

A Study of the Role of the PDF uncertainty on the LHC *W*-boson mass measurement

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The determination of the W -boson mass through an analysis of the decay charged-lepton transverse momentum distribution has a sizable uncertainty due to the imperfect knowledge of the relevant parton distribution functions (PDFs). In this paper, a quantitative assessment of the W -boson mass uncertainty at the LHC resulting from the PDF uncertainty is examined. We use the **CT14 NNLO** PDFs with a NNLL + NNLO calculation (**ResBos**) to simulate the W -boson production and decay kinematics. The uncertainty of the W -boson mass determination is then determined as a function of the boson and lepton kinematics. For W^+ production using $P_T^W < 15$ GeV and $35 < P_T^l$ (GeV) < 45 , PDF uncertainties lead (at the 68% CL) to a mass uncertainty for the M_W^+ of $^{+16.0}_{-17.5}$ MeV, $^{+13.9}_{-14.8}$ MeV, and $^{+12.2}_{-19.2}$ MeV, are determined at 7 TeV, 8 TeV and 13 TeV respectively. The uncertainties of M_W^- for the same cuts are found to be $^{+15.9}_{-15.6}$ MeV, $^{+15.0}_{-12.7}$ MeV and $^{+14.8}_{-15.3}$ MeV, at 7 TeV, 8 TeV and 13 TeV respectively.

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I. INTRODUCTION

A. Importance of the precision W -boson mass determination

The W -boson mass (M_W) is a fundamental parameter of the Standard Model (SM). M_W can be both directly measured, and computed indirectly from electroweak precision constraints. There is an interplay between M_W , the top quark mass (M_t), and the Higgs mass (M_H) when calculating electroweak precision observables [1]. Currently, there is some tension between the best fit value from the electroweak precision data and the direct measurement. Electroweak precision tests are most sensitive to M_W . If a statistically significant discrepancy between the indirect measurement of M_W and the direct measurement is found, this would be strong evidence for new physics beyond the SM such as the Minimal Supersymmetric Standard Model [2]. Therefore, more precise determinations of M_W have been an important goal for collider physics programs at both the Tevatron and the LHC, both for testing the consistency of the SM and for exploring the possibility of new physics beyond the SM.

B. Current limits on M_W and on PDF uncertainty

Currently, the best limits for the W -boson mass come from the measurements at the Tevatron, and from a recent result from ATLAS. The W -boson mass has been measured at CDF (80387 ± 19 MeV) [3] and D0 (80367 ± 26 GeV) [4], with a combined Tevatron average (excluding LEP data) of 80387 ± 16 MeV [5]. A more recent measurement for M_W at ATLAS has yielded a value of 80370 ± 19 MeV [6]. CMS has not reported a result to date. One of the major systematic uncertainties for the direct measurement arises from an imperfect knowledge of the parton distribution functions (PDFs) relevant for W -boson production. Using the M_H measurement [7], the combined value for M_t [8], and the SM precisely measured parameters of the fine-structure constant α , the Fermi constant G_μ , and the Z -boson mass (M_Z), results in a predicted (indirect) mass for the W -boson of 80362 ± 8 MeV [9] and 80358 ± 8 MeV [10]. This indirect uncertainty limit provides a goal for the desired experimental precision for the W -boson mass.

C. PDFs and their uncertainties

Parton distribution functions cannot at present be calculated from first principles, but must be determined by data. This determination requires the use of data from a wide variety of processes and experiments, both for the determination of the central PDFs and of their uncertainties. There are several collaborations dedicated to this endeavor, among them: CTEQ-TEA [11], MMHT [12], and NNPDF [13]. PDFs have been produced at LO, NLO and (more recently) NNLO, in the strong coupling constant α_s . The highest precision predictions at the LHC require the use of NNLO PDFs. Previously, NNLO PDFs from the CTEQ collaboration (CT10 [14]), have been used in a determination of the W -mass uncertainty [15]. In this paper, we update those predictions using the NNLO PDFs of CT14 [11]. The CT14 PDFs include data sets from the LHC at $\sqrt{s} = 7$ TeV for the first time, as well as updated data from the Tevatron and from HERA. The most important aspect of the LHC data sets for this analysis is their ability to impose constraints on the light quark and anti-quark PDFs at parton x-values appropriate for W -boson production at the LHC¹.

