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Study of Isothermal Compressibility and Speed of Sound in the Hadronic Matter Formed in Heavy-Ion Collision Using Unified Formalism

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Abstract: The thermodynamical quantities and response functions are useful to describe the particle production in heavy-ion collisions as they reveal crucial information about the produced system. While the study of isothermal compressibility provides an inference about the viscosity of the medium, speed of sound helps in understanding the equation of state. With an aim towards understanding the system produced in the heavy-ion collision, we have made an attempt to study isothermal compressibility and speed of sound as function of charged particle multiplicity in heavy-ion collisions at $\sqrt{s_{NN}} = 2.76$ TeV, 5.02 TeV, and 5.44 TeV using unified formalism.

Keywords: relativistic heavy ion collision; QGP; unified statistical framework; isothermal compressibility; speed of sound



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1. Introduction

Among the main aim of heavy-ion collision program at present collider experiments such as Relativistic Heavy Ion Collider (RHIC) and Large Hadron Collider (LHC) is to mimic the state that was created few microseconds after the Big Bang. This state of matter, created at extremely high temperature and energy density, is called the Quark-Gluon Plasma (QGP) and is also believed to be present at the core of massive neutron stars. Interaction between quark and gluon, which leads to the formation of QGP, is governed by the Quantum Chromo-dynamics (QCD). Asymptotic freedom, which is an important pillar of QCD, suggests a confinement deconfinement phase transition, during which the hadronic degree of freedom changes to partonic degree of freedom. Whether the phase transition is first-order, second-order or a simple cross-over and the search for the critical point are some of the important questions that are of immediate interest in the particle physics community.

Since the formation of QGP occurs at a very short time scale, it is not possible to directly probe in the experiment using current technologies. Therefore, we rely on the information carried by the final state particles to the detectors to gain insight into the medium created in the heavy-ion collision. Although we only measure the kinematic quantities such as the pseudorapidity η , transverse momentum p_T , energy E etc. of the final state particles in the experiment, a breadth of information about the medium can be extracted by studying these kinematic observables.

Another set of quantities that are not directly observable, however, play an important role in understanding the nature of the medium and the equation of state are the thermodynamical response functions. This includes quantities that express how a system responds to change in some external parameters such as pressure, temperature etc. Isothermal compressibility (κ_T), and speed of sound (c_s) are some of the response function that are of

interest in high-energy physics [1–3]. The isothermal compressibility, κ_T , which exhibits the important property of the medium, tells us how much the volume of the medium changes on the change in pressure at a fixed temperature. This quantity can be used to study how close a medium is to be called perfect fluid. Perfect fluids are ideal fluids that do not possess shear stress, viscosity and also do not conduct heat. The κ_T of perfect fluid is zero and the zero value signifies that the fluid is incompressible. Although the incompressible fluids do not exist in nature, the recent findings of the value of κ_T , as in Ref. [4], are almost close to zero which suggests that the medium created is almost a perfect fluid. Estimation of κ_T will tell us about how close the medium to be a perfect fluid which will lead us an approximation about the viscosity as viscosity describes a fluid's resistance to flow, hence perfect fluid has zero viscosity. Isothermal compressibility and viscosity of a liquid is related to the low-shear viscosity of a liquid dispersion of solid particles [5]. Perfect fluid can also be characterized by the ratio of shear viscosity to entropy density (η/s). Calculation based on *AdS/CFT* correspondence has put up a universal lower bound of $1/4\pi$ for strongly interacting quantum field theories [6]. On the other hand, the value of η/s has been found to be close to the lower bound based on the flow harmonics calculation of the experimental data, indicating the near-perfect behaviour of medium created in heavy-ion collision [7,8].

As explained in Ref. [9], the speed of sound can quantify the nature of the same as it connects and explains the hydrodynamical evolution of the produced matter in the heavy-ion collisions. Fundamentally the speed of sound also gives the information about the equation of state, which relates pressure (P) and the energy density (ϵ). For a non-interacting massless ideal gas, the value of the squared speed of sound c_s^2 is expected to be $1/3$ times speed of light squared [10]. Hence, the comparison with the massless ideal gas will give crucial information about the system dynamics and reveals the nature of the medium [11]. Different studies suggested that for the system created in heavy-ion collision the value of c_s^2 is close to the ideal value [1,3,11–17].

