Black hole thermodynamics in finite time

Christine Gruber*

*Institute of Physics, Carl v. Ossietzky University of Oldenburg,
26129 Oldenburg, Germany*
*E-mail: christine.gruber@uni-oldenburg.de

Finite-time thermodynamics provides the means to revisit ideal thermodynamic equilibrium processes in the light of reality and investigate the energetic "price of haste", i.e. the consequences of carrying out a process in finite time, when perfect equilibrium cannot be awaited due to economic reasons or the nature of the process. Employing the formalism of geometric thermodynamics, a lower bound on the energy dissipated during a process is derived from the thermodynamic length of that process. The notion of length is hereby defined via a metric structure on the space of equilibrium thermodynamics, spanned by a set of thermodynamic variables describing the system. Since the aim of finite-time thermodynamics is to obtain realistic limitations on idealized scenarios, it is a useful tool to reassess the efficiency of thermodynamic processes. We examine its implications for black hole thermodynamics, in particular scenarios inspired by the Penrose process, a thought experiment by which work can be extracted from a rotating black hole. We consider a Kerr black hole which, by some mechanism, is losing mass and angular momentum. Thermodynamically speaking, such a process is described in the equilibrium phase space of the black hole, but in reality, it is neither reversible nor infinitely slow. We thus calculate the dissipated energy due to non-ideal finite-time effects.

Keywords: Black hole thermodynamics; Kerr black hole; finite-time thermodynamics; MG14 Proceedings.

1. Introduction

Even though experimentally black holes have so far remained evasive and mysterious astrophysical objects, black hole thermodynamics is an active and innovative field of research. In the recent years black holes have been approached from an increasingly engineering point of view. Investigations in this direction include studies of black hole stability with respect to thermodynamic fluctuations or phase transitions, and proposals of heat engines involving black holes.

One of the first ideas to utilize black holes as engines to gain energy from was put forward in the 1970s, suggesting the extraction of angular momentum from a rotating black hole[1]. Since then, many different mechanisms have been contemplated in this context, usually employing a specific mechanical process (for a recent review, see Ref. 2). In this work, we suggest an approach based on thermodynamics, retaining maximum generality which does not require fixing a particular method of energy extraction.

We will furthermore include the effects of finite time into our analysis, i.e. consider realistic thermodynamic processes[3] as opposed to idealised reversible ones. In reversible thermodynamics it is assumed implicitly that processes happen ideally and are represented by a sequence of perfect equilibrium states along the path - which requires an infinite amount of time. In reality however, processes are carried out on finite time scales, and thus at each local equilibrium state along the sequence

a small amout of energy is lost, i.e. a small amount of entropy ΔS is produced or energy ΔU is dissipated in each infinitesimal equilibration. In the ideal reversible limit, these losses become zero. A measure of the amount of dissipation along a process is given by the notion of thermodynamic length of a process.

2. Finite-time thermodynamics

The thermodynamic length of a process can be defined in the framework of geometric thermodynamics, using the definition of a metric in the *equilibrium* phase space spanned by the state variables describing the system. We will here use the thermodynamic metric of Ruppeiner[4], which is defined as the Hessian of the entropy of the system,

$$ds_S = \mathbf{g}_S \, dq^a \otimes dq^b = \frac{\partial^2 S}{\partial q^a \partial q^b} \, dq^a \otimes dq^b \,, \tag{1}$$

with respect to its extensive thermodynamic parameters $\mathbf{f} = \{q^a\}$.

The length of a path with respect to the entropy metric (1) in the equilibrium space of the system is then defined as

$$L_S = \int_a^b ds_S = \int_a^b \sqrt{d\mathbf{f}^\top \mathbf{g} \, d\mathbf{f}} \,. \tag{2}$$

or, when parametrized by a parameter ξ,

$$L_S = \int_0^{\xi_{max}} \sqrt{\dot{\mathbf{f}}^\top \mathbf{g} \dot{\mathbf{f}}} \, d\xi \,. \tag{3}$$

The entropy production along a path between the points a and b can be derived as[5]

$$\Delta S = \frac{1}{2} \int_0^{\xi_{max}} \epsilon_S \, \dot{\mathbf{f}}^\top \mathbf{g} \, \dot{\mathbf{f}} \, d\xi =: \bar{\epsilon}_S \, \mathcal{J}_S \,, \tag{4}$$

where $\mathcal{J}_S$ is the generalized action of the path, and ϵ_S ($\bar{\epsilon}_S$) is a measure of the infinitesimal entropy produced in each step (averaged over all steps). By comparison with the definition of the thermodynamic length, and considering the Cauchy-Schwarz inequality, a bound for the entropy generated along the path is found as

$$\mathcal{J}_S \geq \frac{L_S^2(\xi_{max})}{2\xi_{max}} \,. \tag{5}$$

From this bound, we see that for $\xi_{max} \to \infty$, the generalized action and thus the dissipated entropy go to zero, $\mathcal{J}_S \to 0$. Further, the shortest thermodynamic length between two points a and b can be achieved by following the geodesics of $\mathbf{g}$,

$$\mathcal{J}_S \geq \frac{L_S^2}{2\xi_{max}} \geq \frac{L_{S,\,geod.}^2}{2\xi_{max}} \,. \tag{6}$$

Ultimately, the minimum entropy for a fixed path between two points is obtained by following the path with constant speed.

