

STUDIES ON RADIATION SHIELDING  
OF ENERGY DOUBLER MAGNETS (1)

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I Introduction

Several measurements have been made to understand quench problems of superconducting magnets due to beam energy deposition in magnet coils.<sup>1</sup> Some preliminary analysis of quench problems of Energy Doubler magnets due to beam losses at injection, during acceleration and at extraction have already appeared.<sup>2,3</sup> In this report we discuss radiation problems due to beam losses during acceleration and extraction. We also discuss various proposed schemes to minimize such problems.

II. Limits of Energy Deposition

Figure 1 shows an estimate of the relation between limits of energy deposition in superconducting coils and magnet excitation current.<sup>2</sup> These limits appear consistent with measurements.<sup>1</sup> Table I summarizes the limits for extraction beam losses.

Table I. Limits of energy deposition at extraction energy.

Mode	Limits
Slow	$4 \text{ mw.g}^{-1}$
Fast (1ms)	$1 \text{ mj.g}^{-1}.\text{pulse}^{-1}$
Abort (20 $\mu\text{s}$ )	$0.5 \text{ mj.g}^{-1}.\text{pulse}^{-1}$

We assumed that  $I/I_{\max} \approx 0.9$  for the following reasons. Fristly, the highest energy deposited inside the coils always appears near the horizontal median plane because of magnetic field effects on charged secondary particles and because of the curvature of the magnets. The magnetic field strength inside the coils around the horizontal median plane is about 10% lower than the maximum field which appears at the top and bottom corners of the coils. Secondly, it seems to be reasonable to operate the Energy Doubler at a slightly lower current level than the extreme quench level.

### III. Extraction Beam Losses

Extraction beam losses are primarily caused by beam scattering and nuclear interactions in the electrostatic septum wires and Lambertson septum magnets.

#### (A) Radiation Due to Scattering From Electrostatic Septum Wires:

The program CASIM<sup>4</sup> was modified to calculate energy deposited in Doubler magnets.<sup>5</sup> The wire septum and dipole magnet string have been modelled in all revelant geometric detail including the presence of electric and magnetic fields. However, in the present study a string of four quadrupole magnets to be placed upstream of the dipole string were omitted. A schematic drawing of the arrangement is shown in Figure 2. Figure 3 shows the energy deposition in the dipole string as a function of depth for a shallow (0.59 cm thick) radial region located near the inner radius of the coils for a 1000 GeV proton beam (0.2 cm high) uniformly distributed over the lateral dimension of the wire. Curves are drawn to guide the eye. The collimator upstream of the magnet string was omitted. Two distinct peaks are seen. The first peak is due to radiation

through the upstream magnet surface and can be suppressed by placing a collimator in front of the magnet. This is shown in Figure 4. An iron collimator with a hole radius of 2 cm and a length of 2m was placed as shown in Figure 2. The first peak was suppressed by a factor of roughly 100. The second peak is most prominent near  $\phi = \pi$  and is mainly due to high energy photons from neutral pions produced in the forward direction. The collimator reduced the energy density at the second peak in general, but its maximum value (near  $\phi = \pi$ ) was not significantly lowered. This large energy deposition of  $3 \times 10^{-3} \text{ GeV.cm}^{-3} (\text{inc. proton})^{-1}$  therefore presents a more difficult problem. Table II gives the maximum numbers of protons that can be extracted under these conditions. The extraction inefficiency was assumed to be 2.5%, i.e. 2.5% of protons strike septum wires.

Table II. Numbers of protons that can be extracted for the electrostatic septum arrangement shown in Figure 2.

Mode	Number of Protons
Slow	$2.7 \times 10^{12} \text{ protons/sec}$
Fast (1ms)	$0.7 \times 10^{12} \text{ protons/pulse}$

Those limits are unacceptably low, particularly for fast spill. T. Collins proposed a modification of a long straight section to provide a large  $\beta$  region in the horizontal plane<sup>6</sup>, shown in Figure 5. If conventional magnets are used at the middle of the straight section, this will considerably reduce energy deposition in the Doubler magnets. Detailed calculations are in progress. In this arrangement good field tracking between superconducting and conventional magnets is required.

A bump as large as 4 mrad may not be needed to achieve a sufficient reduction of energy deposition density at the Doubler

magnets. Then, as shown in Figure 6, a simpler bump by conventional magnets only which does not require the same high degree of field tracking is more attractive. The direction of the bump can be outward or inward in the radial direction.

(B) Radiation Due to Scattering From Lambertson Magnets:

Protons which are scattered out of the electrostatic septum wires may strike the Lambertson magnet septum region and may cause radiation in downstream Doubler magnet coils. Figure 7 shows the schematic drawing of the arrangement used in the calculation. The energy density distribution in the Doubler magnet is shown in Figure 8. It was assumed that a 1000 GeV proton beam (0.2 cm high) was uniformly distributed over the central (narrowest) region of the septum, and struck the upstream surface of the septum perpendicularly. Nearly all incident protons underwent nuclear interactions in the septum. The large peaks at the upstream end can be suppressed by a collimator as in the electrostatic septum case. Unlike in the latter case the second peak corresponding to neutral pion production does not appear there because the Lambertson magnets have enough absorbing material at the septum region in the forward direction. The maximum energy density is about  $4 \times 10^{-4}$  GeV.cm<sup>-3</sup> (inc. proton on Lambertson magnet).<sup>-1</sup> Table III gives the maximum numbers of protons that can be extracted through the Lambertson magnets. The interaction rate of protons at the Lambertson magnets was assumed to be 0.63% which is obtained if one quarter of protons which strike the electrostatic septum wires intersect the Lambertson magnets.