As already discussed, these quantities are not directly observable in the experiment and we extract them by utilizing the distribution of kinematic observables such as the transverse momentum p_T -spectra, rapidity, angle of emission etc. The p_T -spectra carries sufficient information to study such quantities as it is directly related to the energy of the system. Understanding the distribution of p_T is in itself a tedious task because in the low- p_T region, the QCD coupling strength is very high and hence we cannot apply the perturbative QCD theories to explain the spectra. Several phenomenological models have been developed to tackle this issue, and the most widely accepted are the statistical thermal models. We can utilize the statistical thermal models to extract the thermodynamical quantities such as temperature, number density, energy density etc.

If we assume the purely thermal origin of final state particles, the most natural choice to explain the energy distribution of particles is Boltzmann-Gibbs (BG) statistics [18–20]. However, it has been discussed in many works [21,22] that the BG distribution function deviates significantly from the experimental data because the spectra are more like a power-law rather than the simple exponential. Also, the BG statistics fails to explain the strongly correlated systems [23] in which the long-range correlations are present, and entropy becomes non-additive and non-extensive [24]. The existence of long-range interaction in high-energy heavy-ion collisions is discussed in Ref. [25] motivating to explore beyond the extensive BG regime to study the spectra. In 1988, C. Tsallis proposed a statistics [26–28], introducing an additional parameter q , which takes care of the non-extensivity in the system. It is a thermodynamically consistent [29,30], generalized version of Boltzmann distribution [31]. The power-law behavior of Tsallis distribution makes it a good choice to study the p_T -spectra and it is shown to nicely fit the spectra, particularly in the low- p_T region. Although Tsallis statistics nicely explain the data in the low- p_T region, however, it starts to deviate from the experimental data as we move toward the high- p_T part of the spectra.

Particle spectra in heavy-ion collisions can be divided into two distinct regions, low- p_T regime corresponds to the particle produced in soft processes whereas the hard processes

dominate particle production in the high- p_T region. The limitation of Tsallis statistics in explaining the particle produced in hard processes demands a framework that can consider the effect of both soft and hard processes in the particle spectra. Some modification in Tsallis statistics [32–35] has been proposed to explain the high- p_T part of spectra in the heavy-ion collision, however, more work is required in this direction to get the full benefit from the spectra. To explain both the hard and soft part of particle spectra in a consistent manner, a unified theory using Pearson distribution is introduced in Ref. [21]. It is a generalized form of the Tsallis distribution and is shown to be thermodynamically consistent and backward compatible to the Tsallis statistics within some limit on its parameters [22].

In this work, we have calculated the isothermal compressibility and speed of sound for charged hadrons produced in heavy-ion collisions using the unified statistical framework. For this analysis, we have taken the experimental data of transverse momentum spectra for charged hadrons produced in $Pb - Pb$ collision at $\sqrt{s_{NN}} = 2.76$ TeV [36], 5.02 TeV [37], and $Xe - Xe$ collision at 5.44 TeV [38] measured by the ALICE experiment.

2. Methodology

The basic thermodynamic quantities that are of interest to formulate the isothermal compressibility and speed of sound include energy density ϵ , number density n and pressure P . From the standard thermodynamics, the number of particles N in a system and its total energy E can be calculated as:

$$N = \sum_i f_i \quad (1)$$

$$E = \sum_i E_i f_i \quad (2)$$

where E_i is the energy of i^{th} state and f_i is the corresponding distribution function. The standard replacement while going from summation to integration for small energy intervals is given as [29]:

$$\sum_i \rightarrow V \int \frac{d^3 p}{(2\pi)^3} \quad (3)$$

Here, V is the volume and p represent the momentum. So, using the above transformation, the number density n will be of the form:

$$n = \int \frac{d^3 p}{(2\pi)^3} \times f(E) \quad (4)$$

and the corresponding energy density ϵ will be given as:

$$\epsilon = \int \frac{d^3 p}{(2\pi)^3} E \times f(E) \quad (5)$$

Since the momentum distribution of the final state particles are fixed at kinetic freeze-out [12], the pressure of the system could be estimated from the moments of energy distribution. The pressure P is given as:

$$P = \int \frac{d^3 p}{(2\pi)^3} \frac{p^2}{3E} \times f(E) \quad (6)$$

Among all the quantities discussed above, one common factor is the energy distribution of the particles ($f(E)$). Energy is related to the transverse mass m_T and pseudorapidity y as $E = m_T \cosh(y)$ and the transverse mass is defined in term of transverse momentum p_T and mass of particle m as $m_T = \sqrt{p_T^2 + m^2}$. So, the distribution of transverse momenta acts as a proxy for the energy distribution. Hence, the proper parameterization of transverse

momentum spectra is crucial to understand the thermodynamics of the system created in high-energy collisions. In the present work, we have used the unified statistical framework to explain the p_T -spectra and extract the thermodynamical quantities such as temperature T , non-extensive parameter q .

In the seminal work [39], Karl Pearson discussed a family of the curve, based on the first four moments (mean, variance, skewness, and kurtosis), called Pearson distribution. Before the introduction of Pearson formalism in 1895, all probability distribution was only constructed based on mean and variance and did not take care of skewness and kurtosis. Pearson introduced a new probability distribution function where skewness and kurtosis can also be adjusted along with the mean and variance of a distribution. An important characteristic of this distribution is that depending on the limit on its parameters, it reduces to different distribution function such as Gaussian, normal, Student's T, Gamma distribution etc. The differential form of a Pearson distribution function, $p(x)$, for a variable x is expressed as [40]:

$$\frac{1}{p(x)} \frac{dp(x)}{dx} + \frac{a+x}{b_0 + b_1 x + b_2 x^2} = 0 \quad (7)$$

where a , b_0 , b_1 , and b_2 are related to first four moments of the distribution. By integrating this differential equation, one can get,

$$p(x) = \exp \left(- \int \frac{x+a}{b_2 x^2 + b_1 x + b_0} dx \right) \quad (8)$$

Solving above equation we get the general solution of the form:

$$p(x) = C(e+x)^f(g+x)^h \quad (9)$$

$$p(x) = B \left(1 + \frac{x}{e} \right)^f \left(1 + \frac{x}{g} \right)^h \quad (10)$$

upto some normalization constant $B = Ce^f g^h$. Here, C , e , f , g & h are the parameters of the equation and can be related to the physical parameters such as temperature T , non-extensivity parameter q etc.

Distribution function, in case of unified statistical framework, obtained from the above Pearson distribution, is given as [22]:

$$f_i = (B f_E)^{1/q} f_{Ta} \quad (11)$$

where

$$B = \frac{C}{(p_0)^n} \left(\frac{T}{q-1} \right)^{\frac{-q}{q-1}} \quad (12)$$

$$f_E = \frac{1}{E} \left(1 + \frac{E}{p_0} \right)^{-n} \quad (13)$$

and

$$f_{Ta} = \left[1 + (q-1) \frac{p_T}{T} \right]^{\frac{-1}{q-1}} \quad (14)$$

In the above equation, p_0 and n are the free parameters with parameter n is related to the second order flow coefficient [21]. This formalism reduces to Tsallis statistics within the limit $n = -1$ and $p_0 = 0$. Therefore, it can be considered as a generalized version of the Tsallis function and explains both soft and hard process contributions to p_T -spectra.

