

# First Extraction of Polarized Sea Asymmetry from Weak Boson Production in Proton-Proton Collisions

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We present a global QCD analysis of spin-dependent parton distribution functions (PDFs) that includes the latest polarized  $W$ -lepton production data from the STAR collaboration at RHIC. These data allow the first data-driven extraction of a nonzero polarized light quark sea asymmetry  $\Delta\bar{u} - \Delta\bar{d}$  within a global QCD framework with minimal theoretical assumptions.

**KEYWORDS:** Global QCD analysis, antiquark PDFs

## 1. Introduction

Understanding the detailed decomposition of the proton spin into its constituent quark and gluon helicity and orbital angular momentum components promises to be one of the most significant accomplishments in nuclear and particle physics of this generation [1–3]. While the total light quark contributions to the helicity are well determined from polarized inclusive deep-inelastic scattering (DIS) data, and jet production in polarized  $pp$  collisions provides constraints on the gluon helicity, far less is known about the polarization of the antiquark sea. There have been some intriguing hints of a polarized antiquark asymmetry,  $\Delta\bar{u} - \Delta\bar{d}$ , from polarized semi-inclusive DIS (SIDIS) measurements, in analogy with the spin-averaged  $\bar{u} - \bar{d}$  asymmetry inferred from unpolarized DIS and Drell-Yan measurements.

Recently, more probes of antiquark polarization have been possible through  $W$ -lepton production in polarized  $pp$  collisions. In particular, the STAR collaboration at RHIC [4] has used polarized  $pp$  collisions at center of mass energy  $\sqrt{s} = 510$  GeV to measure the longitudinal single-spin asymmetry  $A_L = (\sigma_+ - \sigma_-)/(\sigma_+ + \sigma_-)$ , where  $\sigma_+$  ( $\sigma_-$ ) is the cross section for positive (negative) proton helicity, for the leptonic decay channels  $W^+ \rightarrow e^+ \nu$  and  $W^- \rightarrow e^- \bar{\nu}$ . At leading order, these can be written as

$$A_L^{W^+} \propto \frac{\Delta\bar{d}(x_1)u(x_2) - \Delta u(x_1)\bar{d}(x_2)}{\bar{d}(x_1)u(x_2) + u(x_1)\bar{d}(x_2)}, \quad (1a)$$

$$A_L^{W^-} \propto \frac{\Delta\bar{u}(x_1)d(x_2) - \Delta d(x_1)\bar{u}(x_2)}{\bar{u}(x_1)d(x_2) + d(x_1)\bar{u}(x_2)}, \quad (1b)$$

where  $\Delta f$  ( $f$ ) represents a polarized (unpolarized) PDF evaluated at momentum fraction  $x_1$  ( $x_2$ ) carried by the parton in the polarized (unpolarized) proton. Combined with the DIS observables, these asymmetries provide a vital new handle on the extraction of the polarized antiquark distributions  $\Delta\bar{u}$  and  $\Delta\bar{d}$ .

Previous global analyses [5, 6] have sought to extract the asymmetry under various assumptions and with different methods for estimating uncertainties. De Florian *et al.* (DSSV) [5] extracted a positive  $\Delta\bar{u} - \Delta\bar{d}$  asymmetry from spin-dependent data with fixed input for unpolarized PDFs and fragmentation functions (FFs), assuming PDF positivity and SU(3) symmetry for axial-vector charges

within errors. The impact of the latter assumptions was examined in a simultaneous analysis of spin PDFs and FFs by the JAM collaboration [7], who found polarized light antiquark and strange PDFs consistent with zero when the constraints were relaxed. The Monte Carlo analysis by the NNPDF collaboration [6] generated prior samples from the DSSV fit [5], thus inheriting the corresponding biases. The NNPDF analysis also used a reweighting procedure involving  $\chi^2$ -based weights, which is inconsistent with the Gaussian likelihood used in the generation of the replicas [8].

Instead of relying on reweighting prescriptions and assumptions about PDF positivity or SU(3) flavor symmetry, here we present a new simultaneous global QCD analysis of unpolarized and polarized PDFs and FFs, including for the first time STAR  $A_L^W$  data, along with data on inclusive and semi-inclusive polarized lepton-nucleon DIS and jet production in polarized  $pp$  collisions. The Monte Carlo analysis allows us to more reliably quantify the uncertainties on all distributions, and examine the interplay between the sea asymmetry and parametrizations of FFs.

## 2. Theoretical Framework

Our theoretical framework is based on fixed-order collinear factorization for high-energy scattering processes. The cross section for the  $W$ -lepton production process can be written as differential in the lepton pseudorapidity,  $\eta_\ell$ , and its transverse momentum,  $p_T^\ell$ . The renormalization and factorization scales are chosen to be the mass of the  $W$  boson,  $\mu_R = \mu_F = M_W$ , and the NLO expressions for the hard scattering kernels are found in Ref. [9].