PDF uncertainties are calculated for CT14 through the use of Hessian eigenvectors [17]. CT14 NNLO has 56 Hessian error PDFs, corresponding to 28 eigenvectors, which result from the 28 free PDF parameters in the CT14 fit. In comparison, CT10 NNLO had 50 error sets, resulting from using 25 free parameters. The increased number of parameters is a result of a more flexible parametrization used in the CT14 global PDF fit. The purpose of the present analysis is to provide a quantitative assessment of the M_W uncertainty resulting from PDF uncertainties using the CT14 NNLO PDFs.

II. METHODS FOR DETERMINING M_W

The production of W/Z bosons is one of the most well-studied examples of hard scattering processes at hadron colliders [18]. After production, the W boson decays into jets, or into a lepton-neutrino pair. Decays into the former channel are difficult to resolve due to large

¹ It is often useful to include information from more than one PDF set. CT14 was included in such a combination (PDF4LHC15), along with MMHT2014 and NNPDF3.0 [16]. This combination in principle provides more information on the PDF uncertainty for the W -boson production. However, as the reference indicates, the central PDFs and the PDF uncertainties for the three PDFs included in the combination are consistent for the W -boson production initial states in the relevant kinematic region.

contributions from QCD dijet background processes. On the other hand, the $W \rightarrow e\nu_e$ and $W \rightarrow \mu\nu_\mu$ channels allow for precise measurements of prompt, energetic, and isolated electrons and muons. The decay of $W \rightarrow \tau\nu_\tau$ is also not included for precision measurements due to the complex nature of the τ decay.

The only observables directly measured by the detectors are the momenta of the leptons (P^l) and of the hadrons (P^{hadrons}) produced in association with the W -boson; the latter is referred to as the hadronic recoil. In addition, the hadronic recoil's transverse momentum defines the negative transverse momentum of the W -boson itself (P_T^W). Since the neutrinos escape the detector without interaction, M_W cannot be reconstructed on an event-by-event basis. However, the sum of the transverse momentum of all particles in the event should sum to zero in the absence of any particles evading detection and detector resolution effects. Therefore, the neutrino's transverse momentum (P_T^ν) can be inferred indirectly from the transverse energy missing from the event, $\cancel{E}_T \equiv -(P_T^l + P_T^{\text{hadrons}})$. The transverse mass of the W -boson (M_T^W), introduced in [19, 20], is defined as:

$$M_T^W = \sqrt{2P_T^l \cancel{E}_T (1 - \cos \Delta\phi_{l\nu})}, \quad (1)$$

where $(\Delta\phi_{l\nu})$ is the azimuthal angle between the lepton and the neutrino (or the missing transverse energy).

It is potentially possible to measure the W -boson mass using any one of the three kinematic variables, M_T^W , P_T^l and P_T^ν . In practice, the experimental resolution for P_T^ν does not allow for a competitive measurement of the W -boson mass (but it can still be useful as a cross-check). The other two variables can, and have been, used though at both the Tevatron and the LHC.

The extraction of M_W is obtained through the use of templates generated from a highly parametrized Monte Carlo (MC) simulation. At leading order (LO), the W -boson is produced with zero transverse momentum, and thus the charged lepton and neutrino are always back-to-back. Therefore, in the LO calculation, in the limit of zero width, and with a perfect detector, M_T^W and P_T^l would have extremely sharp Jacobian peaks exactly at M_W and $M_W/2$ respectively. QCD radiation, and the impact of detector resolution, results in a shift in the location of the Jacobian peak and a broadening the distributions.

A. Differences between Tevatron and LHC

Compared to the Tevatron, the LHC experiments benefit from larger signal and calibration samples. For luminosities in the multi- fb^{-1} range, the data samples are larger by an order of magnitude compared to the corresponding samples used at the Tevatron, and thus the statistical errors are significantly smaller. Moreover, and given the precisely measured value of $M_Z(91187.6 \pm 2.1 \text{ MeV})$ [21] and the clean leptonic final state, the $Z \rightarrow l^+l^-$ processes are able to be used to model the detector's response to $W \rightarrow l\nu$ and to validate the analysis synopsis [22].

