

2.3 The K_L^0 Beam Facility at JLab

Moskov Amaryan

Department of Physics

Old Dominion University

Norfolk, VA 23529, U.S.A.

Abstract

Following a Letter of Intent submitted to PAC43 at JLab in this talk we discuss the possibility to create a secondary K_L^0 beam in Hall D to be used with GlueX detector for spectroscopy of excited hyperons.

1. Introduction

Our current understanding of strong interactions is embedded in Quantum Chromodynamics (QCD). However, QCD being a basic theory, extremely successful in explaining the plethora of experimental data in the perturbative regime, faces significant challenges to describe the properties of hadrons in non-perturbative regime. Constituent Quark Model (CQM) is surprisingly successful in explaining spectra of hadrons, especially in the ground state; however, CQM appears to be too naive to describe properties of excited states. It is natural that excited states are not simply explained with spatial excitations of constituent quarks, but it is an effective representation revealing complicated interactions of quarks and gluons inside. Hadron spectroscopy aims to provide a comprehensive description of hadron structure based on quark and gluon degrees of freedom. Despite many successes in observing hundreds of meson and baryon states experimentally we haven't succeeded to either observe or rule out existence of glueballs, hybrids and multi quark systems; although it is tempting to explain recently observed X, Y, Z [1] states as first evidences of tetraquarks as well as recently observed heavy baryon states at LHCb [2] as charmed pentaquarks.

An extensive experimental program is developed to search for hybrids in the GlueX experiment at JLab. Over the last decade, significant progress in our understanding of baryons made of light (u, d) quarks have been made in CLAS at JLab. However, systematic studies of excited hyperons are very much lacking with only decades old very scarce data filling the world database in many channels. In this experiment we propose to fill this gap and study spectra of excited hyperons using the modern CEBAF facility with the aim to use proposed secondary K_L^0 beam with physics target of the GlueX experiment in Hall D. The goal is to study $K_L - p$ and $K_L - d$ interactions and do the baryon spectroscopy for the strange baryon sector.

Unlike in the cases with pion or photon beams, Kaon beams are crucial to provide the data needed to identify and characterize the properties of hyperon resonances.

Our current experimental knowledge of strange resonances is far worse than our knowledge of N and Δ resonances; however, within the quark model, they are no less fundamental. Clearly there is a need to learn about baryon resonances in the "strange sector" to have a complete understanding of three-quark bound states.

The masses and widths of the lowest mass baryons were determined with Kaon-beam experiments in the 1970s [1]. First determination of pole positions, for instance for $\Lambda(1520)$, were

obtained only recently from analysis of Hall A measurement at JLab [3]. An intense Kaon beam would open a window of opportunity not only to locate missing resonances, but also to establish properties including decay channels systematically for higher excited states.

A comprehensive review of physics opportunities with meson beams is presented in a recent paper [4]. Importance of baryon spectroscopy in strangeness sector was discussed in Ref. [5].

2. Reactions that Could be Studied with K_L^0 Beam

More details about this chapter could be found in a talk by Mark Manley at this workshop.

(a) Elastic and charge-exchange reactions



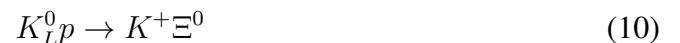
(b) Two-body reactions producing $S = -1$ hyperons



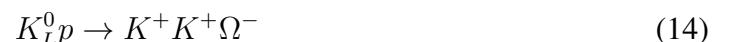
(c) Three-body reactions producing $S = -1$ hyperons



(d) Two- and three-body reactions producing $S = -2$ hyperons



(e) Three-body reactions producing $S = -3$ hyperons



Reactions 10–15 will be discussed in more details below.