Table III. Numbers of protons that can be extracted through the Lambertson magnets.

Mode	Number of Protons
Slow	$7.9 \times 10^{13}$ protons/sec
Fast (1ms)	$2.0 \times 10^{13}$ protons/pulse

These limits are quite acceptable. No special arrangement seems to be required for the Lambertson magnet area. However, the assumption that the beam is uniformly distributed on the septum region (Figure 7) is perhaps questionable. Figure 9 shows the energy density distribution when the protons were incident uniformly distributed over the first 3.2 m on the side surface of the no-field region of the septum with an angle of 30  $\mu$ rad. The energy density is much larger than in Figure 8. Furthermore, a second peak appear around  $\phi = \pi$  as in the electrostatic septum case. If a large fraction of the protons lost at the Lambertson magnet is typified by this mode, then the limits given in Table III must be substantially lowered.

The Lambertson magnet arrangement for the beam abort system has a lower limit of energy deposition because of the single-turn extraction. However, in principle it may have less problems since the extraction can be done more cleanly. Of the order of  $6 \times 10^{10}$  protons are allowed to hit the septum.

#### IV. Beam Losses During Acceleration

Figure 10 shows the schematic drawing of the case in which protons strike the vacuum chamber of the Doubler magnet on the inside edge. The proton energy is 1000 GeV. Figures 11 and 12 show energy density distributions in the Doubler magnet coils as functions of azimuthal angle and coil radius when protons interact with the vacuum chamber at  $\phi = 0$  (at the accelerator inner

radius as indicated in Figure 10). For the vacuum chamber radius of 3.68 cm, the energy density distribution has a sharp peak near the proton interaction point  $\phi = 0$  and reaches about  $2 \text{ GeV cm}^{-3}$  ( $\text{interacting proton})^{-1}$ , which corresponds to about  $4 \times 10^{-8} \text{ mj. g}^{-1}$  ( $\text{interacting proton})^{-1}$ . This peak density is much smaller and has a gentler dependence when the vacuum chamber is placed at 3.18 cm inner radius. The second peak which appears at the downstream outer radius ( $\phi = \pi$ ) is again due to neutral pions produced in the forward direction. The peak values are roughly the same for the both cases.

Figures 13 and 14 show energy density distributions when protons interact with the inside edge of the vacuum chamber at  $\phi = \pi$  (at the outer accelerator radius). Essentially all the characteristics of the energy density distributions are identical to those distributions shown in Figures 11 and 12 with the  $\phi$  angles reversed except for the second peaks corresponding to neutral pions. The latter cases do not have a neutral pion peak for geometrical reasons.

The first peak can be suppressed by long scrapers ( $\gtrsim 3\text{m}$ ), for example at medium straight sections. However, the second peak of about  $1 \times 10^{-9} \text{ mj g}^{-1}$  ( $\text{interacting proton})^{-1}$  when protons interact with the vacuum chamber inside edge at the accelerator inner radius is difficult to reduce by a simple arrangement. If long straight sections are available for beam scraping, an arrangement of a beam scraper and absorbers with a beam dump system of conventional magnets seems to be a good solution, as shown in Figure 6. If no long straight section is available for beam scraping, T. Collins proposed a special beam scraper arrangement at medium straight sections, as shown in Figures 15 and 16. This arrangement requires special dipole magnets of a larger aperture and half the length of the normal Doubler dipole magnet. Detailed studies are in progress.

### V. Conclusions

A special beam bump arrangement is required at the long straight section where the electrostatic septum is installed for extraction. The Lambertson magnet areas for both the normal extraction and the beam abort system may not require any special system other than simple beam collimators.

Beam scraper systems which localize beam losses and minimize radiation to the Doubler magnet coils during acceleration may have to be installed in medium straight sections. Further detailed studies are needed.

### References

1. H. Edwards, C. Rode, and J. McCarthy, IEEE Trans. on Magnetics, Vol. 1, Mag 13, 1977, p 666, and B. Cox, P.O. Mazur, and A. Van Ginneken, "Sensitivity of an Energy Doubler Dipole to Beam Induced Quenches", TM-828A, 1978.
2. H. Edwards, "Beam Abort System", Internal Report, June 3, 1977.
3. T. Collins, "Radiation on Magnet Coils from the Wire Septum", Internal Report, 1978.
4. A. Van Ginneken, "CASIM, A Program to Simulate Hadronic Cascades in Bulk Matter", FN-272, January 1975.
5. A. Van Ginneken, "Distribution of Heat Due to Beam Loss in Energy Doubler/Saver Type Superconducting Magnets", TM-685, September 13, 1976.
6. T. Collins, "High-beta Straight Sections for the Doubler", UPC No. 14, November 14, 1978.
7. T. Collins, "Charged Particle Distributions Estimate", Internal Report, 1978.

Figure Captions

Figure 1. Limits of energy deposition in Doubler superconducting magnet coils.

Figure 2. Schematic drawing of the arrangement of the electrostatic wire septum for extraction at a long straight section. The second unit is tilted by 30  $\mu$ rad away from the normal orbit so that wires of the second unit do not intersect protons. The  $\phi$  angle is measured from the inner accelerator radius in the median plane.

Figure 3. Energy density distributions in Doubler magnet coils due to scattering from the extraction electrostatic wire septum for a shallow radial region of 3.81 cm to 4.40 cm. The incident proton energy was 1000 GeV. No collimator was present.