Uncertainties in the PDFs are the dominant source of error for the extraction of M_W at the Tevatron (a complete list of these uncertainties can be found in Tables XIV and VI in Ref. [3, 4] respectively). The PDF uncertainties for W -boson production were expected to be larger at the LHC [23], due to the smaller parton x sampled at the LHC. There is additional uncertainty from a pp collider vs. a $p\bar{p}$ collider due to larger contributions from sea quarks (see Table 3 of Ref. [6]).

Whereas W -boson production at the Tevatron is charge symmetric, the W^+ -boson production rate at the LHC exceeds that of W^- -boson by about 40%. Moreover, the second generation quarks contribute only approximately 5% of the overall W -boson production rate at the Tevatron, while at the LHC ($\sqrt{s} = 7 \text{ TeV}$) this rate is approximately 25% of the overall W -boson production rate [6]. This fraction continues to increase as the center-of-mass energy increases. The uncertainty on the strange and charm quarks is larger than those on the light quarks(anti-quarks), and thus result in a proportionally larger contribution to the W boson mass determination.

Compared to P_T^l , the M_T^W measurement at the LHC is affected by larger experimental systematic uncertainties, due, for example, to high pile-up energy deposited in the detector from the additional proton-proton interactions in each bunch crossing. This results in a degradation of the resolution of the measurement of P_T^{hadrons} that scales roughly as the square root of the total hadronic energy in the event [24]. The limiting factor to balance, at present, between the experimental uncertainty (M_T^W) and the theoretical uncertainty (P_T^l) depends on the ability to develop pile-up mitigation techniques [25]. Although both techniques have been used at the LHC [6], the greater discriminatory power lies with the use of the P_T^l distribution, as shown in Table 10 and in Fig. 23 of Ref. [6]. Therefore, this

paper will focus on determining the PDF uncertainties for the P_T^l distribution.

B. Resummed QCD

At NLO, the W -boson can recoil against one parton, and at NNLO, against two partons, thus acquiring a non-zero transverse momentum. For P_T^W much less than M_W , soft gluon radiation has a large impact on the P_T^W , and resummation calculations are necessary to provide a good description of the distribution. At very low P_T values, non-perturbative effects also become important, and must also be taken into account in the resummation calculations. In practice, most W -boson events used in the determination of M_W [6] are produced with low P_T^W values (smaller than 30 GeV). This means that, in principle, a resummed generator with non-perturbative effects such as `ResBos` is preferable to fixed-order calculations [26].

Fixed-order QCD predictions work well for W boson production at intermediate and large P_T . However, the W/Z -bosons are predominantly produced at low P_T , where the fixed-order cross section behaves as

$$\frac{d\sigma}{dP_T^2} \sim \frac{1}{P_T^2} \sum_{n=1}^{\infty} {}_k c_n \alpha_s^n(Q) \sum_{k=0}^{2n-1} \ln^k \frac{P_T^2}{Q^2}, \quad (2)$$

in the limit of $P_T \rightarrow 0$. It can be seen that the above equation has an unphysical divergence when $\ln^k(P_T^2/Q^2) \rightarrow 0$. This results in large logarithms ($\alpha_s^n \ln^k(P_T^2/Q^2) > 1$), making perturbative calculations unreliable [27]. This unphysical divergence is addressed through the resummation procedure, in which the logarithmic terms are resummed to all orders in the α_s expansion. The general formalism for transverse momentum resummation was first introduced by Collins, Soper, and Sterman (CSS) [28].

$$\frac{d\sigma}{dQ^2 dy dq_t^2} = \sigma_0 \left(\int \frac{d^2 b}{(2\pi)^2} e^{iq_t \cdot b} \widetilde{W}(b) + Y \right), \quad (3)$$

where

$$\widetilde{W}(b) = e^{-S_{pert}(Q^2, b_*) - S_{NP}(Q, b)} \sum_{i,j} C_{qi} \otimes f_{i/A}(x_1, \mu = C_0/b_*) C_{\bar{q}j} \otimes f_{j/B}(x_2, \mu = C_0/b_*). \quad (4)$$

