

# New study of the neutron rich $^{136}\text{Te}$ isotope through decay spectroscopy

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**Abstract.**  $^{136}\text{Sb}$  is a neutron rich nucleus with one valence proton and three valence neutrons outside the doubly-magic nucleus  $^{132}\text{Sn}$ . It plays an important role in the rapid neutron capture process (r-process) as participates in its path further connected to the nucleosynthesis and the type II supernovae. In this work, a new  $\beta$  decay  $\gamma$  ray data on  $^{136}\text{Te}$  is obtained. The experiment is performed at the Institute Laue-Langevin using the Lohengrin spectrometer and the thermal neutron-induced fission reaction on  $^{235}\text{U}_{143}$ . A specific  $\beta$ - $\beta$  condition method is investigated for the  $\gamma$ -ray detection.

## 1 Introduction

The rapid neutron capture process (r-process) represents series of important nuclear reactions preformed at neutron-rich region in nuclear chart, responsible for creating about half of the elements heavier than iron. Nuclei around  $^{132}\text{Sn}$  play a significant role in r-process, their nuclear structure properties, such as mass ( $A$ ), half life ( $T_{1/2}$ ),  $\beta$  delayed neutron emissions  $P_n$  and  $P_{2n}$  values, are essential to reproduce the r-process abundances and help us to understand the origin of the elements in our universe. Tellurium ( $Z = 52$ ) is found at the second r-process peak ( $A \approx 130$ ) associated with the  $N = 82$  neutron shell closure in the solar system r-process distribution and predominantly produced in the main component of the r-process. The  $^{136}\text{Te}$  isotopes reveal structure dominated by pairs of valence protons and neutrons above the  $^{132}\text{Sn}$  core and nuclei in this neutron-rich region are mostly produced by fission reactions or the  $\beta$  decay of fission products [1]. The previous study about  $\beta$  decay of  $^{136}\text{Sb}$  is incomplete [2], the  $\gamma$  ray spectrum of  $^{136}\text{Te}$  from  $\beta$  decay of the mother  $^{136}\text{Sb}$  isotope is investigated here.

## 2 Experiment

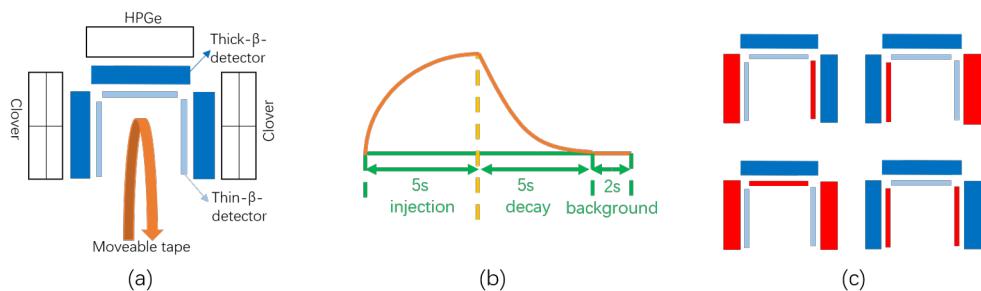
### 2.1 Setup

The experiment is performed at the Institute Laue-Langevin by the thermal neutron-induced fission reaction of  $^{235}\text{U}$ . The fission productions go through the LOHENGRIN separator and are selected by the magnetic and electric field according to their mass ( $A$ ) over ionic charge ( $q$ ) ratio and kinetic energy ( $E_k$ ) over ionic charge ( $q$ ). The target nuclei are implanted into a movable tape which surrounded by  $\beta$  detectors and  $\gamma$  detectors. The main configuration

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of the detection system is presented in Fig.1-(a). The  $\beta$  signals are recorded by two plastic detectors of different thickness: thin and thick, placed in front of the detection area of the  $\gamma$  ray detectors, consisting by two clovers and one coaxial detector. The clover detector is made of four high purity germanium crystals, which allows to improve the statistic of high energy  $\gamma$  rays by add-back procedure. To prevent the affect from the long lived contaminants, time chopper signal is used to manage the measurement cycle. A movable tape is used to take away the injected nuclei after each measurement cycle. According to the lifetime of the  $^{136}\text{Sb}$  [3], each measurement cycle is divided into three parts: the first 5 s for beam injection, the next 5 s we stop the injection and wait the implanted nuclei to decay and the last 2 s tape is moved to take away the rest decay products (background), see Fig.1-(b).



**Figure 1.** The main configuration of the detection system (a). The time chopper information (b). The schematic drawing of  $\beta$  detector combinations used to trigger the  $\gamma$ -ray spectra (c).