3. Energy extraction from a Kerr black hole

The entropy S of a Kerr black hole is related to its mass M and its angular momentum J by the relation

$$S(M, J) = 2M^2 \left(1 + \sqrt{1 - \frac{J^2}{M^4}} \right) , \qquad (7)$$

and it obeys the first law of black hole thermodynamics,

$$dS = \frac{1}{T}dM - \frac{\Omega}{T}dJ , \qquad (8)$$

with the entropy as the fundamental potential, and the extensive thermodynamic parameters $\mathbf{f} = \{M, J\}$.

The thermodynamic length is calculated as

$$L_S = \int_0^{\xi_{max}} \left[-2\dot{M}^2 - \frac{6M^2\dot{M}^2 - \dot{J}^2}{\sqrt{M^4 - J^2}} + \frac{4M^6\dot{M}^2 - J^2\dot{J}^2}{\sqrt{M^4 - J^2}^3} \right]^{1/2} d\xi . \qquad (9)$$

In the following, we investigate two types of processes and calculate the respective thermodynamic lengths.

Keeping the angular momentum constant, we will first vary the black hole mass. We will consider processes starting from the extremal limit, $M_0 = \sqrt{J}$, and increasing the mass up to the boundary of the stable region for a Kerr black hole, $M_+ \leq M_1 = \sqrt{J/\alpha_{min}}$, where $\alpha = J/M^2$, and $\alpha_{min} = 0.6813$. The thermodynamic length is given by the integral

$$L_S = \int_{M_0}^{M_+} \left[4 - \frac{4M^6 - 12M^2J^2}{\sqrt{M^4 - J^2}^3} \right]^{1/2} dM , \qquad (10)$$

which was then solved numerically. The left graph in Fig. 1 shows the resulting length as a function of M_+ for various values of the angular momentum.

The second type of process considers a constant black hole mass, while decreasing its angular momentum from an upper value $J_+ \leq J_1 = M^2$ down to zero, $J_0 = 0$, and thus describes the classic example of a Penrose process, where all the rotational energy is extracted from a Kerr black hole. The thermodynamic length in this case can be calculated analytically as

$$L_S = \int_{J_0}^{J_+} \left[\frac{M^4}{\sqrt{M^4 - J^2}^3} \right]^{1/2} dJ = \frac{J_+}{M} \, _2F_1 \left[\frac{1}{2}, \frac{3}{4}, \frac{3}{2}, \frac{J_+}{M^2} \right] . \qquad (11)$$

The result is shown in the right graph in Fig. 1, where the length is given as a function of J_+.

From the resulting curves in Fig. 1, we deduce that close to the extremal limit the thermodynamic length grows stronger than further away from it, and thus the dissipation increases more strongly close to the extremal limit. For an efficient energy extraction therefore it is preferrable to operate away from the extremal limit, to minimize the dissipative losses during the process.
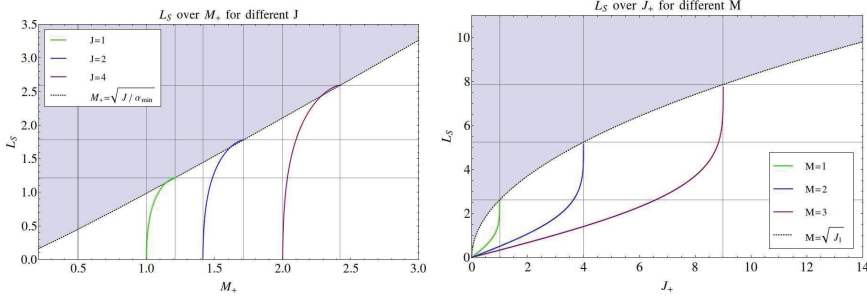
Fig. 1.   (color online) Thermodynamic lengths for processes with constant J (left) and M (right).

4. Conclusions & Outlook

In this work, we considered Penrose-inspired processes, i.e. the extraction of energy fom black holes, from a thermodynamic point of view, taking into account the dissipative losses generated due to finite-time effects. The entropy generated along a process can be bound by the thermodynamic length of that process in the geometric space of equilibrium thermodynamics. We calculate the thermodynamic length of processes with constant angular momentum and mass of a Kerr black hole, and find that dissipative losses grow more strongly close to the extremal limit of the black hole. A next step is to calculate the length of geodesic paths in the thermodynamic geometry, minimizing the dissipative losses, which has been done in Ref. 2.

Acknowledgments

References

1. R. Penrose and R. Floyd, Extraction of rotational energy from a black hole, *Nature* **229**, 177 (1971).
2. A. Bravetti, C. Gruber and C. S. Lopez-Monsalvo, Thermodynamic optimization of a Penrose process: an engineers' approach to black hole thermodynamics, *arXiv:1511.06801 [gr-qc]* .
3. B. Andresen, Current trends in finite-time thermodynamics, *Angewandte Chemie International Edition* **50**, 2690 (2011).
4. G. Ruppeiner, Thermodynamics: A Riemannian geometric model, *Phys. Rev. A* **20**, p. 1608 (1979).
5. P. Salamon and R. S. Berry, Thermodynamic length and dissipated availability, *Physical Review Letters* **51**, p. 1127 (1983).