Here σ_0 is the leading order cross section, Y is the regular piece in the limit $q_t \rightarrow 0$, the C 's contain the hard collinear virtual corrections, S_{pert} is the perturbative Sudakov factor, S_{NP} is the non-perturbative Sudakov factor, $x_{1,2} = \frac{Q}{\sqrt{s}} e^{\pm y}$ represent the momentum fractions

carried by the incoming partons in the given process, and $f_{i/A}, f_{j/B}$ are the PDFs. The non-perturbative Sudakov factor is introduced to handle the Landau pole in QCD in the limit that $b \rightarrow \infty$, or $\mu \rightarrow 0$. Therefore, we adopt the b_* formalism introduced in [28], and defined as:

$$b_* = \frac{b}{\sqrt{1 + b^2/b_{\max}^2}}, \quad b_{\max} < 1/\Lambda_{QCD} \quad (5)$$

The resummed and fixed-order calculations then have to be matched at intermediate P_T , where the fixed-order calculation does not contain large logarithms, in order to obtain QCD predictions for the entire range of P_T . Kinematic restrictions on the decay products of the W -boson can then be applied to mimic the cuts applied to the data.

C. **ResBos** versus parton shower

ResBos, the analytic resummation program proposed and discussed in detail in Refs. [29–31], is used to calculate at next-to-next-to-leading-log and next-to-next-to-leading order (NNLL) + (NNLO) the W boson cross section for the process $pp \rightarrow W^\pm + X \rightarrow l^\pm(\bar{\nu}) + X$, using a renormalization and factorization scale of $\mu_{R/F} = M_W$. **ResBos** uses a NNLO/NLO k-factor to obtain the NNLO correction for the Y -piece of the resummed cross-section. This k-factor is calculated as a function of Q, q_t , and y . The non-perturbative Sudakov form factor, which describes the P_T^W at low P_T , is fit using DY data, and parameterized by the BLNY form [31].

Parton shower Monte Carlo (MC) programs such as **POWHEG** [32] can also be used to simulate W -boson production at the LHC. Parton showers resum the leading tower of logarithms while resummation programs can include higher order logarithms, providing a more accurate description of the P_T^W distribution. The W -boson decay distributions are highly sensitive to the order of resummation in the calculation, and these small differences can potentially lead to an uncertainty in M_W greater than the desired goal of 10 MeV. Therefore, a calculation to the highest precision possible in terms of both higher loop and higher logarithmic resummation effects (better than available in Monte-Carlo Event Generators) would provide a better description of the W -boson kinematics and decay products, thus, allowing for a more precise determination of M_W [30].

The MC event generators output are fully exclusive, i.e. they provide information on all final-state particles, but approximate the resummation effects through the parton shower.

On the other hand, resummation programs are inclusive, that is they provide information only about the W -boson and its decay products, integrating out all additional QCD radiation. However, the information provided by resummation codes is sufficient for a precise analysis of the W -boson mass. Additionally, corrections to the width of W -boson and spin correlations between the initial and final state particles are included in `ResBos` when applicable.

III. TECHNIQUE FOR PDFS ERROR UNCERTAINTY DETERMINATION

A. PDF uncertainty estimation

In order to quantify the PDF uncertainty on M_W , we need to fit the detailed shapes of the P_T^l distributions using the log-likelihood (χ^2) analysis. The steps we have followed are outlined below, following Ref. [15]:

- 1 - While keeping M_W fixed at a given central value ($M_{W,0}=80358$ MeV), we generate the lepton distribution P_T^l for the 56 Hessian error PDFs. Here we choose the mass at 80358 MeV based on the indirect fit obtained in Ref. [10] (since we are only interested in the shift of the mass due to the PDF, and not what the central prediction is, this is an acceptable choice to make).
- 2 - We now leave the PDF fixed to the central value (CT14), and vary M_W in the generator to obtain the lepton transverse momentum distributions (P_T^l) for each of the different masses considered. Here we consider values between 80308 MeV and 80408 MeV, creating templates in steps of 1 MeV.
- 3 - For each template generated in step 1, corresponding to the different error PDFs for a fixed W boson mass (termed P_i^{PDF}), we loop through all the different templates generated in step 2, corresponding to the different W -boson masses using the central CT14 PDF (termed P_j^{Mass}), and compute the corresponding χ_{ij}^2 ;

$$\chi_{ij}^2 = \sum_{k=1}^{N_{\text{bins}}} \frac{(P_i^{\text{PDF}} - P_j^{\text{Mass}})^2}{(\sigma_i^{\text{PDF}})^2 + (\sigma_j^{\text{Mass}})^2}, \quad (6)$$

where the summation is over all of the lepton transverse momentum bins and σ is the statistical uncertainty for the given bin.