The equation for the average number of particles and energy, in the case of unified formalism, remains the same as Tsallis [22]:

$$N = \sum_i f_i^q \quad (15)$$

and, the energy of the system will be:

$$E = \sum_i E_i f_i^q \quad (16)$$

Here, the additional power of q comes from the thermodynamic consistency. In case of the unified formalism, the transverse momentum spectra is defined as:

$$\frac{1}{2\pi p_T} \frac{d^2 N}{dp_T dy} = B' \left(1 + \frac{p_T}{p_0}\right)^{-n} \left[1 + (q-1) \frac{p_T}{T}\right]^{\frac{-q}{q-1}} \quad (17)$$

where $B' = B \times \frac{V}{(2\pi)^3}$, T is temperature and q is non-extensive parameter. Here we considered the chemical potential to be zero because at LHC energy, the net-baryonic number is extremely small at the central rapidity region. Thermodynamic parameters such as T , q and the other quantities can be obtained by fitting the measured transverse momentum spectra with the unified distribution using the Equation (17). These quantities extracted from the spectra can be used to calculate the response function as discussed below.

2.1. Isothermal Compressibility

In the high-energy collider experiment, we only consider a part of phase space because of the limited η acceptance of the detectors. The overall number of particles in the collision is conserved, but number of particles and energy in a particular phase space window may vary. Hence, the system can be considered as a grand canonical ensemble for the estimation of isothermal compressibility. Therefore, the variance of number of particles N , can be written as [41]:

$$\langle (N - \langle N \rangle)^2 \rangle = VT \frac{\partial n}{\partial \mu} \quad (18)$$

And, the isothermal compressibility, κ_T , can be written as:

$$\kappa_T = -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_T \quad (19)$$

Using the expression of variance of N and κ_T [41,42], we can write:

$$\langle (N - \langle N \rangle)^2 \rangle = \text{var}(N) = TVn^2\kappa_T \quad (20)$$

Equation (20) requires an event-by-event information of N to estimation κ_T . On contrary, we can compare Equations (18) and (20) to derive a fluctuation independent formula for isothermal compressibility as:

$$\kappa_T = \frac{\partial n / \partial \mu}{n^2} \quad (21)$$

As per Equation (21), the estimation of κ_T does not depend on the fluctuation of N and it makes the estimation possible without having event-by-event information of particle number.

Number density, n , in case of unified formalism, is of the form:

$$n = \int \frac{d^3 p}{(2\pi)^3} \times \frac{B}{E} \left(1 + \frac{E}{p_0}\right)^{-n} \left[1 + (q-1) \frac{(E-\mu)}{T}\right]^{\frac{-q}{q-1}} \quad (22)$$

and,

$$\frac{\partial n}{\partial \mu} = \int \frac{d^3 p}{(2\pi)^3} \times \frac{q}{T} \times \frac{B}{E} \left(1 + \frac{E}{p_0}\right)^{-n} \left[1 + (q-1) \frac{(E-\mu)}{T}\right]^{\frac{1-2q}{q-1}} \quad (23)$$

By using the above equations, we have estimated the values of κ_T for heavy-ion collisions at different energies.

2.2. Speed of Sound

For a thermodynamic system at temperature T and volume V , the squared speed of sound is given by,

$$c_s^2 = \left(\frac{\partial P}{\partial \epsilon}\right)_s \quad (24)$$

where P is pressure and ϵ is energy density of the system. As discussed in Ref. [43], the propagation of sound wave in a medium is an adiabatic process and entropy is constant in such process, hence the squared speed of sound is estimated at constant entropy density (s). Above equation can be further reduced to:

$$c_s^2 = \frac{\frac{\partial P}{\partial T}}{\frac{\partial \epsilon}{\partial T}} \quad (25)$$

where

$$P = \int \frac{d^3 p}{(2\pi)^3} \times B \times \frac{p^2}{3E^2} \left(1 + \frac{E}{p_0}\right)^{-n} \left[1 + (q-1) \frac{E}{T}\right]^{\frac{-q}{q-1}} \quad (26)$$

and,

$$\epsilon = \int \frac{d^3 p}{(2\pi)^3} \times B \left(1 + \frac{E}{p_0}\right)^{-n} \left[1 + (q-1) \frac{E}{T}\right]^{\frac{-q}{q-1}} \quad (27)$$

By using the above equations, the squared speed of sound c_s^2 reduces to

$$c_s^2 = \frac{\int \frac{p^2 d^3 p}{3E^2} \left(1 + \frac{E}{p_0}\right)^{-n} \left[\frac{T}{q-1} + E\right]^{\frac{1-2q}{q-1}}}{\int d^3 p \left(1 + \frac{E}{p_0}\right)^{-n} \left[\frac{T}{q-1} + E\right]^{\frac{1-2q}{q-1}}} \quad (28)$$

We have used the Equation (28) to estimate the squared speed of sound in the medium created in heavy-ion collision at three different energies.

