properties, either experimentally or theoretically.

Precise measurements of the Higgs boson production and decay channels provide critical tests of the SM and are vital in the exploration of new physics beyond the SM. Over the past decade, many new measurements on the Higgs boson properties have been performed by the collaborations at the LHC. Some new Higgs factories such as the International Linear Collider (ILC) [3], the Circular Electron Positron Collider (CEPC) [4] and the Future Circular Collider [5] have been designed to further improve the experimental precision on the Higgs properties. Thus, the Higgs boson is being moved from the object of a search to an exploration tool. Till now, almost all of the related measurements have been in agreement with the SM predictions within errors. As one of the most important decay channels of the Higgs boson, it has been shown that the process $H \rightarrow \gamma\gamma$ has an observable fraction $(2.50 \pm 0.20) \times 10^{-3}$ [6], which plays an important role in Higgs phenomenology.

Because the photon is massless, the process $H \rightarrow \gamma\gamma$ is a loop-induced process even at the leading-order level, whose amplitude can be decomposed into a bosonic contribution,

stemming from the W boson, and the fermionic contributions, respectively. More explicitly, its decay width can be written as

$$\Gamma(H \rightarrow \gamma\gamma) = \frac{M_H^3}{64\pi} \left| A_W + \sum_f A_f \right|^2, \quad (1)$$

where M_H is the Higgs mass, A_W is the contribution from the purely bosonic diagrams, and A_f is the contribution from the amplitudes with $f = (t, b, c, \tau)$, which corresponds to the top quark, the bottom quark, the charm quark, and the τ lepton, accordingly. The above equation can be further rewritten as [7]

$$\Gamma(H \rightarrow \gamma\gamma) = \frac{M_H^3}{64\pi} \left(A_{\text{LO}}^2 + A_{\text{EW}} \frac{\alpha}{\pi} \right) + R_n, \quad (2)$$

where α is the fine-structure constant, A_{EW} is the electroweak (EW) correction [8,9], A_{LO} is the leading-order (LO) contribution, and R_n represents the QCD corrections, in which n represents the QCD correction calculated up to the n_{th} -loop level. At present, the LO, the next-to-leading order (NLO), the $N^2\text{LO}$, the $N^3\text{LO}$, and the $N^4\text{LO}$ perturbative QCD (pQCD) corrections for $\Gamma(H \rightarrow \gamma\gamma)$ have been calculated in Refs. [8–24] under various approaches. In particular, the fermionic contribution which forms a gauge-invariant subset has been calculated up to the $N^4\text{LO}$ level in the large top-quark limit with $M_H \ll 2m_t$ [21]. Those improvements give us a good basis for achieving a precise pQCD prediction on $\Gamma(H \rightarrow \gamma\gamma)$. On the other hand, future precise measurements on the Higgs boson decay may determine the branching fraction of its decay into two photons up to a high precision of one percent [25]. Thus, to fully exploit future precise measurements, it is important to achieve a high-precision theoretical prediction as much as possible, as is the purpose of the present paper.

2. The $N^4\text{LO}$ -Level Prediction R_4 under the PMC and the Higher-Order Contribution Using a Bayesian Analysis

A valid prediction for a physical observable from quantum field theory should be independent of the choice of renormalization scheme and scale—this is the primary requirement of renormalization group invariance (RGI). Satisfying the RGI is a challenging problem for pQCD, since a truncated perturbation series does not automatically satisfy the requirements of the renormalization group. In the following we will take the higher-order QCD corrections to the process $H \rightarrow \gamma\gamma$ as an explicit example.

The perturbative series of the QCD correction R_4 up to the $\mathcal{O}(\alpha_s^5)$ level can be read from Refs. [20,21], which is given in n_f series with n_f being the active number of quark flavors. For the later convenience of applying the renormalization group equation (RGE) to set the effective magnitude of α_s , we express it as a $\{\beta_i\}$ series by using the general degeneracy relations of the QCD theory among different orders, e.g.,

$$R_4 = \sum_{i=1}^4 r_i (\mu_r^2 / Q^2) a^i(\mu_r) \quad (3)$$

$$\begin{aligned} &= r_{1,0} a(\mu_r) + [r_{2,0} + \beta_0 r_{2,1}] a^2(\mu_r) \\ &+ [r_{3,0} + \beta_1 r_{2,1} + 2\beta_0 r_{3,1} + \beta_0^2 r_{3,2}] a^3(\mu_r) \\ &+ [r_{4,0} + \beta_2 r_{2,1} + 2\beta_1 r_{3,1} + \frac{5}{2}\beta_1\beta_0 r_{3,2} + 3\beta_0 r_{4,1} \\ &+ 3\beta_0^2 r_{4,2} + \beta_0^3 r_{4,3}] a^4(\mu_r) + \mathcal{O}(a^5), \end{aligned} \quad (4)$$

where $a = \alpha_s / \pi$ and $Q = M_t$ (M_t being the top-quark pole mass), which represents the typical momentum flow of the process. The $\{\beta_i\}$ -functions have been calculated up to the five-loop level in the $\overline{\text{MS}}$ scheme [26–37]. The expansion coefficients $r_{i,j}$ in Equation (3) can be derived from the ones of Refs. [20,21] via proper transformations. In Refs. [20,21], the