4 - For a fixed PDF j , the value of i that minimizes the χ_{ij}^2 distribution corresponds to the mass that would be predicted by the given PDF error set. In other words, the mass predicted by PDF j would be mass i ($M_{W,i}$), if the corresponding χ_{ij}^2 was the smallest.

The PDF Hessian uncertainty on M_W then is calculated using the master equation proposed in Ref [33, 34] for asymmetric uncertainties as follows:

$$\begin{aligned}\Delta M_W^+ &= \sqrt{\sum_{i=1}^{28} [\max(\{M_{W,i+} - M_{W,0}\}, \{M_{W,i-} - M_{W,0}\}, 0)]^2} \\ \Delta M_W^- &= \sqrt{\sum_{i=1}^{28} [\max(\{M_{W,0} - M_{W,i+}\}, \{M_{W,0} - M_{W,i-}\}, 0)]^2}\end{aligned}\tag{7}$$

where $M_{W,i}^\pm$ represents the best fit value for the M_W corresponding to the PDF set i induced by a change of ± 1 standard deviations of each independent parameter describing the PDF set.

Normalizing the P_T^l templates to use the shape of the distribution, instead of the overall rate, in the fit region can substantially reduce the size of the PDF induced uncertainties without losing the sensitivity to the value of M_W [24]. Thus, in this study we use normalized distributions.

B. Fit parameters and kinematic cuts

ResBos is used to calculate the W^\pm boson kinematics at NNLL+NNLO for the processes $pp \rightarrow W^\pm + X \rightarrow l^\pm \bar{\nu} + X$, based on the CT14 NNLO sets, at $\sqrt{s} = 7, 8$ and 13 TeV. A charm pole mass of 1.3 GeV is used, as in the CT10 NNLO PDFs. The PDFs for up, down, strange (anti) quarks and the gluon are parametrized at an initial scale of 1.295 GeV. The central PDF sets are obtained using a central value of $\alpha_s(M_Z)$ of 0.118 , as recommended by the PDF4LHC group [16]. The CT14 NNLO PDFs uncertainties are provided as 90% confidence level (CL) intervals, and then are scaled by a multiplicative factor of $(1/1.642)$ to provide 68% CL intervals. This scaling is appropriate if the χ^2 distribution is suitably quadratic.

The W -boson signal in the data is extracted by selecting events with one central isolated, high P_T^l lepton, large missing energy, and low hadronic recoil. In the **ResBos**-generated

events, the cuts described below are implemented, mimicking the cuts used in the experimental analyses. The missing transverse energy (\cancel{E}_T) is required to be greater than 20 GeV and the absolute value of the lepton pseudo-rapidity ($|\eta|$) is restricted to the region less than 2.5. For comparative purposes, different ranges for P_T^l (in GeV) are used: ($35 < P_T^l < 45$, $30 < P_T^l < 50$, and $20 < P_T^l < 60$). All of these ranges have equally spaced bins of 0.5 GeV. Additional boson transverse momentum cuts ($P_T^W < 15$ GeV, $P_T^W < 30$ GeV, $P_T^W < 60$ GeV and $P_T^W < 300$ GeV) are applied as well for comparison². We generate approximately 100M events for each template to minimize the effects of statistical fluctuations.

IV. RESULTS

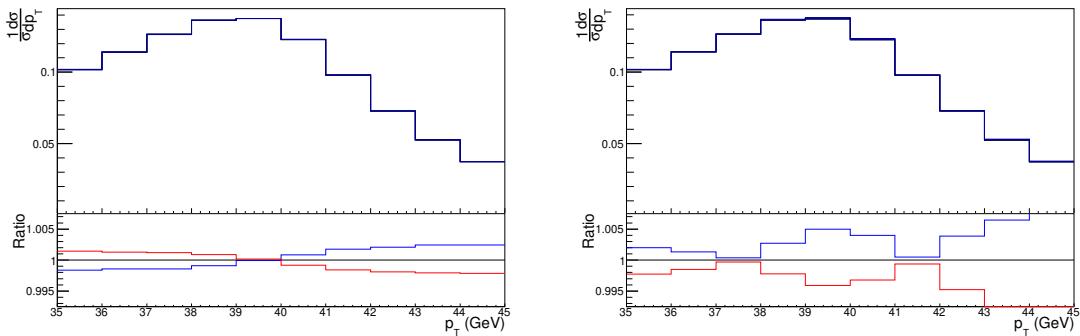


FIG. 1: P_T^l (in GeV) for the different P_j^{Mass} (left) and P_i^{PDF} (right) PDF errors for W^+ at $\sqrt{s} = 7$ TeV.