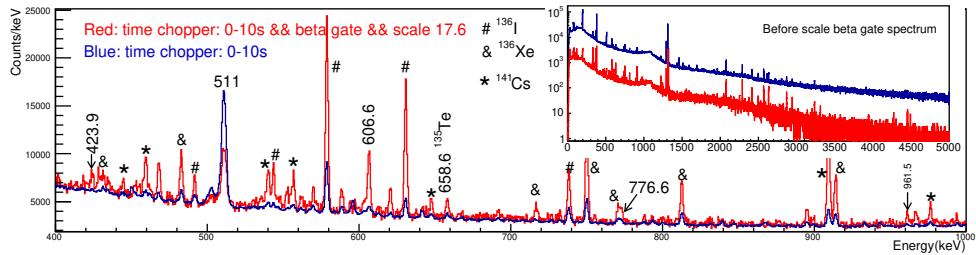
## 2.2 Data analysis

In order to obtain the  $\gamma$  ray spectrum of  $^{136}\text{Te}$  after  $\beta$  decay of  $^{136}\text{Sb}$ , we used the time chopper information between 0 s and 10 s. The plastic detectors serve as a  $\beta$  trigger only while their  $\beta$  energy information was not used. Each  $\gamma$  ray is triggered by the combinations of thin and/or the opposite thin and/or thick plastic detector as shown in Fig.1-(c). The schematically drawn combinations represent the used thick/thin  $\beta$ - $\beta$  coincidence condition. For example, when a set of thick/thin plastic detectors registers a  $\beta$  signal within the same time window (marked in red), this is considered as a  $\beta$ - $\beta$  coincidence [4] and the  $\gamma$  rays are respectively registered by any of the HpGe detectors. This kind of  $\beta$ - $\beta$  coincidence condition between two or three  $\beta$  detectors strongly suppresses the background without reducing the statistics of transitions of interest. On the contrary, one can use an anti-coincidence with the  $\beta$  signals to artificially enhance background and exploit background lines.

## 3 Result

Fig.2 shows part of the energy spectrum after  $\beta$  decay of  $^{136}\text{Sb}$  (in blue) obtained in the 0-10 s of each measurement cycle. Adding the  $\beta$ - $\beta$  coincidence condition by all combinations represents the  $\beta$ -gated spectrum (in red). One can see from the inset, an order of magnitude suppression of background using this condition. In order to find the effect of the  $\beta$  trigger, the  $\beta$ -gated spectrum is compared to the normalised unconditioned  $\gamma$  spectrum. The  $\gamma$  rays from  $^{136}\text{Sb}$   $\beta$  decay to  $^{136}\text{Te}$  and  $\beta$ -n decay to  $^{135}\text{Te}$  are marked with their energy values. For example, the 606.6 keV and 658.6 keV transitions correspond to the de-excitation of the first

excited states in  $^{136}\text{Te}$  [2] and  $^{135}\text{Te}$  [5], respectively. It is obvious that the application of this  $\beta$ - $\beta$  coincidence procedure is very significant in the case of weak transitions. To demonstrate this, one can point to such weak transitions in  $^{136}\text{Te}$  as the 423.9 keV and 961.5 keV lines, which are not visible before the coincidence condition, while they can be easily identified after the  $\beta$ - $\beta$  coincidence trigger. In the Fig.2 we can also see the main contaminants coming from the grand-daughter nuclei  $^{136}\text{I}$  and  $^{136}\text{Xe}$ . The  $^{141}\text{Cs}$  is also detected by our system.



**Figure 2.** A  $\beta$ -gated (red)  $\gamma$  ray spectrum for  $^{136}\text{Te}$  is compared with its un-gated  $\gamma$  ray spectrum (in blue). The transitions in  $^{136}\text{Te}$  are marked with their energy values.

The efficiency of the  $\beta$  detectors could be measured by comparing the intensity of  $\gamma$  rays in the Ge spectra with and without condition on  $\beta$  detectors. These measurements could be executed by implanting beam with certain mass number ( $A$ ) e.g.  $A = 137$  into the tape system so, that  $^{137}\text{Te}$  and  $^{137}\text{I}$  decay quickly to create a pure  $^{137}\text{Xe}$  source. Based on the strongest 445 keV ( $^{137}\text{Cs}$ ) [6] and 333 keV ( $^{136}\text{I}$ ) [7] transitions, the results of the efficiency measurements amount to 47(5)% for this experiment. Note that similar value is obtained for the energy range up to about 1.4 MeV.

## 4 Summary

A specific  $\beta$ - $\beta$  coincidence trigger is used for the study of  $\beta$  delayed  $\gamma$  ray spectrum. By applying this method to  $^{136}\text{Sb}$   $\beta$  decay experiment, one can conform weak transitions from previous work [2] and identify new transitions. The  $\beta$  efficiency is found to be 47(5)% for  $Q_\beta$  value of around 5 MeV. Note that high energy  $\beta$  rays may also be deposited in the HPGe detectors. Such signals would be added to  $\gamma$  ray detection and may cause an energy shift for these  $\gamma$  rays. As next step one may design a new detector system to reject these kind of synthetic signals as well as to improve the  $\beta$  detection efficiency.

## References

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