3. Results and Discussion

This study presents a formalism to calculate κ_T and c_s^2 using the non-extensive unified statistical framework discussed in Ref. [22]. We have estimated the κ_T/V and c_s^2 in the medium formed of charged hadrons as a function of charged particle multiplicity for different collision systems. The data for charged particle multiplicity ($\langle \frac{dN_{ch}}{d\eta} \rangle$) corresponding to a particular centrality is taken from the experimental results Refs. [44–46].

For this analysis, we have considered the transverse momentum spectra of charged hadrons produced in $Pb - Pb$ collision at 2.76 [36] and 5.02 TeV [37] and $Xe - Xe$ collision at 5.44 TeV [38]. The p_T range is restricted to $p_T < 5$ GeV/c since we are trying to study bulk properties and the majority of high p_T particles are produced from hard processes. The experimental data for the p_T -spectra for all the energies used in the paper belongs to the pseudorapidity range $|\eta| < 0.8$. The unified function fit to the p_T -spectra at 2.76, 5.02 and 5.44 TeV is provided in the Refs. [21,47,48]. The numerical value of T , q and the other fitting parameters are calculated by fitting the measured transverse momentum spectra with the unified formalism as in Equation (17) and the best fit value of the parameters are provided in Table 1.

Table 1. Numerical values of the fit parameters T (GeV), q , p_0 (GeV/c) and n obtained by fitting the experimental data of p_T -spectra fitted with the unified formalism Equation (17).

Centrality	2.76 TeV				5.02 TeV				5.44 TeV			
	T	q	p_0	n	T	q	p_0	n	T	q	p_0	n
0–5%	0.393 ±0.05	1.048 ±0.004	0.105 ±0.21	0.749 ±0.36	0.407 ±0.003	1.048 ±0.000	0.0018 ±0.030	0.562 ±0.05	-	-	-	-
5–10%	0.386 ±0.04	1.053 ±0.041	0.0877 ±0.191	0.700 ±0.32	0.415 ±0.004	1.049 ±0.000	0.0167 ±0.033	0.604 ±0.05	-	-	-	-
10–20%	0.370 ±0.07	1.060 ±0.006	0.0600 ±0.18	0.619 ±0.30	0.422 ±0.004	1.052 ±0.000	0.0394 ±0.033	0.659 ±0.06	0.409 ±0.01	1.072 ±0.001	0.0977 ±0.08	0.720 ±0.15
20–30%	0.351 ±0.08	1.070 ±0.008	0.0385 ±0.18	0.548 ±0.30	0.424 ±0.012	1.059 ±0.001	0.0812 ±0.042	0.744 ±0.07	0.460 ±0.03	1.067 ±0.003	0.225 ±0.11	1.101 ±0.24
30–40%	0.331 ±0.07	1.081 ±0.008	0.0256 ±0.20	0.489 ±0.34	0.412 ±0.013	1.068 ±0.001	0.0824 ±0.038	0.749 ±0.07	0.447 ±0.04	1.079 ±0.004	0.2286 ±0.10	1.112 ±0.24
40–50%	0.311 ±0.08	1.093 ±0.008	0.0341 ±0.25	0.474 ±0.46	0.369 ±0.018	1.085 ±0.002	0.05 ±0.042	0.614 ±0.08	0.455 ±0.05	1.091 ±0.005	0.2881 ±0.13	1.306 ±0.36
50–60%	0.292 ±0.08	1.106 ±0.008	0.0457 ±0.32	0.468 ±0.61	0.34 ±0.023	1.101 ±0.002	0.0527 ±0.051	0.578 ±0.11	0.434 ±0.08	1.108 ±0.008	0.2904 ±0.15	1.317 ±0.51
60–70%	0.273 ±0.11	1.121 ±0.012	0.0747 ±0.49	0.487 ±1.03	0.311 ±0.025	1.118 ±0.002	0.0658 ±0.071	0.557 ±0.17	0.357 ±0.07	1.123 ±0.006	0.1977 ±0.17	0.943 ±0.53
70–80%	-	-	-	-	0.329 ±0.034	1.131 ±0.003	0.1565 ±0.094	0.855 ±0.29	0.338 ±0.09	1.139 ±0.011	0.2060 ±0.26	0.974 ±0.87