Figure 1 shows sample lepton transverse momentum distributions for the W -boson mass variations (left) and for the error PDF variations (right) for W^+ at $\sqrt{s} = 7$ TeV. Figure 2 shows the same lepton transverse momentum distributions for W^- . The variations, too small to easily observe in the top distributions, are more easily seen in the ratio plots underneath.

For the ratio plots to $M_{W,0}$ for the W -boson mass variations (left side of Figs. 1,2), the

² In the ResBos Monte Carlo generator, the W -boson transverse momentum is known exactly. It is difficult to reconstruct the transverse momentum experimentally with precision, but it is possible (see for example [35, 36]), and restrictions on the W -boson transverse momentum can lead to reduced values for W -boson mass uncertainty.

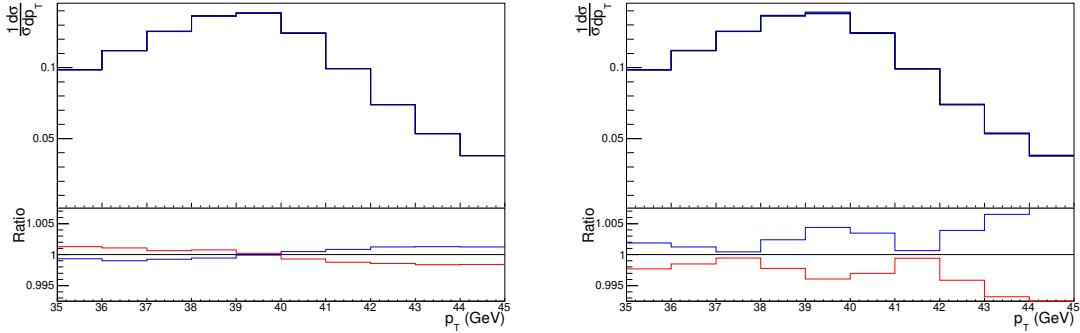


FIG. 2: P_T^l (in GeV) for the different P_j^{Mass} (left) and P_i^{PDF} (right) PDF errors for W^- at $\sqrt{s} = 7$ TeV.

red (blue) curve corresponds to the maximum (minimum) allowed mass from the χ^2 fit. The crossing of the two curves occurs due to the fact that the result is normalized to unity. Similarly, the ratio plots for the PDF variations (right side of Figs. 1,2), the upper (lower) curve corresponds to the total positive (negative) PDF uncertainty obtained from the master equation for the transverse momentum of the lepton. Here the curves do not cross due to the fact that the upper curve always corresponds to the positive direction uncertainty, while the lower corresponds to the negative direction uncertainty. For the variation of the error PDFs, the positive and negative ratios are approximately mirror images of each other, indicating that the uncertainty on the W mass from the PDFs is reasonably symmetric.

As an example, the $\Delta\chi_{ij}^2$ profiles as a function of the W^+ -boson mass at $\sqrt{s} = 7$ TeV are shown in Figure 3. The $\Delta\chi_{ij}^2$ profiles are reproduced for the central PDF, for the PDF that produces the highest mass W -boson, and the PDF that produces the lowest mass W -boson. Note that the $\Delta\chi_{ij}^2$ distributions are parabolic, indicating that the 90% errors can be scaled down to 68% by applying a scaling factor. Similar curves are obtained at center-of-mass energies of 8 and 13 TeV.

The Hessian PDF uncertainties, as a function of various kinematic cuts, for both W^\pm at $\sqrt{s} = 7$ TeV, are listed in Table I. Those for $\sqrt{s} = 8$ TeV are listed in Table II, and those for $\sqrt{s} = 13$ TeV are shown in Table III.