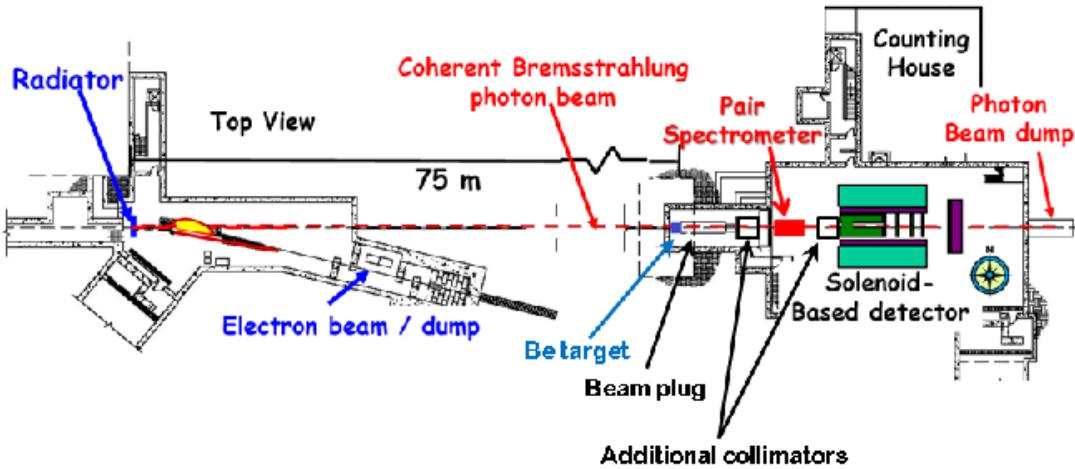


Figure 1: Schematic view of Hall D beamline. See a text for explanation.

3. The K_L^0 Beam in Hall D

In this Section, we describe photo-production of secondary K_L^0 beam in Hall D. There are few points that need to be decided. To produce intensive photon beam one needs to increase radiation length of the radiator up to 10% radiation length. In a first scenario $E_e = 12$ GeV, electrons produced at CEBAF will scatter in a radiator in the tagger vault, generating intensive beam of bremsstrahlung photons. This may will then require removal of all tagger counters and electronics and very careful design of radiation shielding, which is very hard to optimize and design. In a second scenario one may use Compact Photon Source design (for more details see a talk by Pavel Degtarenko at this workshop) installed after the tagger magnet, which will produce bremsstrahlung photons and dump electron beam inside the source shielding the radiation inside. At the second stage, bremsstrahlung photons interact with Be target placed on a distance 16 m upstream of liquid hydrogen (LH_2) target of GlueX experiment in Hall D producing K_L^0 beam. To stop photons a 30 radiation length lead absorber will be installed in the beamline followed by a sweeping magnet to deflect the flow of charged particles. The flux of K_L on LH_2 target of GlueX experiment in Hall D will be measured with pair spectrometer upstream the target. Details of this part of the beamline (for a details see a talk by Ilya Larin at this workshop). Momenta of K_L particles will be measured using the time-of-flight between RF signal of CEBAF and start counters surrounding LH_2 target. Schematic view of beamline is presented in Fig. 1. The bremsstrahlung photons, created by electrons at a distance about 75 m upstream, hit the Be target and produce K_L^0 mesons along

with neutrons and charged particles. The lead absorber of ~ 30 radiation length is installed to absorb photons exiting Be target. The sweeping magnet deflects any remaining charged particles (leptons or hadrons) remaining after the absorber. The pair spectrometer will monitor the flux of K_L^0 through the decay rate of Kaons at given distance about 10 m from Be target. The beam flux could also be monitored by installing nuclear foil in front of pair spectrometer to measure a rate of K_S^0 due to regeneration process $K_L + p \rightarrow K_S + p$ as it was done at NINA (for details see a talk by Michael Albrow at this workshop).

Here, we outline experimental conditions and simulated flux of K_L^0 based on GEANT4 and known cross sections of underlying subprocesses [6–8].

- An electron beam with energy $E_e = 12$ GeV and current $I_e = 5 \mu A$ (maximum possible, limited by the Hall D beam dump).
- A thickness of radiator 5 % radiation length.
- Primary Be target with $R = 4$ cm, $L = 40$ cm.
- LH_2 target with $R = 2$ cm, $L = 30$ cm.
- Distance between Be and LH_2 targets 24 m.

The expected flux of K_L^0 mesons integrated in the range of momenta $P = 0.3 - 10$ GeV/c will be $\approx 2 \times 10^3 K_L^0/\text{sec}$ on the physics target of the GlueX setup.

In a more aggressive scenario with

- A thickness of radiator 10%.
- Be target with a length $L = 60$ cm.
- LH_2 target with $R = 3$ cm.

The expected flux of K_L^0 mesons integrated over the same momentum range will increase to $\approx 10^4 K_L^0/\text{sec}$.

In addition to these requirements it will require lower repetition rate of electron beam with ~ 40 ns spacing between bunches to have enough time to measure time-of-flight of the beam momenta and to avoid an overlap of events produced from alternating pulses. Lower repetition rate was already successfully used by G0 experiment in Hall C at JLab [9].