perturbative expressions are given in the form of the $\overline{\text{MS}}$ -scheme top-quark running mass (m_t). Following the arguments of Ref. [38], we transform it into the perturbative series over the top-quark pole mass (M_t) with the help of the $\mathcal{O}(\alpha_s^5)$ -level relation between m_t and M_t [22] in order to avoid the confusion of applying the PMC-scale-setting procedures, e.g., only the RGE-involved β_i terms remain and are adopted for fixing the correct magnitude of the strong coupling and its argument, e.g., the PMC scale Q^* . The coefficients $r_{i,0}$ are conformal ones which are free of the renormalization scale μ_r , and the nonconformal coefficients $r_{i,j}(\neq 0)$ are functions of μ_r which can be expressed as

$$r_{i,j} = \sum_{k=0}^j C_j^k \hat{r}_{i-k,j-k} \ln^k(\mu_r^2/Q^2), \quad (5)$$

where $\hat{r}_{i,j} = r_{i,j}|_{\mu_r=Q}$. The RGE determines the running behavior of α_s and is scheme-dependent. By applying the principle of maximum conformality (PMC) [39–43], which applies the RGE recursively to set the effective magnitude of α_s of the process, the remaining conformal coefficients are well matched with the corresponding α_s at each order, leading to a scheme-and-scale invariant and convergent perturbative series free of divergent renormalization terms, cf. the reviews [44–46]. These reviews show that the PMC predictions respect all features of the renormalization group, and its prediction satisfies all the requirements of RGI; The commensurate scale relations, which relate physical observables to each other, ensure that PMC predictions are independent of the choice of renormalization scheme for any observable; And the transitivity and symmetry properties of the commensurate scales are the scale transformations of the renormalization “group”. Moreover, the PMC reduces in the Abelian limit to the Gell–Mann–Low method [47], and it provides a solid way to extend the well-known Brodsky–Lepage–Mackenzie (BLM) method [48] to all orders.

The PMC single-scale approach (PMCs) [49] determines an overall effective α_s (its argument is called the PMC scale) for the fixed-order predictions, and the resultant perturbative series provide a good basis for demonstrating that the PMC series is free of renormalization scale-and-scheme ambiguities up to any fixed order, being consistent with the fundamental renormalization group approaches [50,51]. It is noted that the single-scale approaches suggested in Ref. [49] are different from each other, but it has been demonstrated that the resultant pQCD series for both approaches are exactly the same. This equivalence indicates that by using the RGE to fix the value of the effective coupling is equivalent to requiring each loop’s terms to be scale-invariant simultaneously. Following the PMCs procedures [49], all the RGE-involved nonconformal terms of the above conventional series (4) of $R_4(\mu_r)$ should be removed from the series and adopted for fixing the correct magnitude of α_s of the process; one then obtains a scale-invariant conformal series. Up to the $N^4\text{LO}$ level, we have

$$R_4|_{\text{PMCs}} = \sum_{i=1}^4 \hat{r}_{i,0} a^i(Q_*) + \mathcal{O}(a^5), \quad (6)$$

where Q_* is the PMC scale, which can be determined by the following equation

$$\ln \frac{Q_{*,N^l\text{LL}}^2}{Q^2} = - \frac{\sum_{k=1}^{l+2} \sum_{i=1}^{l-k+2} \left[(-1)^i \Delta_{n,k}^{(i-1)} \hat{r}_{k+i,i} (n+k-1) a^k(Q_{*,N^l\text{LL}}) \right]}{\sum_{\eta=1}^{l+1} \sum_{k=1}^{l+2} \sum_{i=\eta}^{l-k+2} \left[(-1)^i (n+k-1) C_i^\eta \Delta_{n,k}^{(i-1)} \hat{r}_{k+i-\eta,i-\eta} L_{Q_{*,N^{l-1}\text{LL}}}^{\eta-1} a^k(Q_{*,N^l\text{LL}}) \right]} \quad (7)$$

$$= \sum_{i=0}^2 S_i a^i(Q_{*,N^l\text{LL}}), \quad (8)$$

where $L_{Q_{*,N^{l-1}\text{LL}}} = \ln Q_{*,N^{l-1}\text{LL}}^2/Q^2$. In the second line, e.g., Equation (8), we have expanded the series in the nominator and denominator as power series over $a = \alpha_s/\pi$, and their precision depend on how many loop terms for the pQCD approximation R_n are known. That is, by using R_2 , R_3 , and R_4 accordingly, the PMC scale can be fixed at the LL accuracy,