Figure 4. Energy density distributions in Doubler magnet coils due to scattering from the extraction electrostatic wire septum for a shallow radial region of 3.81 cm to 4.40 cm. The incident proton energy was 1000 GeV. An iron collimator of a hole radius of 2 cm and a length of 2 m was placed upstream of the Doubler magnets.

Figure 5. Schematic drawing of a modified high- $\beta$  long straight section by T. Collins. (See Reference 3.)

Figure 6. Schematic drawing of a bump arrangement of conventional B-2 magnets at a long straight section.

Figure 7. Schematic drawing of the Lambertson septum magnet arrangement for extraction.

Figure 8. Energy density distributions in Doubler magnet coils due to scattering from the extraction Lambertson septum magnets for a shallow radial region of 3.81 cm to 4.40 cm. It was assumed that 1000 GeV protons were distributed uniformly over the central region of the septum and struck perpendicularly the upstream surface of the septum. No collimator was present.

Figure 10. Energy density distributions in Doubler magnet coils due to scattering from the extraction Lambertson septum magnets for a shallow radial region of 3.81 cm to 4.40 cm. It was assumed that 1000 GeV protons struck the side surface of the septum from the normal beam orbit side with an angle of 30  $\mu$ rad, uniformly distributed over the first 3.2 m of the septum. No collimator was present.

Figure 11. Schematic drawing of the case in which 1000 GeV protons strike the vacuum chamber of a Doubler magnet. The proton beam was assumed to have a distribution of a  $\delta$ -function in the vertical direction and interact the inside edge of the vacuum chamber in the median plane. The vacuum chamber radius is from 3.18 to 3.31 cm and 3.68 to 3.81 cm. The  $\phi$  angle is measured from the inner accelerator radius in the median plane.

Figure 12. Energy density distributions in Doubler magnet coils due to beam losses on the vacuum chamber. The proton beam of 1000 GeV struck the inside edge of the vacuum chamber at the inner accelerator

radius ( $\phi = 0$ ). The vacuum chamber radius was from 3.68 cm to 3.81 cm. The inner radius of the superconducting coils was 3.81 cm.

Figure 13. Energy density distributions in Doubler magnet coils due to beam losses on the vacuum chamber. The proton beam of 1000 GeV struck the inside edge of the vacuum chamber at the inner accelerator radius ( $\phi = 0$ ). The vacuum chamber radius was from 3.18 to 3.31 cm. The inner radius of the superconducting coils was 3.81 cm.

Figure 14. Energy density distributions in Doubler magnet coils due to beam losses on the vacuum chamber. The proton beam of 1000 GeV struck the inside edge of the vacuum chamber at the outer accelerator radius ( $\phi = \pi$ ). The vacuum chamber radius was from 3.68 cm to 3.81 cm. The inner radius of superconducting coils was 3.81 cm.

Figure 15. Energy density distributions in Doubler magnet coils due to beam losses on the vacuum chamber. The proton beam struck the inside edge of the vacuum chamber at the outer accelerator radius ( $\phi = \pi$ ). The vacuum chamber radius was from 3.18 cm to 3.31 cm. The inner radius of superconducting coils was 3.81 cm.

Figure 15. Schematic drawing of a modified medium straight section with two half-bending magnets and beam scraper arrangement proposed by T. Collins (See Reference 3.).

Figure 16. Trajectories of neutral particles scattered from edges of the beam scraper in a Collins' modified medium straight section (See Reference 3.).