The radiation length of the radiator needs further studies in order to estimate the level of radiation and required shielding in the tagger region. During this experiment all photon beam tagging detector systems and electronics will be removed.

The final flux of K_L^0 is presented with 10% radiator, corresponding to maximal rate .

In the production of a beam of neutral Kaons, an important factor is the rate of neutrons as background. As it is well known, the ratio $R = N_n/N_{K_L^0}$ is on the order 10^3 from primary proton beams [10], the same ratio with primary electromagnetic interactions is much lower. This is illustrated in Fig. 2, which presents the rate of Kaons and neutrons as a function of the momentum, which resembles similar behavior as it was measured at SLAC [11].

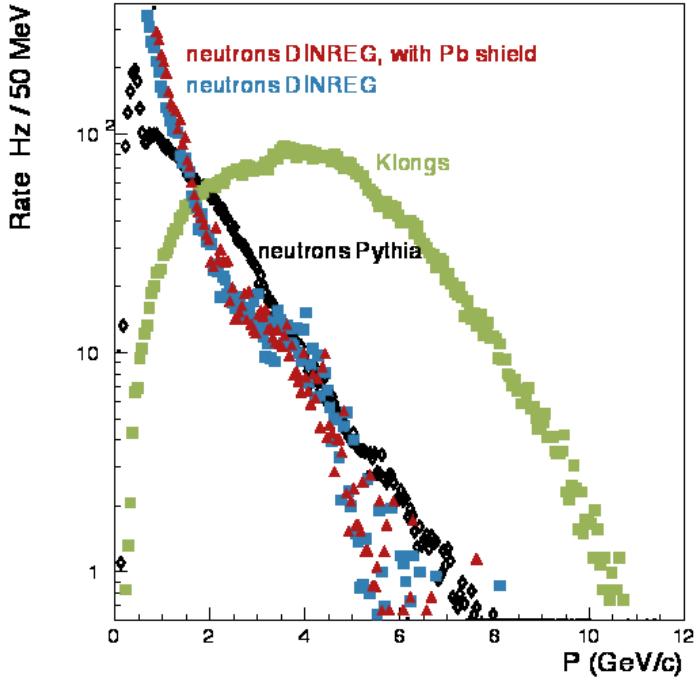


Figure 2: The rate of neutrons (open symbols) and K_L^0 (full squares) on LH_2 target of Hall D as a function of their momenta simulated with different MC generators with $10^4 K_L^0/\text{sec}$.

Shielding of the low energy neutrons in the collimator cave and flux of neutrons has been estimated to be affordable, however detailed simulations are under way to show the level of radiation along the beamline.

The response of GlueX setup, reconstruction efficiency and resolution are presented in a talk by Simon Taylor at this workshop.

4. Expected Rates

In this Section, we discuss expected rates of events for some selected reactions. The production of Ξ hyperons has been measured only with charged Kaons with very low statistical precision and never with primary K_L^0 beam. In Fig. 3 panel a) shows existing data for the octet ground state Ξ 's with theoretical model predictions for W (the reaction center of mass energy) distribution, panel b) shows the same model prediction [12] presented with expected experimental points and statistical error for 10 days of running with our proposed setup with a beam intensity $2 \times 10^3 K_L^0/\text{sec}$ using missing mass of K^+ in the reaction $K_L^0 + p \rightarrow K^+ \Xi^0$ without detection of any of decay products of Ξ^0 (for more details on this topic see a talk by Kanzo Nakayama at this workshop).

The physics of excited hyperons is not well explored, remaining essentially at the pioneering stages of '70s-'80s. This is especially true for $\Xi^*(S = -2)$ and $\Omega^*(S = -3)$ hyperons. For example, the $SU(3)$ flavor symmetry allows as many $S = -2$ baryon resonances, as