NLL accuracy, and N²LL accuracy, respectively. A similar PMC analysis on the N²LO level R_2 was conducted in Ref. [52], in which the LL-accuracy PMC scale was given. Up to the N⁴LO level, we need to know the first three functions $\Delta_{n,k}^{(0,1,2)}$, which are

$$\begin{aligned}\Delta_{n,k}^{(0)} &= 1, \\ \Delta_{n,k}^{(1)} &= -\frac{1}{2} \sum_{i=0}^{+\infty} (n+k+i) \beta_i a^{i+1}, \\ \Delta_{n,k}^{(2)} &= \frac{1}{3!} \sum_{i=0}^{+\infty} \sum_{j=0}^{+\infty} (n+k+i)(n+i+j+k+1) \\ &\quad \times \beta_i \beta_j a^{i+j+2}.\end{aligned}\quad (9)$$

And the functions S_i with $i = (0, 1, 2)$ are

$$S_0 = -\frac{\hat{r}_{2,1}}{\hat{r}_{1,0}} \quad (10)$$

$$S_1 = \frac{2(\hat{r}_{2,0}\hat{r}_{2,1} - \hat{r}_{1,0}\hat{r}_{3,1})}{\hat{r}_{1,0}^2} + \frac{(\hat{r}_{2,1}^2 - \hat{r}_{1,0}\hat{r}_{3,2})}{\hat{r}_{1,0}^2} \beta_0 \quad (11)$$

$$\begin{aligned}S_2 &= \frac{4(\hat{r}_{1,0}\hat{r}_{2,0}\hat{r}_{3,1} - \hat{r}_{2,0}^2\hat{r}_{2,1}) + 3(\hat{r}_{1,0}\hat{r}_{2,1}\hat{r}_{3,0} - \hat{r}_{1,0}^2\hat{r}_{4,1})}{\hat{r}_{1,0}^3} \\ &\quad + \frac{3(\hat{r}_{2,1}^2 - \hat{r}_{1,0}\hat{r}_{3,2})}{2\hat{r}_{1,0}^2} \beta_1 - \left[\frac{2\hat{r}_{2,0}^2\hat{r}_{2,1} - \hat{r}_{1,0}\hat{r}_{2,0}(6\hat{r}_{3,1} + 2\hat{r}_{3,2})}{\hat{r}_{1,0}^3} \right. \\ &\quad \left. - \frac{3(\hat{r}_{2,0}\hat{r}_{2,1}^2 + \hat{r}_{1,0}^2\hat{r}_{4,2})}{\hat{r}_{1,0}^3} \right] \beta_0 + \left[\frac{(\hat{r}_{1,0}\hat{r}_{2,0}\hat{r}_{3,2} - \hat{r}_{1,0}^2\hat{r}_{4,3})}{\hat{r}_{1,0}^3} \right. \\ &\quad \left. + \frac{2(\hat{r}_{1,0}\hat{r}_{2,0}\hat{r}_{3,2} - \hat{r}_{2,1}^3)}{\hat{r}_{1,0}^3} \right] \beta_0^2\end{aligned}\quad (12)$$

The predictive power of pQCD prediction also depends on a reliable way of estimating the contributions of unknown higher-order (UHO) terms. A Bayesian-based approach provides such a way of estimating the UHO contribution, which predicts the magnitude of the UHO terms based on an optimized analysis of the probability density distribution. The Bayesian analysis constructs probability distributions in which Bayes' theorem is used to iteratively update the probability as new information becomes available [53–57]. The interested reader may turn to Ref. [57] to learn the recent progresses on Bayesian analysis. We put the key formulas in the following for self-consistency.

If the perturbative approximation starts at the initial order $O(\alpha_s^l)$ and stops at the k_{th} order $O(\alpha_s^k)$, the corresponding perturbatively calculable physical observable can be schematically represented as

$$\rho_k = \sum_{i=l}^k c_i \alpha_s^i, \quad (13)$$

where the c_i 's are expansion coefficients. Replacing $\rho_k \rightarrow R_n$, $l \rightarrow 1$, and $c_i \rightarrow r_i(\hat{r}_{i,0})$ in the following formulas, we obtain the required formulas for the conventional (PMC) series of R_n . By taking three reasonable hypotheses, we obtain the probability density function (p.d.f) for the unknown higher-order coefficient c_n ,

$$f_c(c_n | c_l, \dots, c_k) = \begin{cases} \frac{n_c}{2(n_c+1)\bar{c}_{(k)}}, & |c_n| \leq \bar{c}_{(k)} \\ \frac{n_c \bar{c}_{(k)}^{n_c}}{2(n_c+1)|c_n|^{n_c+1}}, & |c_n| > \bar{c}_{(k)} \end{cases}. \quad (14)$$

where $\bar{c}_{(k)} = \text{Max}\{|c_l|, \dots, |c_k|\}$, and $n_c = k - l + 1$, which represents the number of known perturbative coefficients, c_l, \dots, c_k . Using Equation (14), one then derives the conditional p.d.f. for the uncalculated higher-order term $\delta_n = c_n \alpha_s^n$, ($n > k$). Especially for the one-order higher UHO-term with $n = k + 1$, the conditional p.d.f. of δ_{k+1} and ρ_{k+1} with given coefficients c_l, \dots, c_k , denoted by $f_\delta(\delta_{k+1}|c_l, \dots, c_k)$ and $f_\rho(\rho_{k+1}|c_l, \dots, c_k)$, respectively, read