In Figure 1, we have plotted the isothermal compressibility over volume calculated using the Equations (21)–(23). It is observed that there is a decline in the values of κ_T/V with an increase in the multiplicity. At higher charged-particle multiplicity, κ_T/V becomes the lowest, which suggests that the system move toward near-ideal behaviour with the increase in multiplicity. This trend is in line with the expectation as higher multiplicity class contains a larger number of particles and hence a higher pressure is required to attain a small change in volume. Similar values of κ_T/V for different collision systems show an indication of similar dynamics of the produced medium. It is worth mentioning here that the ideal fluid is incompressible, hence $\kappa_T = 0$, implying that the volume cannot be changed by applying pressure. For water, the corresponding value is several order of magnitude higher than what is obtained in case of heavy-ion collision. The values for κ_T/V obtained in the case of heavy-ion collision using the unified formalism is in the range from 10^{-3} to 10^{-5} GeV $^{-1}$.

A proper estimation of volume is required to extract the value of κ_T (fm 3 /GeV). Different techniques have been developed and tested on diverse datasets to extract the volume parameter [32,49–55]. Although the numerical values vary greatly in different models, all of them are in the order of 10^3 – 10^4 fm 3 and hence utilizing the value of volume from these models will give us the value of κ_T in the order of 1 – 10 fm 3 /GeV. This range of the value of κ_T matches well with the values obtained by other techniques in the Ref. [4,56]. The obtained value of κ_T is very low as compared to the water and other materials, indicating that the compressibility of the system created in the heavy-ion collision is very close to an ideal fluid. Proper estimation of volume is still an undergoing field of research, hence, we did not select a particular model and instead, we presented the value in terms of κ_T/V .

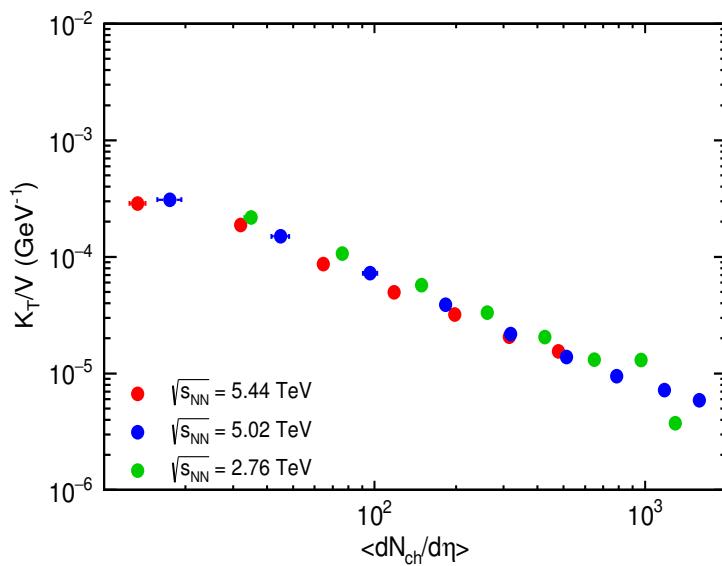


Figure 1. (Color online) Variation of isothermal compressibility over volume (κ_T/V) with the average charged particle multiplicity ($\langle \frac{dN_{ch}}{d\eta} \rangle$) for $Pb - Pb$ collision at $\sqrt{s_{NN}} = 2.76$ TeV, $Pb - Pb$ collision at $\sqrt{s_{NN}} = 5.02$ TeV and $Xe - Xe$ collision at $\sqrt{s_{NN}} = 5.44$ TeV using Unified formalism Equations (21)–(23).