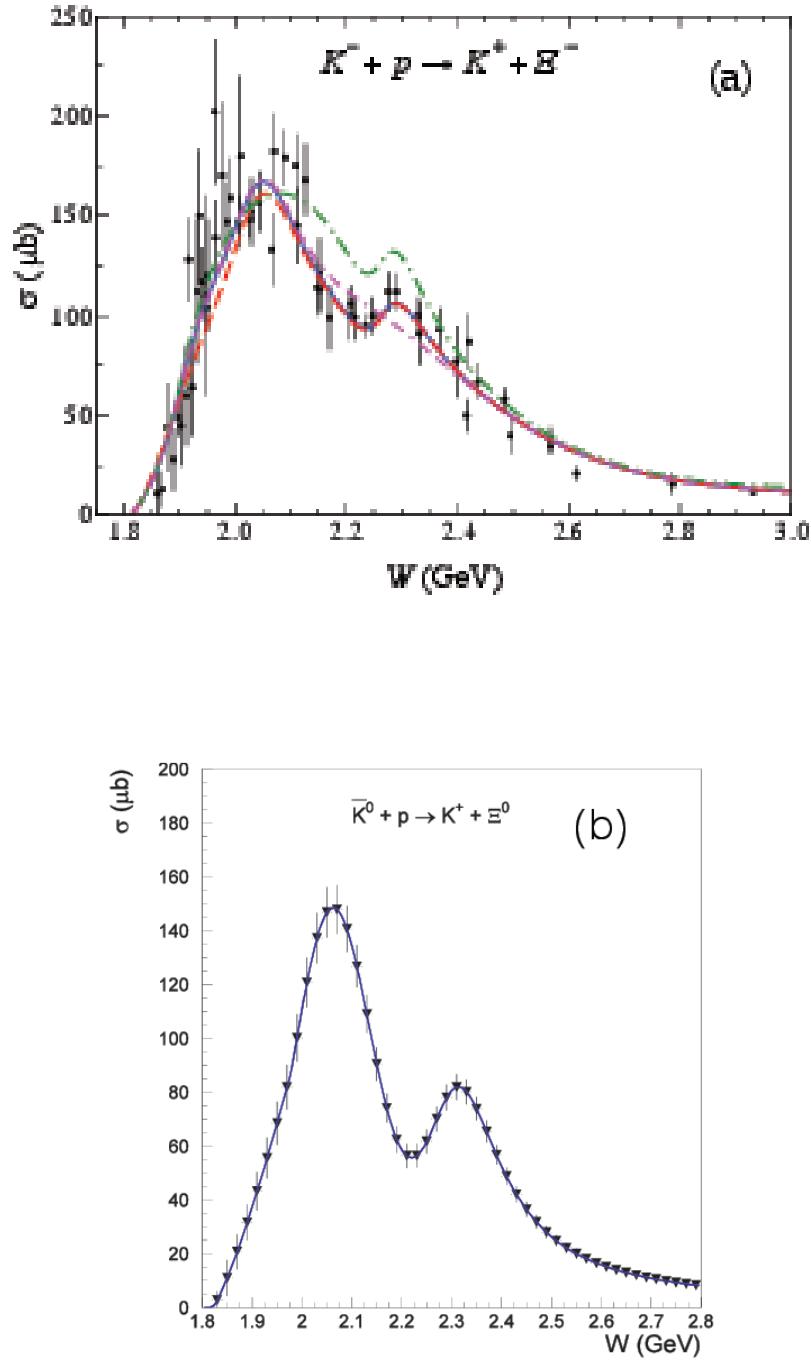


Figure 3: a) Cross section for existing world data on $K^- + p \rightarrow K^+ \Xi^-$ reaction with model predictions from [12]; b) expected statistical precision for the reaction $\bar{K}^0 + p \rightarrow K^+ \Xi^0$ in 10 days of running with a beam intensity $2 \times 10^3 K_L/\text{sec}$ overlaid on theoretical prediction [12].

there are N and Δ resonances combined (≈ 27); however, until now only three [ground state $\Xi(1382)1/2^+$, $\Xi(1538)3/2^+$, and $\Xi(1820)3/2^-$] have their quantum numbers assigned and

few more states have been observed [1]. The status of Ξ baryons is summarized in a table presented in Fig. 4 together with quark model predicted states [13].