$$f_\delta(\delta_{k+1}|c_l, \dots, c_k) = \left(\frac{n_c}{n_c + 1} \right) \frac{1}{2\alpha_s^{k+1} \bar{c}_{(k)}} \begin{cases} 1, & |\delta_{k+1}| \leq \alpha_s^{k+1} \bar{c}_{(k)} \\ \left(\frac{\alpha_s^{k+1} \bar{c}_{(k)}}{|\delta_{k+1}|} \right)^{n_c+1}, & |\delta_{k+1}| > \alpha_s^{k+1} \bar{c}_{(k)} \end{cases}, \quad (15)$$

$$f_\rho(\rho_{k+1}|c_l, \dots, c_k) = \left(\frac{n_c}{n_c + 1} \right) \frac{1}{2\alpha_s^{k+1} \bar{c}_{(k)}} \begin{cases} 1, & |\rho_{k+1} - \rho_k| \leq \alpha_s^{k+1} \bar{c}_{(k)} \\ \left(\frac{\alpha_s^{k+1} \bar{c}_{(k)}}{|\rho_{k+1} - \rho_k|} \right)^{n_c+1}, & |\rho_{k+1} - \rho_k| > \alpha_s^{k+1} \bar{c}_{(k)} \end{cases}. \quad (16)$$

One usually estimates the central value of ρ_{k+1} to be its expectation value $E(\rho_{k+1})$ and takes its uncertainty as its standard deviation, σ_{k+1} . The expectation value $E(\rho_{k+1})$ can be related to the expectation value of δ_{k+1} , i.e., $E(\rho_{k+1}) = E(\delta_{k+1}) + \rho_k$. For the present prior distribution, $E(\delta_{k+1}) = 0$, due to the fact that the symmetric probability distribution (15) is centered at zero. To predict the magnitude of δ_{k+1} consistently, it is useful to define a critical degree of belief (DoB), $p_c\%$, which equals the least value of $p\%$ that satisfies the following equations:

$$\rho_{i-1} + c_i^{(p)} \alpha_s^i \geq \rho_i + c_{i+1}^{(p)} \alpha_s^{i+1}, \quad (i = l + 1, \dots, k), \quad (17)$$

$$\rho_{i-1} - c_i^{(p)} \alpha_s^i \leq \rho_i - c_{i+1}^{(p)} \alpha_s^{i+1}, \quad (i = l + 1, \dots, k). \quad (18)$$

Thus, for any $p\% \geq p_c\%$, the error bars determined by the $p\%$ -credible intervals (CIs) provide consistent estimates for the magnitude of δ_{k+1} . The value of $p_c\%$ is nondecreasing when k increases. Practically, we adopt the smallest $p_s\%$ -CI so as to obtain a consistent and high DoB estimation, i.e.,

$$[E(\rho_{k+1}) - c_{k+1}^{(p_s)} \alpha_s^{k+1}, E(\rho_{k+1}) + c_{k+1}^{(p_s)} \alpha_s^{k+1}], \quad (19)$$

as a final estimate for ρ_{k+1} , where $p_s\% = \text{Max}\{p_c\%, p_\sigma\%\}$. Here, $p_\sigma\%$ represents the DoB for the 1σ interval, and $\rho_{k+1} \in [E(\rho_{k+1}) - \sigma_{k+1}, E(\rho_{k+1}) + \sigma_{k+1}]$.

2.1. Basic Numerical Results and Discussions

To perform the numerical calculation, we took the values of the input parameters from the Particle Data Group [6], e.g., the W -boson mass $M_W = 80.377$ GeV, the τ -lepton mass $M_\tau = 1.7769$ GeV, the b -quark pole mass $M_b = 4.78$ GeV, the c -quark pole mass $M_c = 1.67$ GeV, the t -quark pole mass $M_t = 172.69$ GeV, and the Higgs mass $M_H = 125.25$ GeV. The Fermi constant was $G_F = 1.1664 \times 10^{-5}$ GeV $^{-2}$ and the fine-structure constant was $\alpha = 1/137.036$. We assumed the running of α_s was at the four-loop level; the QCD asymptotic scale Λ_{QCD} was determined by using $\alpha_s(M_Z) = 0.1179$, which gave $\Lambda_{\text{QCD}}^{n_f=5} = 0.2072$ GeV.

For the process $H \rightarrow \gamma\gamma$, its QCD correction R_n under the $\overline{\text{MS}}$ -scheme was calculated up to the $N^4\text{LO}$ level. The initial fixed-order pQCD series was scheme-and-scale-dependent. (A way of achieving a scheme-and-scale-invariant prediction directly from the initial series, which is called the principle of minimum sensitivity (PMS), has been suggested in the literature. It assumes that all uncalculated higher-order terms give zero contribution and determines the optimal scheme and scale by requiring the slope of the pQCD series over the scheme-and-scale choices vanish. Since the PMS breaks the standard renormalization group invariance [45], it cannot be treated as a strict solution of conventional scheme-and-scale ambiguities, which, however, could be treated as an effective treatment.) As has been

discussed above, after applying the PMC, the resultant conformal series becomes scheme-and-scale-invariant. We present the scale-invariant conformal coefficients $\hat{r}_{i,0}$ ($i = 1, \dots, 4$) in Table 1, where the scale-dependent coefficients r_i at $\mu_r = M_H/2$, M_H and $2M_H$ are also presented for comparison.