A number of generalizations can be made from this information. As expected, restricting the W -boson transverse momentum, P_T^W , to be < 15 GeV results in the smallest PDF uncertainties, as does restricting the P_T^l range to a narrow band ($35 < P_T^l$ (GeV) < 45) about

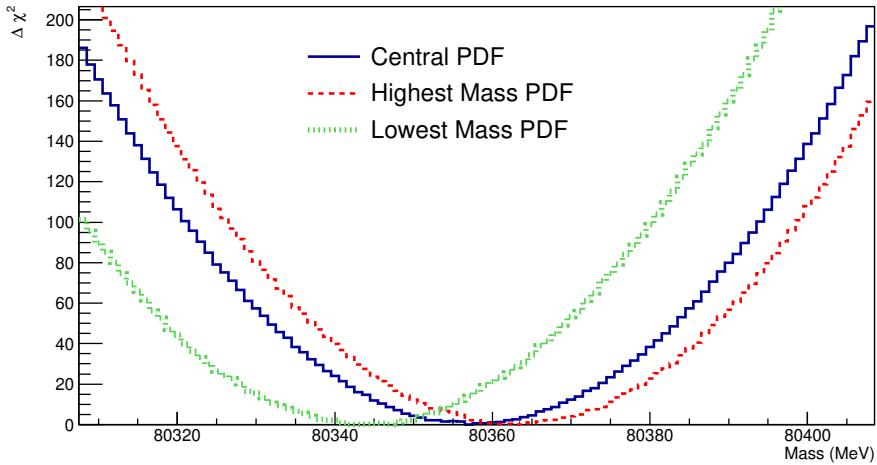


FIG. 3: The χ_{ij}^2 profiles as a function of M_W for W^+ at $\sqrt{s} = 7$ TeV. Here the highest mass PDF is that which gives the largest M_W , and similarly for the smallest. The PDFs used here are from the CT14 error set, representing a 90% C.L.

the Jacobian peak. The PDF uncertainty (and the $+/ -$ asymmetry for the uncertainty) tends to grow as the P_T^W cut increases. The PDF uncertainties are similar for W^+ and W^- production.

A previous analysis [15], using the techniques outlined here, estimated the PDF uncertainty for the W mass using a selection of PDF sets including CT10 NNLO. The analysis framework was based on the **POWHEG** MC interfaced with **PYTHIA** PS [37]. This analysis obtained a global uncertainty that ranged between ± 18 and ± 24 MeV, depending on the final state, collider energy and the PDF set. For CT10 NNLO, Ref. [15] reported that at $\sqrt{s} = 13$ TeV and with $P_T^W < 15$ GeV and $29 < P_T^l$ (GeV) < 49 , the errors are $^{+20}_{-17}$ MeV and $^{+17}_{-9}$ MeV for W^+ and W^- respectively. The results in this paper for a similar lepton transverse momentum range ($30 < P_T^l$ (GeV) < 50) are $^{+17.8}_{-18.3}$ MeV and $^{+15.7}_{-17.1}$ MeV respectively. At $\sqrt{s} = 8$ TeV, Ref. [15] reported errors of $^{+17}_{-18}$ MeV for W^+ and $^{+16}_{-11}$ MeV for W^- , similar to what we have observed.

$35 < P_T^l$ (GeV) < 45	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15$ GeV	+ 16.0 - 17.5	+ 15.9 - 15.6
$P_T^W < 30$ GeV	+ 18.1 - 22.4	+ 22.4 - 18.1
$P_T^W < 60$ GeV	+ 14.2 - 24.7	+ 21.1 - 20.5
$P_T^W < 300$ GeV	+ 12.5 - 32.3	+ 28.5 - 17.2
$30 < P_T^l$ (GeV) < 50	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15$ GeV	+ 17.4 - 18.7	+ 15.0 - 17.7
$P_T^W < 30$ GeV	+ 20.1 - 27.6	+ 23.0 - 24.6
$P_T^W < 60$ GeV	+ 21.1 - 27.8	+ 19.4 - 29.7
$P_T^W < 300$ GeV	+ 17.8 - 30.0	+ 19.2 - 29.3
$20 < P_T^l$ (GeV) < 60	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15$ GeV	+ 20.0 - 19.3	+ 20.7 - 14.1
$P_T^W < 30$ GeV	+ 23.1 - 31.3	+ 31.9 - 21.5
$P_T^W < 60$ GeV	+ 19.9 - 39.4	+ 14.1 - 42.8
$P_T^W < 300$ GeV	+ 17.5 - 43.8	+ 10.0 - 51.3

TABLE I: The impact of different cuts on the PDF induced uncertainty of M_W at $\sqrt{s} = 7$ TeV.