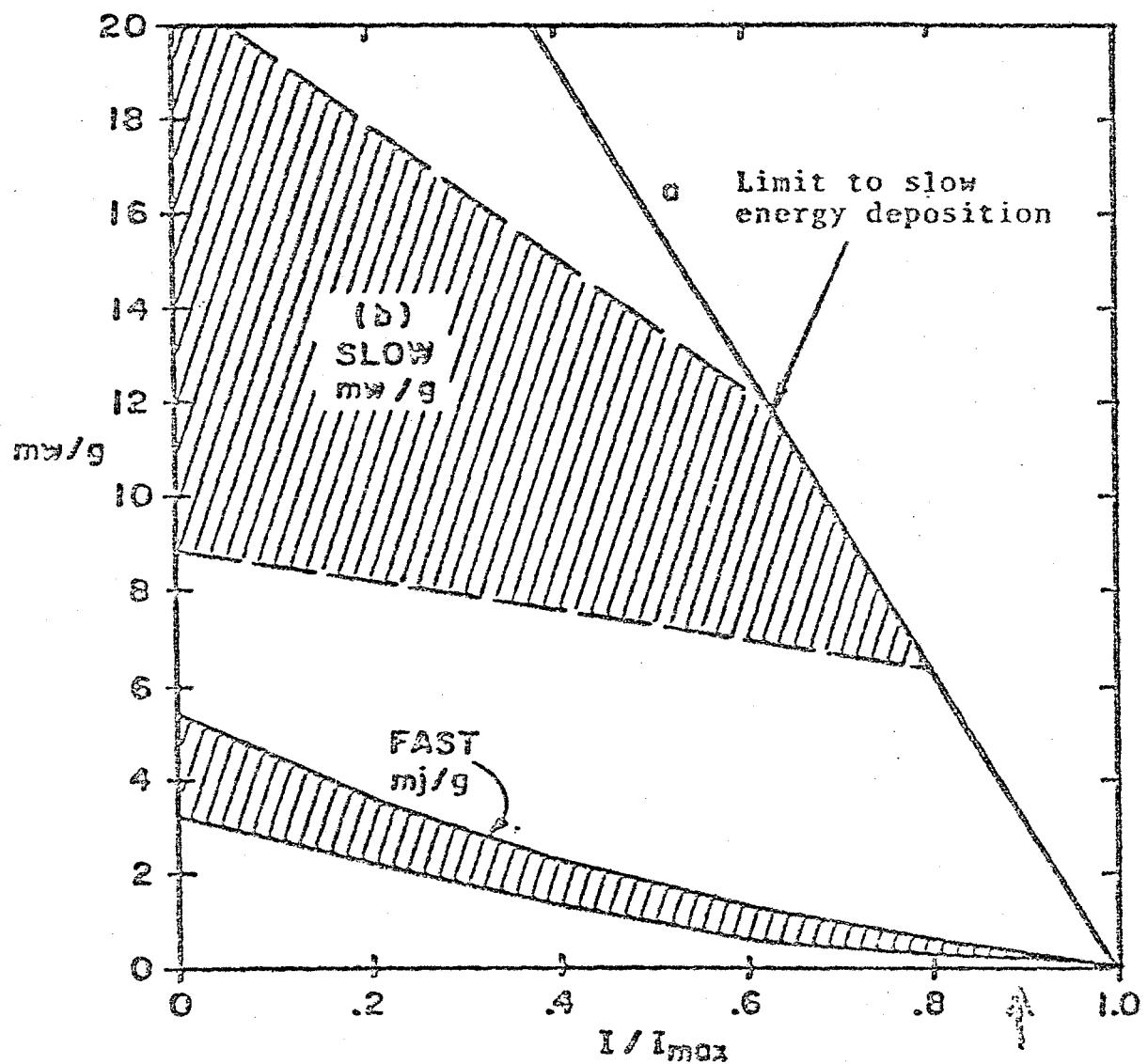


Figure 1. LIMITS OF ENERGY DEPOSITION IN DOUBLER SUPERCONDUCTING MAGNETS.