Table 1. The $\overline{\text{MS}}$ coefficients $\hat{r}_{i,0}$ and r_i for R_4 . The coefficients r_i are also scale-dependent, and their values under three typical scale choices, e.g., $\mu_r = M_H/2$, M_H , and $2M_H$, are given for comparison.

	$i = 1$	$i = 2$	$i = 3$	$i = 4$
$r_i(\mu_r = M_H/2)$	1.4070	-0.9874	-0.4084	3.3437
$r_i(\mu_r = M_H)$	1.4070	0.2024	-1.6545	-0.3693
$r_i(\mu_r = 2M_H)$	1.4070	1.5282	-0.3456	-2.4065
$\hat{r}_{i,0}$	1.4070	1.3387	-3.6304	4.5695

Using the expansion coefficients of the QCD corrections R_2 , R_3 , and R_4 , the PMC scale can be fixed at the LL accuracy, NLL accuracy and N^2LL accuracy, respectively, and we obtain

$$Q_{*,\text{LL}} = 242.791 \text{ GeV}, \quad (20)$$

$$Q_{*,\text{NLL}} = 193.457 \text{ GeV}, \quad (21)$$

$$Q_{*,\text{N}^2\text{LL}} = 213.603 \text{ GeV}. \quad (22)$$

$|Q_{*,\text{N}^2\text{LL}} - Q_{*,\text{NLL}}| < |Q_{*,\text{NLL}} - Q_{*,\text{LL}}|$ indicates that the expansion series of $\ln Q_*^2/Q^2$ has a perturbative nature. Together with the fact that its higher-order terms will suffer from both α_s -power suppression and exponential suppression, the residual scale dependence of Q_* due to even higher-order terms of R_n will be highly suppressed, whose effects on the magnitude of α_s is negligible. The PMC predictions of R_2 , R_3 , and R_4 are

$$R_2|_{\text{PMC}} = 0.159493 \text{ keV}, \quad (23)$$

$$R_3|_{\text{PMC}} = 0.159969 \text{ keV}, \quad (24)$$

$$R_4|_{\text{PMC}} = 0.158517 \text{ keV}. \quad (25)$$

Table 2 shows the N^4LO QCD corrections $R_4 = \sum_{i=1}^4 \Delta_i$ under conventional and PMC scale settings, where Δ_i represents the individual decay width at the NLO, the N^2LO , the N^3LO or the N^4LO level, respectively. Three typical scales $\mu_r = M_H/2$, M_H , and $2M_H$ were adopted to show the conventional renormalization scale uncertainty. Table 2 shows that under conventional scale setting, the separate decay widths Δ_i are highly scale-dependent, and due to the large cancellation among different orders, the net scale dependence of the N^4LO prediction R_4 becomes small $\sim 1.28\%$ for $\mu_r \in [M_H/2, 2M_H]$. After applying the PMC, both Δ_i and R_4 are scale-independent. This confirms the observation that if the correct magnitude of α_s of a pQCD series has been determined by using the RG-involved $\{\beta_i\}$ -terms, indicating a good match of α_s with its expansion coefficients, one will achieve a precise scale-independent pQCD prediction. Such scale-independent nature of the pQCD approximation can be treated as its intrinsic perturbative property. Due to the good perturbative nature of the PMC series of R_n , the difference between the magnitudes of R_n and R_{n-1} becomes smaller with the increment of the given loop numbers.

Under the Bayesian approach, we predicted the magnitude of the unknown coefficient c_{i+1} from the known ones $\{c_1, \dots, c_i\}$ with $c_i \rightarrow r_i$ ($\hat{r}_{i,0}$) for conventional (PMC) series, respectively. Our results are listed in Table 3. From Tables 1 and 3, we can see that the exact values of $r_{i,0}$ ($i = 2, 3, 4$) and r_i ($i = 2, 3, 4$) lay within the predicted 95.5% CIs. Moreover, we obtained the smallest 95.5% credible intervals (CIs) for the perturbative coefficients $r_5(\mu_r = M_H)$ and $r_{5,0}$, which were $r_5 \in [-2.4023, 2.4023]$ and $r_{5,0} \in [-6.6348, 6.6348]$, respectively. The values from given series (“ECs”) are presented for comparison.

Table 2. The N^4 LO QCD corrections $R_4 = \sum_{i=1}^4 \Delta_i$ of $\Gamma(H \rightarrow \gamma\gamma)$ under conventional (Conv.) and PMC scale settings, respectively. Δ_i represents the individual decay width at the NLO, N^2 LO, N^3 LO or N^4 LO level, respectively. Three typical values $\mu_r = M_H/2$, M_H , and $2M_H$ are adopted to show the renormalization scale uncertainty.