The scale dependence of the PDFs is determined according to the DGLAP evolution equations, with the PDFs and  $\alpha_s$  evolved at next-to-leading logarithmic accuracy with the boundary condition  $\alpha_s(M_Z) = 0.118$ . For light as well as heavy quarks the PDFs are evolved using the zero-mass variable flavor number scheme. The values of the heavy quark mass thresholds for the evolution are taken from the PDG as  $m_c = 1.28$  GeV and  $m_b = 4.18$  GeV in the  $\overline{\text{MS}}$  scheme [10].

Our PDF extraction procedure is based on Bayesian inference using the Monte Carlo techniques developed in previous JAM analyses [7]. The parameterization of the unpolarized PDFs is discussed in Ref. [11], while for the polarized PDFs and FFs at the input scale  $\mu_0 = m_c$  we use the form

$$f(x, \mu_0) = Nx^\alpha(1-x)^\beta(1+\eta x), \quad (2)$$

where  $N$ ,  $\alpha$ ,  $\beta$ , and  $\eta$  are fit parameters. The polarized light quark PDFs  $\Delta u$  and  $\Delta d$  are parameterized as a sum of a valence and a sea component. For the sea quark  $\Delta \bar{u}$ ,  $\Delta \bar{d}$ ,  $\Delta s$ , and  $\Delta \bar{s}$  PDFs we use two functions of the form (2), one of which is unique to each flavor while the other describes the low- $x$  region and is shared between all four distributions. The same template (2) is used for FFs, but with  $x$  replaced by the momentum fraction  $z$  of the parton carried by the hadron, and with  $\eta = 0$ . More details about the FF parametrization, and a discussion on how it can influence the results on the polarized sea asymmetry, can be found in [12].

Recently the question of PDF positivity beyond leading order in  $\alpha_s$  in the  $\overline{\text{MS}}$  scheme has been debated [13, 14]. Such a constraint would require  $|\Delta f(x, Q^2)| \leq f(x, Q^2)$  to hold for all flavors at all  $x$  and  $Q^2$ . To explore this question phenomenologically, we perform analyses with and without the positivity constraints. The baseline analysis, referred to in the following as “JAM”, does not enforce positivity; however, when included, the positivity constraints are enforced approximately on each Monte Carlo replica by imposing a penalty on the  $\chi^2$  function when the bounds are violated.

## 3. Quality of Fit

Our analysis includes measurements of the DIS asymmetries  $A_{||}$  and  $A_1$  for the proton, deuteron, and  $^3\text{He}$ , with the restrictions on the four-momentum squared  $Q^2 > m_c^2$  and on the hadronic final state masses  $W^2 > 10 \text{ GeV}^2$  to ensure that the asymmetries are dominated by the leading-twist  $g_1$  structure

function. With the same cuts we include pion and kaon SIDIS measurements on polarized proton and deuteron targets, with the fragmentation variable restricted to  $0.2 < z < 0.8$  to ensure the applicability of the leading-power formalism and avoid threshold corrections.

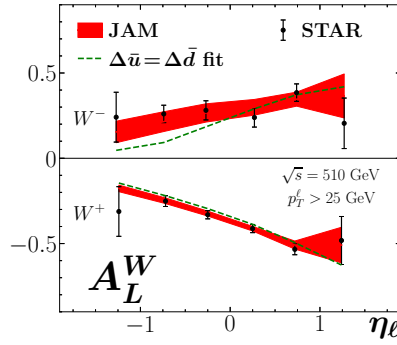


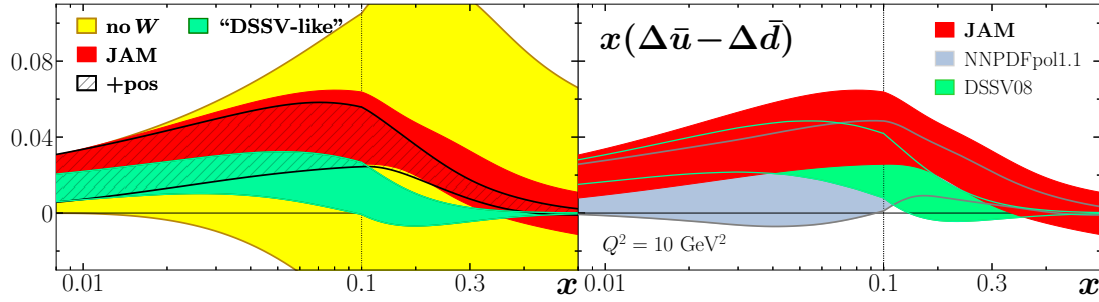
Fig. 1 Single-spin asymmetries  $A_L^W$  versus pseudorapidity  $\eta_\ell$  from STAR [4] (black circles) at  $\sqrt{s} = 510$  GeV and integrated over  $p_T^\ell > 25$  GeV, compared with the full JAM fit (red solid lines and  $1\sigma$  uncertainty bands) and with a fit where  $\Delta\bar{u}$  is set equal to  $\Delta\bar{d}$  (green dashed lines).