ELECTROSTATIC WIRE SEPTUM

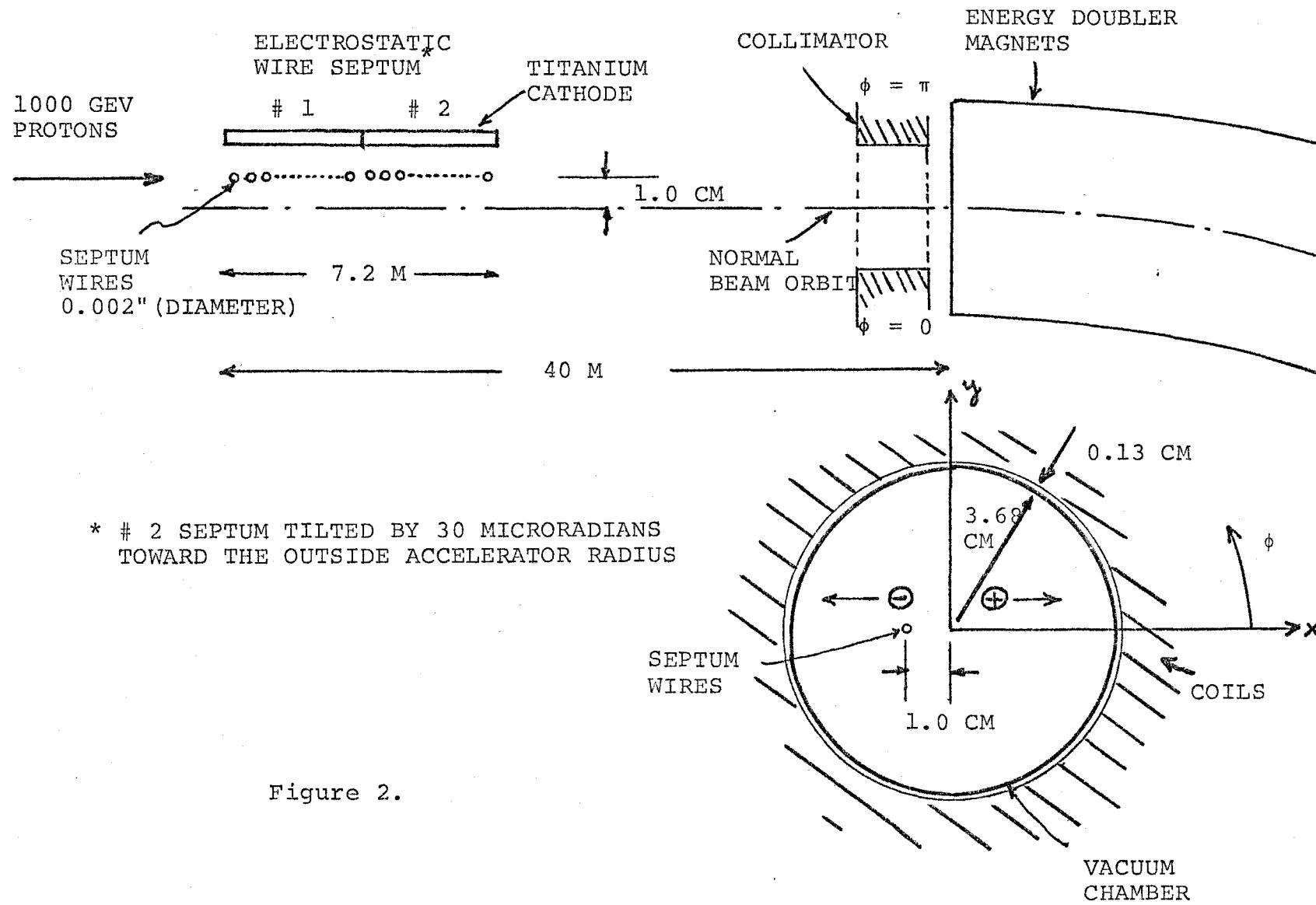
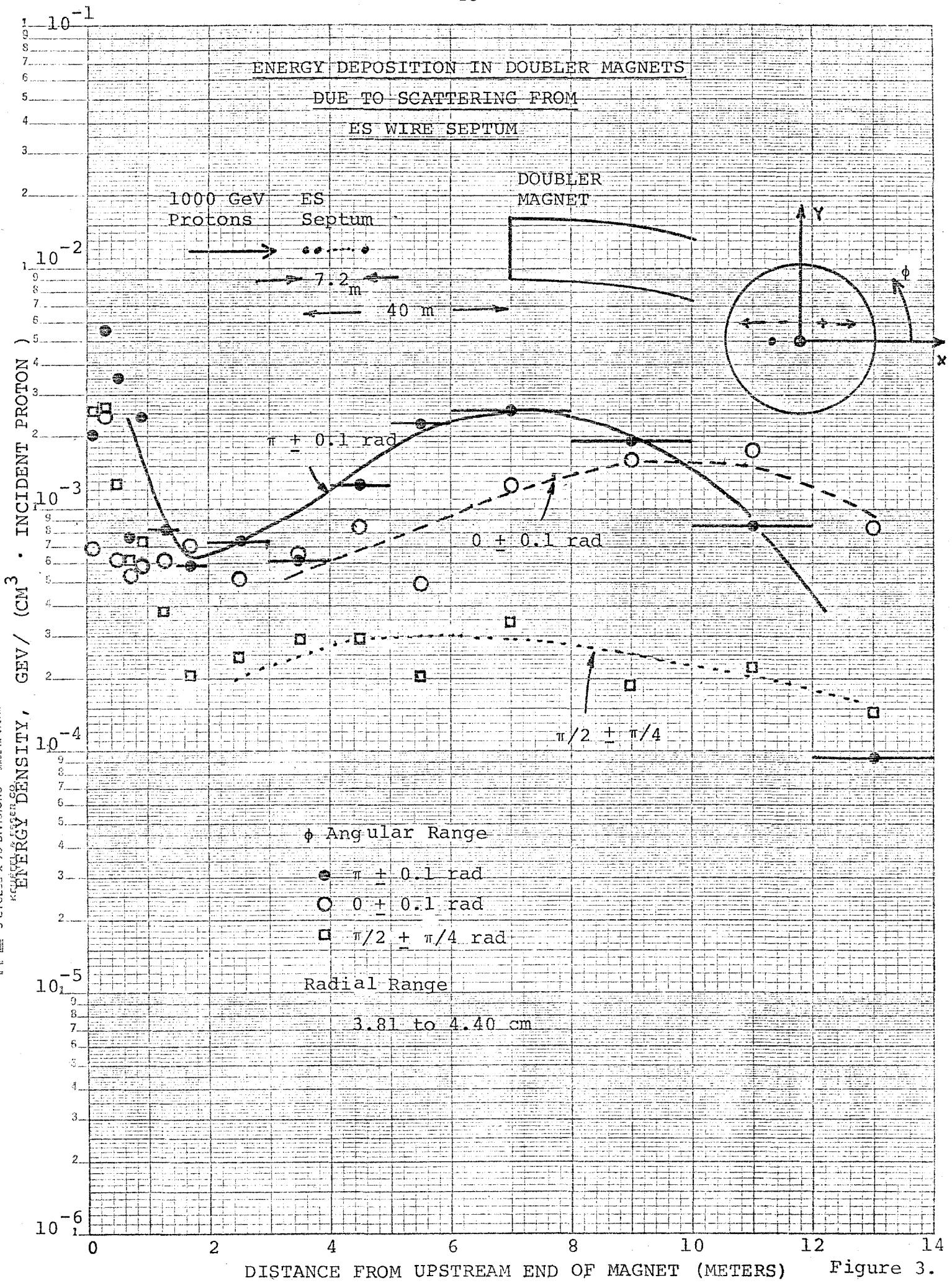


Figure 2.



$10^{-1}$ 

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM

ES WIRE SEPTUM

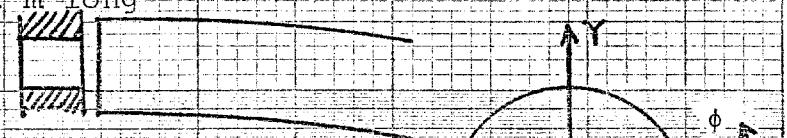
Collimator

2 cm radius

2 m long

1000-GeV  
ProtonsES  
Septum

7.2 m

40 m Doubler  
MagnetENERGY DENSITY, GEV / ( CM<sup>2</sup> · INCIDENT PROTON ) $10^{-2}$  $10^{-3}$  $10^{-4}$  $10^{-5}$  $10^{-6}$  $\pi \pm 0.1 \text{ rad}$  $0 \pm 0.1 \text{ rad}$  $\pi/2 \pm \pi/4$ 

Angular Range

 $\bullet \pi \pm 0.1 \text{ rad}$  $\circ 0 \pm 0.1 \text{ rad}$  $\square \pi/2 \pm \pi/4 \text{ rad}$ 

Radial Range

3.81 to 4.40 cm

0

DISTANCE FROM UPSTREAM END OF MAGNET (METERS)

Figure 4.