		$i = 1$	$i = 2$	$i = 3$	$i = 4$	$R_4(\mu_r)$
$\Delta_i(\text{KeV}) _{\text{Conv.}}$	$\mu_r = M_H/2$	0.17589	-0.01543	-0.00080	0.00082	0.16048
	$\mu_r = M_H$	0.15830	0.00256	-0.00236	-0.00006	0.15845
	$\mu_r = 2M_H$	0.14467	0.01616	-0.00038	-0.00027	0.16018
$\Delta_i(\text{KeV}) _{\text{PMC}}$	$\mu_r \in [Q/2, 2Q]$	0.14744	0.01470	-0.00418	0.00055	0.15852

Table 3. The predicted smallest 95.5% CIs for the scale-dependent conventional coefficients $r_i(\mu_r)$ at the scale $\mu_r = M_H$ and the scale-invariant coefficients $\hat{r}_{i,0}$ ($i = 3, 4, 5$) of $R_n(\mu_r = M_H)$ via the Bayesian approach, where $M_H = 125.25$ GeV. The values from given series (“ECs”) are presented for comparison.

	$r_2(M_H)$	$r_3(M_H)$	$r_4(M_H)$	$r_5(M_H)$
CI	$[-15.6334, 15.6334]$	$[-3.8294, 3.8294]$	$[-2.9303, 2.9303]$	$[-2.4023, 2.4023]$
	0.2024	-1.6545	-0.3693	-
EC	$\hat{r}_{2,0}$	$\hat{r}_{3,0}$	$\hat{r}_{4,0}$	$\hat{r}_{5,0}$
	$[-15.6334, 15.6334]$	$[-3.8294, 3.8294]$	$[-6.4298, 6.4298]$	$[-6.6348, 6.6348]$
	1.3387	-3.6304	4.5695	-

Using the estimated $r_5(M_H)$ and $\hat{r}_{5,0}$, the error of Γ_H caused by the UHO-terms for conventional series and PMC series under the Bayesian approach (B.A.) were

$$\Delta\Gamma_H|_{\text{Conv.}}^{\text{UHO}} = \pm 8.523 \times 10^{-5} \text{ keV}, \quad (26)$$

$$\Delta\Gamma_H|_{\text{PMC}}^{\text{UHO}} = \pm 1.65 \times 10^{-4} \text{ keV}. \quad (27)$$

By further taking $\mu_r \in [M_H/2, 2M_H]$, the conventional series also had the following scale uncertainty

$$\Delta\Gamma_H|_{\text{Conv.}}^{\mu_r} = (+2.03 \times 10^{-3}) \text{ keV}. \quad (28)$$

Then, as a combination, the net errors caused by the N^5 LO UHO terms in conventional and PMC series were

$$\Delta\Gamma_H|_{\text{Conv.}} = (+2.03 \times 10^{-3}) \text{ keV} \quad (29)$$

$$\Delta\Gamma_H|_{\text{PMC}} = \pm 1.65 \times 10^{-4} \text{ keV} \quad (30)$$

where $\mu_r \in [M_H/2, 2M_H]$.

In addition, for the more precise PMC series, we also adopted another usual way of estimating UHO contributions, e.g., the Padé approximation approach (PAA) [58–60] to estimate the UHO terms of R_n . The PAA works when we know enough higher orders, e.g., $n \geq 2$ for the present case. The PAA has an intrinsic error due to the existence of different types of generating functions [61], and we took the result of the $[0/n-1]$ type as its central value and the results of other types were treated as its uncertainty. More explicitly, to estimate the N^3 LO magnitude from the given N^2 LO series, we used the $[0/1]$ -type generating function; to estimate the N^4 LO magnitude from the given N^3 LO series, we used the $[0/2]$ -type and $[1/1]$ -type generating functions; to estimate the N^5 LO magnitude from the given N^4 LO series, we used the $[0/3]$ -type, $[1/2]$ -type, and $[2/1]$ -type generating functions, etc. We put the results in Figure 1, where the “Exact Values” together with the Bayesian approach (B.A.) and Padé approximation approach (PAA) ones are presented.

Figure 1 shows that for the B.A. approach, the “exact” value are always within the predicted error band, the predicted one-order-higher UHO error band is always within the predicted one-order-lower UHO error band, and the predicted UHO values become more accurate when more loop terms are known. Thus, if one has enough higher-order information to tame the probability density function, one may determine the precise contribution of the UHO terms. For the PAA, the “exact” N^4 LO value is outside of the predicted error bar, and the predicted N^5 LO error bar becomes better and is consistent with the B.A. one. In this sense, at least for the present case, the B.A. approach is more effective than the PAA.