Beyond polarized lepton scattering, we describe jet production data in polarized  $pp$  collisions with a cut on the jet transverse momentum of 8 GeV. We also include for the first time single-spin asymmetry  $A_L^W$  data from STAR [4] and  $A_L^{W/Z}$  from PHENIX, which provide the most direct constraints on the antiquark polarization. For unpolarized processes, we use data from inclusive DIS, Drell-Yan lepton-pair production, and inclusive  $W^\pm$ ,  $Z$  and jet production in hadronic collisions, as in Ref. [15]. The FFs are constrained mainly by semi-inclusive pion and kaon production data in  $e^+e^-$  scattering.

The quality of our global analysis is summarized in Table I, which shows a global  $\chi^2/N_{\text{dat}} = 1.05$  for  $N_{\text{dat}} = 5495$  data points (587 for polarized, 4464 for unpolarized, and 444 for SIA). The  $\chi^2/N_{\text{dat}}$  for each experiment is generally stable whether PDF positivity constraints are imposed or not. When enforcing  $\Delta\bar{u} = \Delta\bar{d}$ , there are significant increases in  $\chi^2/N_{\text{dat}}$  for the STAR  $W$  data (from 0.53 to 2.51), PHENIX  $W/Z$  data at mid rapidity (0.18 to 1.49), and for the COMPASS  $A_{1p}^{\pi^-}$  data (0.72 to 1.40, as observed in Ref. [7]). The STAR  $A_L^W$  measurement is compared with the JAM fit in Fig. 1 versus the pseudorapidity  $\eta_\ell$ . When the asymmetry is forced to vanish, the quality of the fit suffers the most for  $A_L^{W^-}$  at low  $\eta_\ell$ . This can be understood from Eq. (1), which shows that the asymmetries are most sensitive to  $\Delta\bar{u}$  and  $\Delta\bar{d}$  at backward rapidity, where the first terms dominate due to  $x_2$  being large and thus  $q(x_2) \gg \bar{q}(x_2)$  for  $q = u, d$ .

**Table I.** Summary of  $\chi^2$  values per number of points  $N_{\text{dat}}$  for the various datasets used in this analysis.

process	$N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$	process	$N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$	process	$N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$
<b>polarized</b>			<b>unpolarized</b>			<b>SIA</b>		
inclusive DIS	365	0.93	inclusive DIS	3908	1.11	$\pi^\pm$	231	0.85
inclusive jets	83	0.81	inclusive jets	198	1.11	$K^\pm$	213	0.49
SIDIS ( $\pi^+, \pi^-$ )	64	0.93	Drell-Yan	205	1.19			
SIDIS ( $K^+, K^-$ )	57	0.36	$W/Z$ production	153	0.99			
STAR $W^\pm$	12	0.53						
PHENIX $W^\pm/Z$	6	0.63						
<b>total</b>	<b>587</b>	<b>0.85</b>	<b>total</b>	<b>4464</b>	<b>1.11</b>	<b>total</b>	<b>444</b>	<b>0.68</b>
			<b>total</b>	<b>5495</b>	<b>1.05</b>			



**Fig. 2.** Polarized sea quark asymmetry  $x(\Delta\bar{u} - \Delta\bar{d})$  from JAM (red  $1\sigma$  bands) at  $Q^2 = 10 \text{ GeV}^2$  compared with: [left panel] fit without RHIC W/Z data (yellow band), the result with positivity constraints (black lines), and a “DSSV-like” analysis (see text) (green band), and [right panel] the NNPDFpol1.1 [6] and DSSV08 [5] analyses.

#### 4. QCD Analysis

The extracted unpolarized PDFs are nearly identical to those from the recent unpolarized JAM analyses [11, 15], while the pion and kaon FFs are consistent with those from Ref. [7]. In this work we focus on the polarized PDFs, extracted from an analysis of over 900 Monte Carlo samples. The polarized antiquark asymmetry is shown in Fig. 2 and indicates a clear nonzero sea asymmetry for  $0.01 < x < 0.3$ . The inclusion of positivity constraints significantly reduces the uncertainties at  $x \gtrsim 0.1$ . Compared to the results without the RHIC  $W$  data, our analysis shows that  $\Delta\bar{u} - \Delta\bar{d}$  constrained by polarized DIS and SIDIS data alone is consistent with zero, as previously found in Ref. [7]. This further emphasizes the importance of the STAR  $W$  data for the extraction of the polarized antiquark asymmetry.