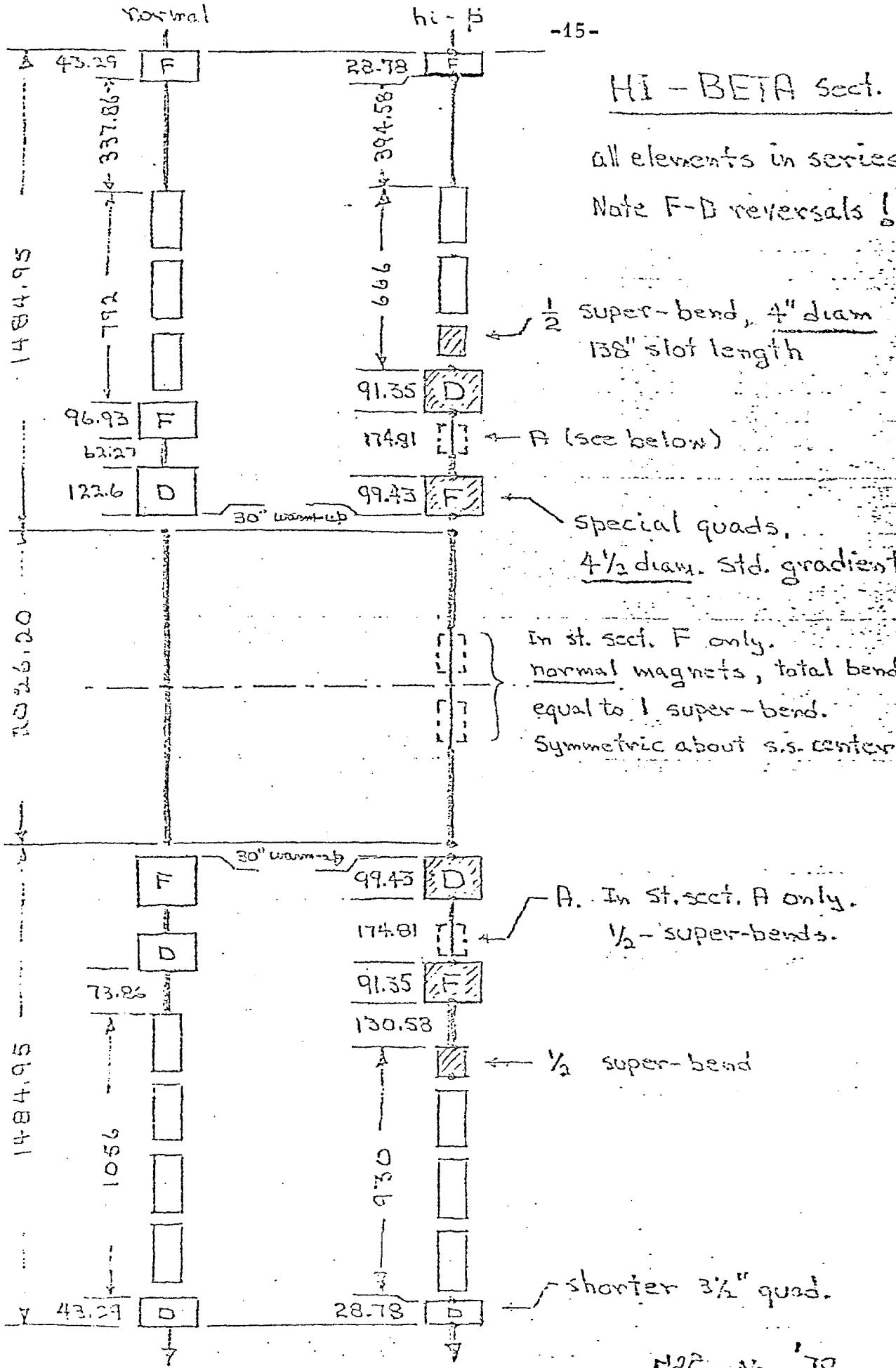
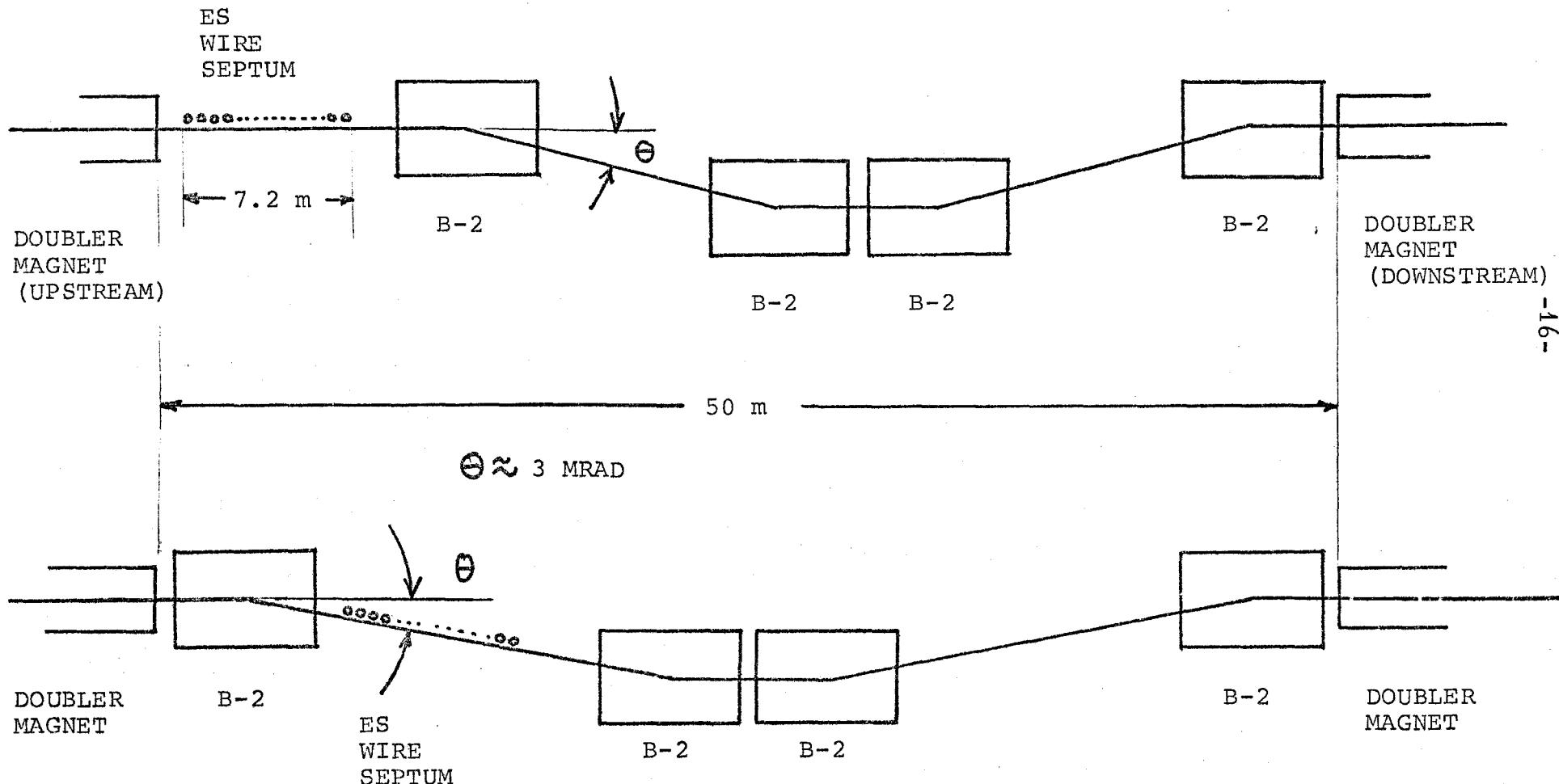