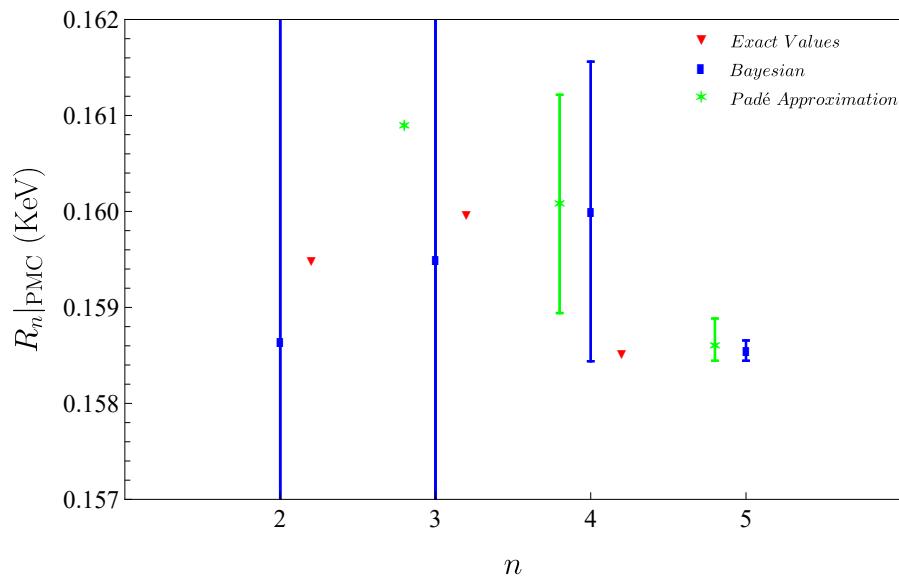


Figure 1. The predicted values for the pQCD correction $R_n|_{PMC}$ under the Padé approximation approach (PAA) and Bayesian approach (B.A.) at different orders, respectively. The blue rectangles together with the error bars are for B.A., the green error bars are brought by different types of PAAs, and the exact values of the $R_n(M_H)|_{PMC}$ at different orders, respectively.

>From Equation (2), there are other error sources such as ΔM_H , Δm_t and $\Delta \alpha_s(M_Z)$ for the total decay width $\Gamma(H \rightarrow \gamma\gamma)$. For this purpose, we took $\Delta M_H = \pm 0.17$ GeV, $\Delta m_t = \pm 0.30$ GeV, and $\Delta \alpha_s(M_Z) = \pm 0.0009$ GeV [6] to show their effects. When discussing the error caused by one parameter, the other parameters were fixed as their center values. Thus, we had

$$\Delta \Gamma_H|_{Conv.}^{\Delta M_H} = (+5.455 \times 10^{-2}) \text{ keV}, \quad (31)$$

$$\Delta \Gamma_H|_{PMC}^{\Delta M_H} = (+5.453 \times 10^{-2}) \text{ keV}, \quad (32)$$

$$\Delta \Gamma_H|_{Conv.}^{\Delta m_t} = (+6.999 \times 10^{-4}) \text{ keV}, \quad (33)$$

$$\Delta \Gamma_H|_{PMC}^{\Delta m_t} = (+7.004 \times 10^{-4}) \text{ keV}, \quad (34)$$

$$\Delta \Gamma_H|_{Conv.}^{\Delta \alpha_s(M_Z)} = (+1.071 \times 10^{-3}) \text{ keV}, \quad (35)$$

$$\Delta \Gamma_H|_{PMC}^{\Delta \alpha_s(M_Z)} = (+1.061 \times 10^{-3}) \text{ keV}. \quad (36)$$

By adding all the mentioned errors in quadrature, our final results for the total decay Γ_H of $H \rightarrow \gamma\gamma$ using the B.A. approach were

$$\Gamma_H|_{Conv.}^{B.A.} = 9.56497^{+0.05461}_{-0.05424} \text{ keV}, \quad (37)$$

$$\Gamma_H|_{PMC}^{B.A.} = 9.56504^{+0.05455}_{-0.05422} \text{ keV} \quad (38)$$

whose net errors were 1.138% and 1.137%. This shows that since the QCD correction was calculated up to the N^4 LO level, the main errors were dominated by ΔM_H . (As a rough

estimation, by setting all input parameters to be their central values, the magnitudes of the EW or the QCD-EW mixing correction are $\sim \mathcal{O}(-2.1 \times 10^{-1}\text{keV})$ or $\sim \mathcal{O}(+1.6 \times 10^{-1}\text{keV})$, and then the errors caused by one-order-higher EW or QCD-EW corrections should be $\sim \mathcal{O}(-1.5 \times 10^{-3}\text{keV})$ or $\sim \mathcal{O}(+1.2 \times 10^{-3}\text{keV})$ because of the α suppression. Furthermore, due to different sign of the EW and QCD-EW mixing corrections, their net error becomes $\sim \mathcal{O}(3.0 \times 10^{-4}\text{keV})$. Those magnitudes are also smaller than the dominant error caused by ΔM_H .