In Fig. 2 we also compare our results to the asymmetries from the DSSV [5] and NNPDF [6] groups. Interestingly, the DSSV analysis [5], which included SIDIS data, found a positive asymmetry for  $x \lesssim 0.1$ . We find that such an extraction depends strongly on positivity constraints, as well as on the propagation of FF uncertainties, and polarized PDF parametrization choice. To demonstrate this, we have carried out a “DSSV-like” analysis that excludes RHIC W/Z data, imposes positivity constraints and SU(3) symmetry, fixes the FFs to the DSS fit [16], and uses the polarized PDF parametrization from Ref. [5]. These choices, none of which involve the addition of data, greatly reduce the uncertainties on the asymmetry and generate a positive  $\Delta\bar{u} - \Delta\bar{d}$  for  $x$  below  $\approx 0.1$ , suggesting that the asymmetry from the DSSV analysis [5] may be driven by parametrization bias and theory assumptions rather than by data.

The NNPDF result [6], on the other hand, shows only a slight deviation from zero at high values of  $x$ . This is consistent with this fit taking the DSSV result [5] as the prior for  $\Delta\bar{u}$  and  $\Delta\bar{d}$ , but with  $4\sigma$  uncertainty, and including the older STAR  $W$  data [17] in their reweighting analysis. Our analysis is thus the first data-driven extraction of a nonzero polarized antiquark asymmetry.

In Fig. 3 we show the truncated integral  $\int_{0.01}^1 dx \Delta q(x)$  at  $Q^2 = 4 \text{ GeV}^2$  for the light quarks and antiquarks before and after including the RHIC  $W$  data. The lower limit of integration is chosen to roughly match the lower  $x$  limit of the data. We see an improvement for  $\Delta u^+$  and  $\Delta d^+$  of roughly 25 to 50%, while for the light antiquarks the improvement is as much as 80%. Prior to the inclusion of RHIC  $W$  data the sign of the antiquark contributions to the proton spin was unknown, but after including these data we find that  $\Delta\bar{u}$  ( $\Delta\bar{d}$ ) provides a small positive (negative) contribution to the proton spin. Prior to the inclusion of the RHIC data, the results for  $\Delta\bar{u}$  and  $\Delta\bar{d}$  depend heavily on the inclusion of positivity constraints. When the RHIC data are included, however, this dependence is significantly reduced, allowing for an extraction that is far less dependent on theoretical assumptions.

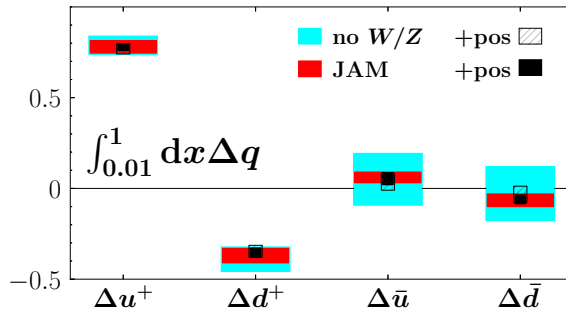


Fig. 3 Truncated integrals  $\int_{0.01}^1 dx \Delta q(x)$  at  $Q^2 = 4 \text{ GeV}^2$  for  $\Delta u^+$ ,  $\Delta d^+$ ,  $\Delta \bar{u}$  and  $\Delta \bar{d}$  from this analysis (red rectangles) compared with the fit without the RHIC W/Z data (cyan) and with positivity constraints (small hatched squares without RHIC, and black squares with RHIC). The vertical height of the bands represents  $1\sigma$  uncertainty.

## 5. Outlook

Our analysis provides the first data-driven extraction of a nonzero polarized sea asymmetry, using the latest  $W$ -lepton data from RHIC, within a simultaneous global QCD analysis of polarized PDFs, unpolarized PDFs, and pion and kaon FFs. This also provides the first self-consistent extraction of the light quark polarizations and shows a nonzero contribution to the proton's spin from the light antiquarks.

With the Jefferson Lab 12 GeV upgrade and the Electron-Ion Collider (EIC), future experiments will access new information on the spin structure of the proton. In particular, the high-luminosity CLAS12 SIDIS experiment using  $K$  production [18] will provide precise SIDIS data to complement the  $W$ -lepton production data from RHIC. The EIC should bring forth new information on all polarized PDFs, in particular the strange and gluon PDFs [19], while also extending the kinematic coverage of polarized DIS experiments to lower  $x$  and higher  $Q^2$ .

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