V. SUMMARY

The data samples for W -boson production at the LHC are large, for all center-of-mass energies. However, one of the limiting factors for determination of M_W is not statistics, but rather our imprecise knowledge of PDFs. In this study, we investigated this dominant source of uncertainty. Additionally, the imprecise knowledge of the p_T^W distribution of the W -boson adds additional uncertainty which is left to the experimentalists to accurately determine.

To improve the PDF uncertainty on the M_W determination, a better knowledge of the relevant PDFs is needed. Such improvement may be possible as new generations of PDFs (such as CT18) include more LHC data, such as precision measurements of the W and Z boson cross sections.

$35 < P_T^l(\text{GeV}) < 45$	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15 \text{ GeV}$	+ 13.9 - 14.8	+ 15.0 - 12.7
$P_T^W < 30 \text{ GeV}$	+ 21.7 - 16.3	+ 18.9 - 17.2
$P_T^W < 60 \text{ GeV}$	+ 21.5 - 21.7	+ 22.4 - 19.7
$P_T^W < 300 \text{ GeV}$	+ 20.0 - 25.2	+ 26.3 - 15.9
$30 < P_T^l(\text{GeV}) < 50$	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15 \text{ GeV}$	+ 16.1 - 16.1	+ 17.6 - 15.4
$P_T^W < 30 \text{ GeV}$	+ 23.5 - 20.2	+ 23.0 - 21.2
$P_T^W < 60 \text{ GeV}$	+ 24.5 - 21.8	+ 29.1 - 19.7
$P_T^W < 300 \text{ GeV}$	+ 21.8 - 23.5	+ 28.3 - 20.0
$20 < P_T^l(\text{GeV}) < 60$	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15 \text{ GeV}$	+ 13.6 - 21.9	+ 17.9 - 16.6
$P_T^W < 30 \text{ GeV}$	+ 20.9 - 32.0	+ 28.7 - 21.9
$P_T^W < 60 \text{ GeV}$	+ 25.1 - 34.5	+ 32.4 - 27.9
$P_T^W < 300 \text{ GeV}$	+ 34.2 - 26.8	+ 34.7 - 27.8

TABLE II: The impact of different cuts on the PDF induced uncertainty of M_W at $\sqrt{s} = 8 \text{ TeV}$.

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$35 < P_T^l \text{ (GeV)} < 45$	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15 \text{ GeV}$	+ 12.2 - 19.2	+ 14.8 - 15.3
$P_T^W < 30 \text{ GeV}$	+ 16.3 - 25.9	+ 14.7 - 24.8
$P_T^W < 60 \text{ GeV}$	+ 19.5 - 27.1	+ 20.7 - 24.3
$P_T^W < 300 \text{ GeV}$	+ 21.8 - 27.6	+ 16.4 - 27.3
$30 < P_T^l \text{ (GeV)} < 50$	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15 \text{ GeV}$	+ 17.8 - 18.3	+ 15.7 - 17.1
$P_T^W < 30 \text{ GeV}$	+ 24.2 - 24.6	+ 20.2 - 25.4
$P_T^W < 60 \text{ GeV}$	+ 24.0 - 25.8	+ 23.8 - 26.7
$P_T^W < 300 \text{ GeV}$	+ 20.4 - 29.3	+ 21.0 - 27.6
$20 < P_T^l \text{ (GeV)} < 60$	W^+ Uncertainty (MeV)	W^- Uncertainty (MeV)
$P_T^W < 15 \text{ GeV}$	+ 18.2 - 18.2	+ 16.9 - 17.2
$P_T^W < 30 \text{ GeV}$	+ 27.6 - 28.2	+ 23.6 - 27.1
$P_T^W < 60 \text{ GeV}$	+ 27.5 - 33.7	+ 30.4 - 30.6
$P_T^W < 300 \text{ GeV}$	+ 26.4 - 33.5	+ 24.6 - 30.3

TABLE III: The impact of different cuts on the PDF induced uncertainty of M_W at $\sqrt{s} = 13$ TeV.

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