We have also attempted to study the speed of sound for different collision systems in order to explore the properties of matter. The speed of sound in a medium reveals the properties of the medium via the equation of state. In Figure 2, we have plotted the squared speed of sound with charged-particle multiplicity for three different energies estimated using the Equation (28). It is observed that the value of the squared speed of sound is very close to $1/3$ times the speed of light squared, and there is an increase in the value with increasing $\langle \frac{dN_{ch}}{d\eta} \rangle$, suggesting that the system becomes more ideal at larger multiplicity. This observation complements the near-ideal behaviour already indicated from the measurement of isothermal compressibility.

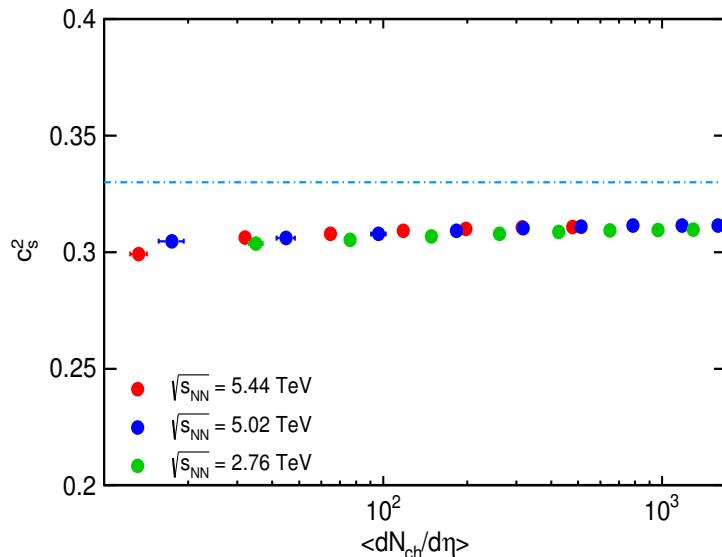


Figure 2. (Color online) Variation of squared speed of sound (c_s^2) as a function of $\langle \frac{dN_{ch}}{d\eta} \rangle$ for $Pb - Pb$ collision at $\sqrt{s_{NN}} = 2.76$ TeV, $Pb - Pb$ collision at $\sqrt{s_{NN}} = 5.02$ TeV and $Xe - Xe$ collision at $\sqrt{s_{NN}} = 5.44$ TeV using Unified formalism as Equation (28). The dotted line represents the theoretical value for ideal gas system.

4. Conclusions

With an aim towards understanding the system produced in the heavy-ion collision, we have made an attempt to study some thermodynamic response functions such as isothermal compressibility and speed of sound. Since transverse momentum spectra carries information about the system, we have analyzed spectra of charged hadrons at three different LHC energies using the unified formalism and used the extracted value of thermodynamical parameters to study the isothermal compressibility and speed of sound. The p_T -spectra of charged hadrons produced in $Pb - Pb$ collision at 2.76 TeV, 5.02 TeV and $Xe - Xe$ collision at 5.44 TeV are taken with p_T range upto 5 GeV/c . We have estimated the value of κ_T/V and c_s^2 and studied their variation as a function of charged particle multiplicity. We observed that while the value of κ_T/V decreases with respect to increase in multiplicity, the values of c_s^2 approaches to 1/3.

These estimations of κ_T/V and c_s^2 using unified formalism represent that the medium tends to move toward a near-ideal behavior with an increase in charged particle multiplicity. In conclusion, we have presented the theoretical formalism to study some of the thermodynamical response functions within the unified statistical framework discussed in the Ref. [22]. The extracted values point toward the creation of a near-ideal medium in high-energy collision and the system approach the ideal behavior as we move from peripheral to the central collision.

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