2.2. The Fiducial Cross Section of $\sigma_{\text{fid}}(pp \rightarrow H \rightarrow \gamma\gamma)$

As an application of the $H \rightarrow \gamma\gamma$ decay width, we estimated the “fiducial cross section” of the process $pp \rightarrow H \rightarrow \gamma\gamma$. The fiducial cross section σ_{fid} can be written as

$$\sigma_{\text{fid}}(pp \rightarrow H \rightarrow \gamma\gamma) = \sigma_{\text{Incl}} \mathcal{B}_{H \rightarrow \gamma\gamma} A \quad (39)$$

where A is the acceptance factor, whose value for different collision energies can be found in Ref. [62]. $\mathcal{B}_{H \rightarrow \gamma\gamma}$ represents the branching ratio of $H \rightarrow \gamma\gamma$. By using $\Gamma(H \rightarrow \gamma\gamma)$ with a conventional-scale-setting approach, the LHC-XS group gave $\mathcal{B}_{H \rightarrow \gamma\gamma} = 0.00227^{+0.00206}_{-0.00208}$ [63]. The inclusive cross section σ_{Incl} predicted by the LHC-XS group is given in Ref. [64]. The results are $\sigma_{\text{fid}}(pp \rightarrow H \rightarrow \gamma\gamma)|_{\text{LHC-XS}} = 24.63^{+2.55}_{-2.50} \text{ fb}$, $30.93^{+3.44}_{-3.33} \text{ fb}$, and $65.86^{+6.58}_{-6.33} \text{ fb}$ for the proton–proton center-of-mass collision energy $\sqrt{S} = 7, 8$, and 13 TeV , respectively, which has been measured by the ATLAS and CMS collaborations with increasing integrated luminosities [62,65–70]. Taking the same inputs as those of Refs. [63,64,71], e.g., $M_H = 125 \text{ GeV}$ and $M_t = 173.3 \text{ GeV}$, and using the QCD corrections up to the N^4LO level, we obtained $\sigma_{\text{fid}}(pp \rightarrow H \rightarrow \gamma\gamma)|_{\text{PMC}} = 30.1^{+2.3}_{-2.2} \text{ fb}$, $38.3^{+2.9}_{-2.8} \text{ fb}$, and $85.5^{+5.7}_{-5.3} \text{ fb}$ for the proton–proton center-of-mass collision energy $\sqrt{S} = 7, 8$, and 13 TeV , respectively. As an intuitive comparison of the experimental data and theoretical results, we present the results in Figure 2 (Due to the resummation of the given $\{\beta_i\}$ -type terms, which determine the precise magnitudes of α_s at different proton–proton center-of-mass collision energies, the PMC predictions are always larger than the LHC-XS predictions using conventional series.). It shows that when $\sqrt{S} = 7$ or 8 TeV , the theoretical results are consistent with the experimental measurements, and when $\sqrt{S} = 13 \text{ TeV}$, the measured values of ATLAS and CMS differ significantly, and the theoretical results are closer to the data of CMS.

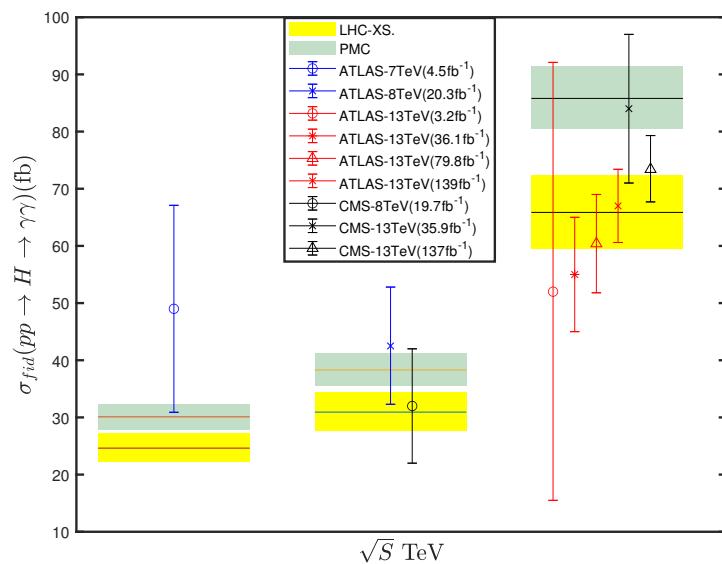


Figure 2. The fiducial cross section $\sigma_{\text{fid}}(pp \rightarrow H \rightarrow \gamma\gamma)$ using $\Gamma(H \rightarrow \gamma\gamma)$ up to the N^4LO level. The LHC-XS prediction, the ATLAS measurements [62,65–67], and the CMS measurement [68–70] are presented as a comparison.

3. Summary

By using the PMC-scale-setting approaches, all nonconformal terms were adopted to set the correct magnitude of α_s with the help of the RGE, and the resultant pQCD series became more precise without the conventional scheme-and-scale independence. In this paper, we calculated the decay width $\Gamma(H \rightarrow \gamma\gamma)$ up to the $N^4\text{LO}$ QCD corrections. A Bayesian approach was applied to estimate the uncalculated $N^5\text{LO}$ contribution, which was only about $\pm 1.65 \times 10^{-4}$ keV for the case of the smallest 95.5% credible interval. After taking all the mentioned errors into consideration, we predicted $\Gamma_H|_{\text{PMC}}^{\text{B.A.}} = 9.56504^{+0.05455}_{-0.05422}$ keV. Thus, by using the Bayesian approach, one can consistently obtain high-reliability estimations of UHO contributions by using convergent and scale-independent PMC series, greatly improving the prediction ability of pQCD.

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