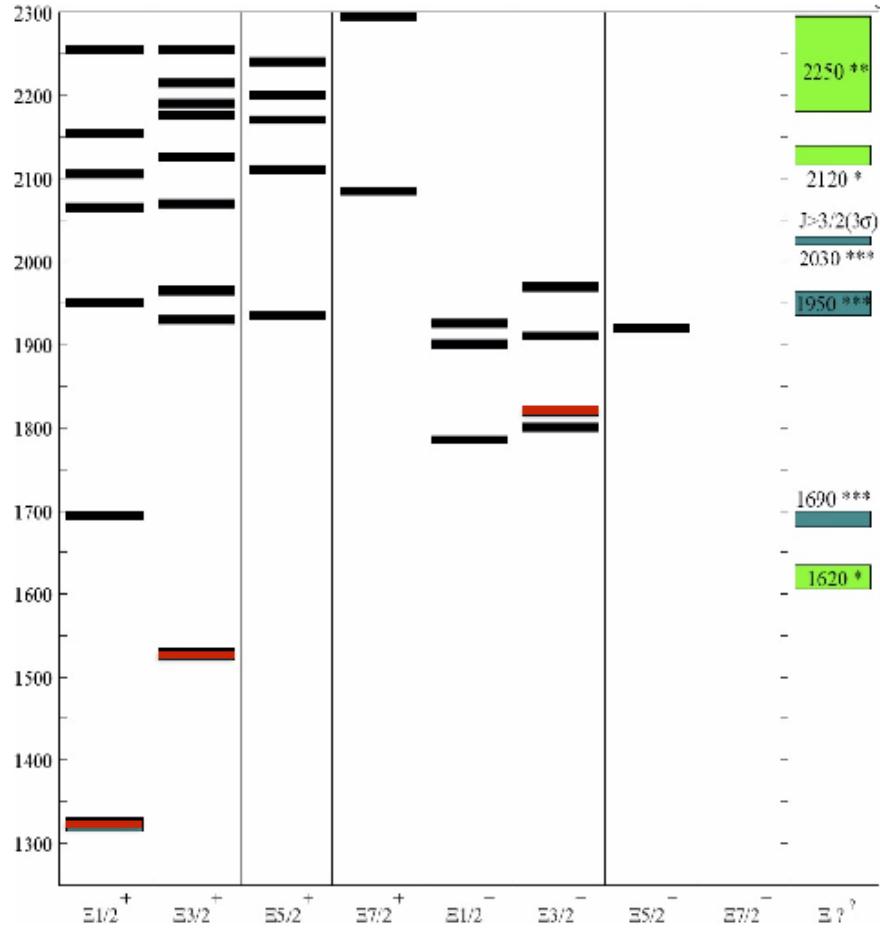


Figure 4: Black bars: Predicted Ξ spectrum based on the quark model calculation [13]. Colored bars: Observed states. The two ground octet and decuplet states together with $\Xi(1820)$ in the column $J^P = 3/2^-$ are shown in red color. Other observed states with unidentified spin-parity are plotted in the rightmost column.

Historically the Ξ^* states were intensively searched for mainly in bubble chamber experiments using the K^-p reaction in '60s-'70s. The cross section was estimated to be on the order of 1–10 μb at the beam momenta up to 10 GeV/c . In '80s-'90s, the mass or width of ground and some of excited states were measured with a spectrometer in the CERN hyperon beam experiment. Few experiments have studied cascade baryons with the missing mass technique. In 1983, the production of Ξ^* resonances up to 2.5 GeV were reported from $p(K^-, K^+)$ reaction from the measurement of the missing mass of K^+ [14]. The ex-

perimental situation with Ω^{-*} 's is even worse than the Ξ^* case, there are very few data for excited states. The main reason for such a scarce dataset in multi strange hyperon domain is mainly due to very low cross section in indirect production with pion or, in particular, - photon beams. Currently only ground state Ω^- quantum numbers are identified. Recently significant progress is made in lattice QCD calculations of excited baryon states [15, 16] which poses a challenge to experiments to map out all predicted states (for more details see a talk by David Richards at this workshop). The advantage of baryons containing one or more strange quarks for lattice calculations is that then number of open decay channels is in general smaller than for baryons comprising only the light u and d quarks. Moreover, lattice calculations show that there are many states with strong gluonic content in positive parity sector for all baryons. The reason why hybrid baryons have not attracted the same attention as hybrid mesons is mainly due to the fact that they lack manifest "exotic" character. Although it is difficult to distinguish hybrid baryon states, there is significant theoretical insight to be gained from studying spectra of excited baryons, particularly in a framework that can simultaneously calculate properties of hybrid mesons. Therefore this program will be very much complementary to the GlueX physics program of hybrid mesons.

The proposed experiment with a beam intensity $10^4 K_L/\text{sec}$ will result in about $2 \times 10^5 \Xi^*$'s and $4 \times 10^3 \Omega^*$'s per month.

A similar program for KN scattering is under development at J-PARC with charged Kaon beams [17]. The current maximum momentum of secondary beamline of $2 \text{ GeV}/c$ is available at the K1.8 beamline. The beam momentum of $2 \text{ GeV}/c$ corresponds to $\sqrt{s}=2.2 \text{ GeV}$ in the K^-p reaction which is not enough to generate even the first excited Ξ^* state predicted in the quark model. However, there are plans to create high energy beamline in the momentum range $5 - 15 \text{ GeV}/c$ to be used with the spectrometer commonly used with the J-PARC $E50$ experiment which will lead to expected yield of $(3 - 4) \times 10^5 \Xi^*$'s and $10^3 \Omega^*$'s per month.

Statistical power of proposed experiment with K_L beam at JLab will be of the same order as that in J-PARC with charged Kaon beam.

An experimental program with Kaon beams will be much richer and allow to perform a complete experiment using polarized target and measuring recoil polarization of hyperons. This studies are under way to find an optimal solution for GlueX setup.

5. Summary

In summary, we intend to create high intensity K_L beam using photoproduction processes from a secondary Be target. A flux as high as $10^4 K_L/\text{sec}$ could be achieved. Momenta of K_L beam particles will be measured with time of flight. The flux of Kaon beam will be measured through partial detection of $\pi^+\pi^-$ decay products from their decay to $\pi^+\pi^-\pi^0$ by exploiting similar procedure used by LASS experiment at SLAC [11]. Besides using unpolarized LH_2 target currently installed in GlueX experiment additional studies are needed to find the optimal choice of polarized targets. This proposal will allow to measure KN scattering with different final states including production of strange and multi strange baryons with unprecedented statistical precision to test QCD in non perturbative domain. It has a potential to distinguish between different quark models and test lattice QCD predictions for excited baryon states with strong hybrid content.

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References

- [1] K.A. Olive *et al.* (Particle Data Group), Chin. Phys. C **38**, 090001 (2014).
- [2] R. Aaij *et al.* (LHCb Collaboration), Phys. Rev. Lett. **115**, 072001 (2015).
- [3] Y. Qiang, Ya.I. Azimov, I.I. Strakovsky, W.J. Briscoe, H. Gao, D.W. Higinbotham, and V.V. Nelyubin, Phys. Lett. B **694**, 123 (2011).
- [4] W.J. Briscoe, M. Döring, H. Haberzettl, D.M. Manley, M. Naruki, I.I. Strakovsky, and E.S. Swanson, Eur. Phys. J. A **51**, 129 (2015).
- [5] U. Al-Binni *et al.*, *Project X: Physics Opportunities*, in *Community Summer Study 2013: Snowmass on the Mississippi* (CSS2013), Minneapolis, MN, USA, July–August, 2013, edited by A.S. Kronfeld and R.S. Tschirhart.
- [6] H. Seraydaryan *et al.* (CLAS Collaboration), Phys. Rev. C **89**, 055206 (2014).
- [7] A.I. Titov and T.-S.H. Lee, Phys. Rev. C **67**, 065205 (2003).
- [8] G. Mclellan *et al.*, Phys. Rev. Lett. **26**, 1593 (1971).
- [9] D. Androic *et al.* (G0 Collaboration), Nucl. Instrum. Meth. A **646**, 59 (2011).
- [10] W.E. Cleland *et al.*, Phys. Rev. D **12**, 1247 (1975).
- [11] G.W. Brandenburg *et al.*, Phys. Rev. D **7**, 708 (1973).
- [12] B.C. Jackson, Y. Oh, H. Haberzettl, and K. Nakayama, Phys. Rev. C **91**, 065208 (2015).
- [13] K.-T. Chao, N. Isgur, and G. Karl, Phys. Rev. D **23**, 155 (1981).
- [14] C.M. Jenkins *et al.*, Phys. Rev. Lett. **51**, 951 (1983).
- [15] R.G. Edwards, N. Mathur, D.G. Richards, and S.J. Wallace (Hadron Spectrum Collaboration), Phys. Rev. D **87**, 054506 (2013).
- [16] G.P. Engel, C.B. Lang, D. Mohler, and A. Schfer (BGR Collaboration), Phys. Rev. D **87**, 074504 (2013).
- [17] H. Takahashi, Proceedings of the *11th International Conference on Hypernuclear and Strange Particle Physics* (HYP 2012), Nucl. Phys. A **914**, 553 (2013).