Figure 5.

SIMPLE BUMP BY CONVENTIONAL MAGNETS ( B-2 )

AT F0 LONG STRAIGHT SECTION\*



\* COLLIMATORS PLACED INSIDE AND/OR IN FRONT OF B-2 MAGNETS  
DOWNSTREAM OF THE ES WIRE SEPTUM.

Figure 6.

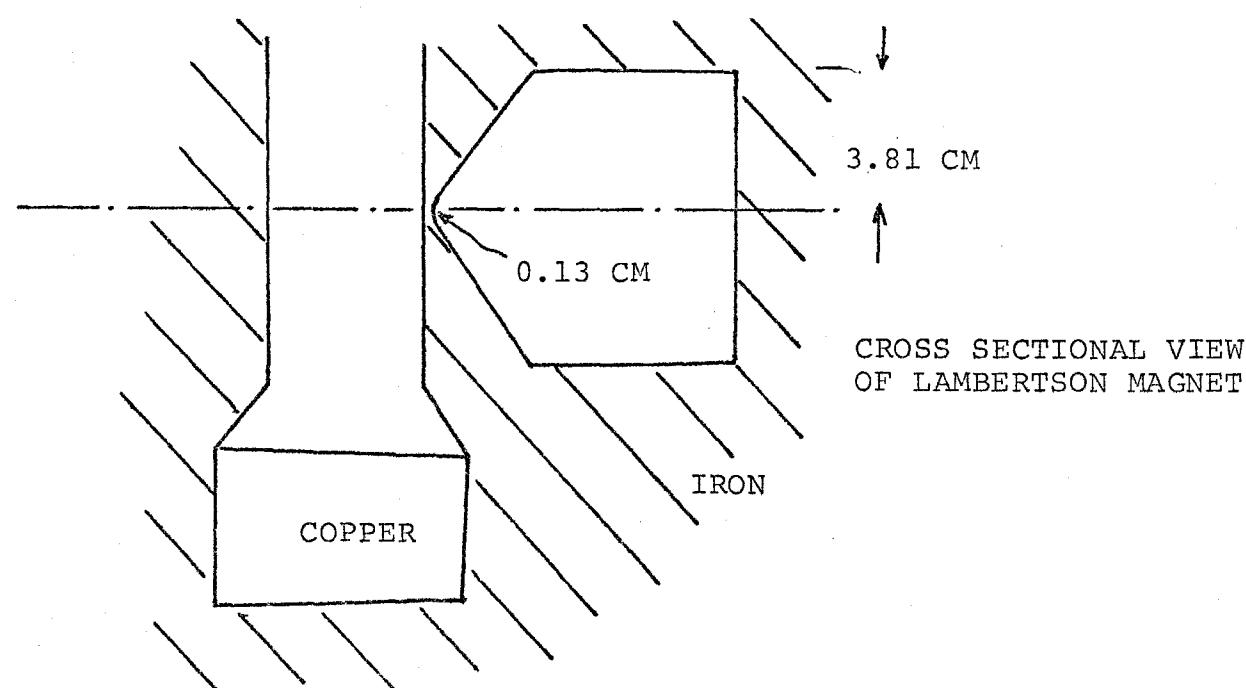
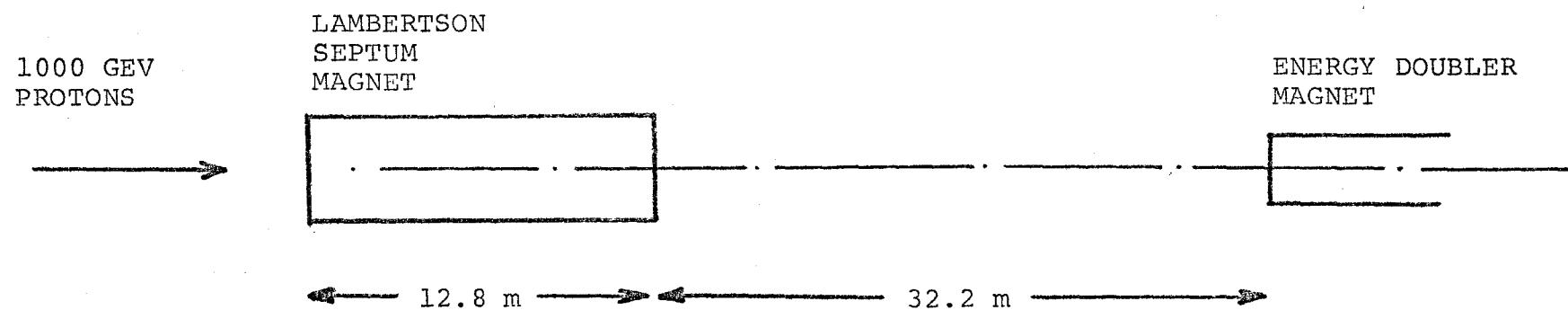
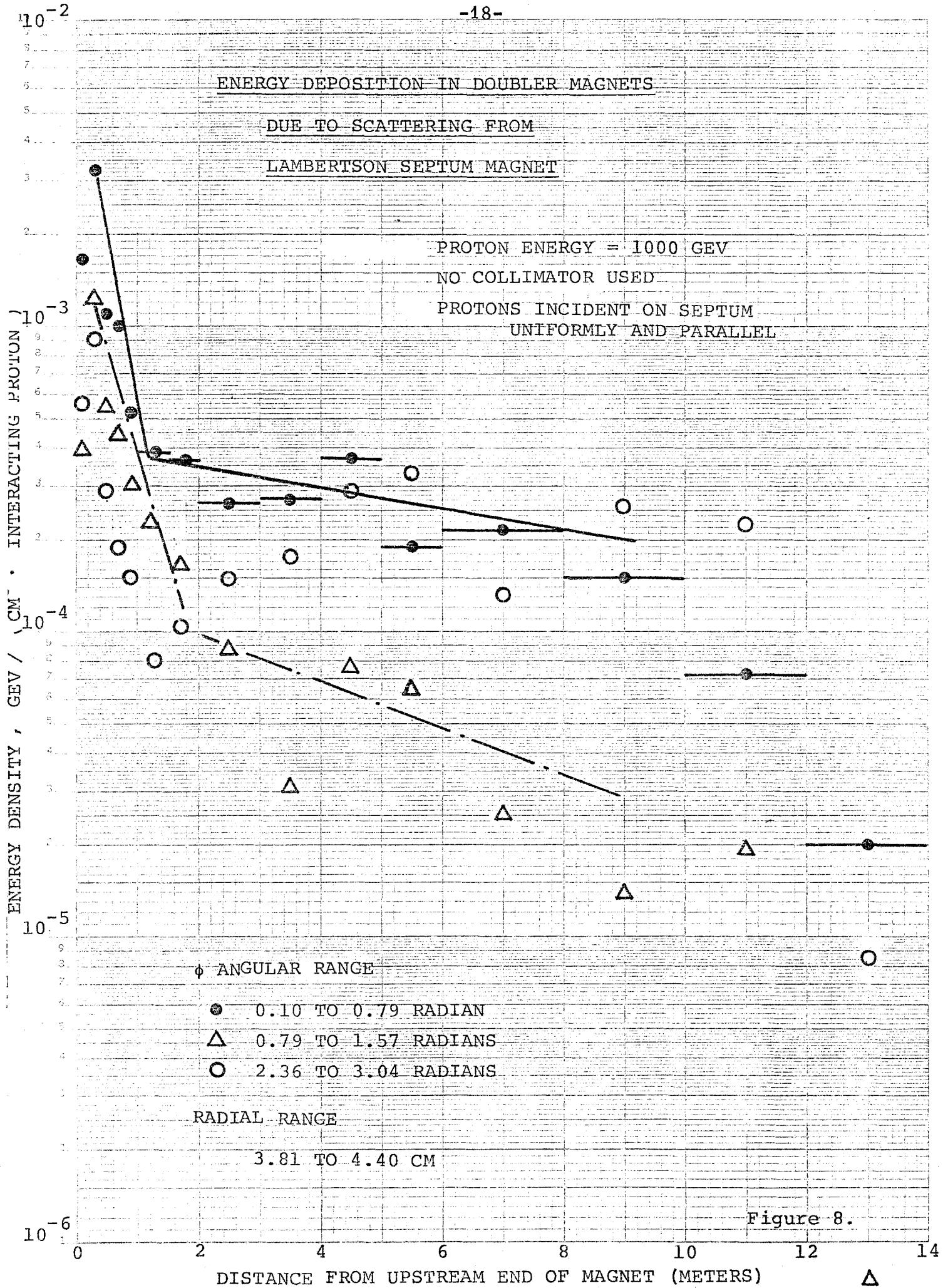


Figure 7.



ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
LAMBERTSON SEPTUM MAGNET

PROTON ENERGY = 1000 GEV

NO COLLIMATOR USED

PROTONS INCIDENT FROM THE SIDE SURFACE  
OF NO FIELD REGION WITH AN ANGLE  
OF 30 MICRORADIAN

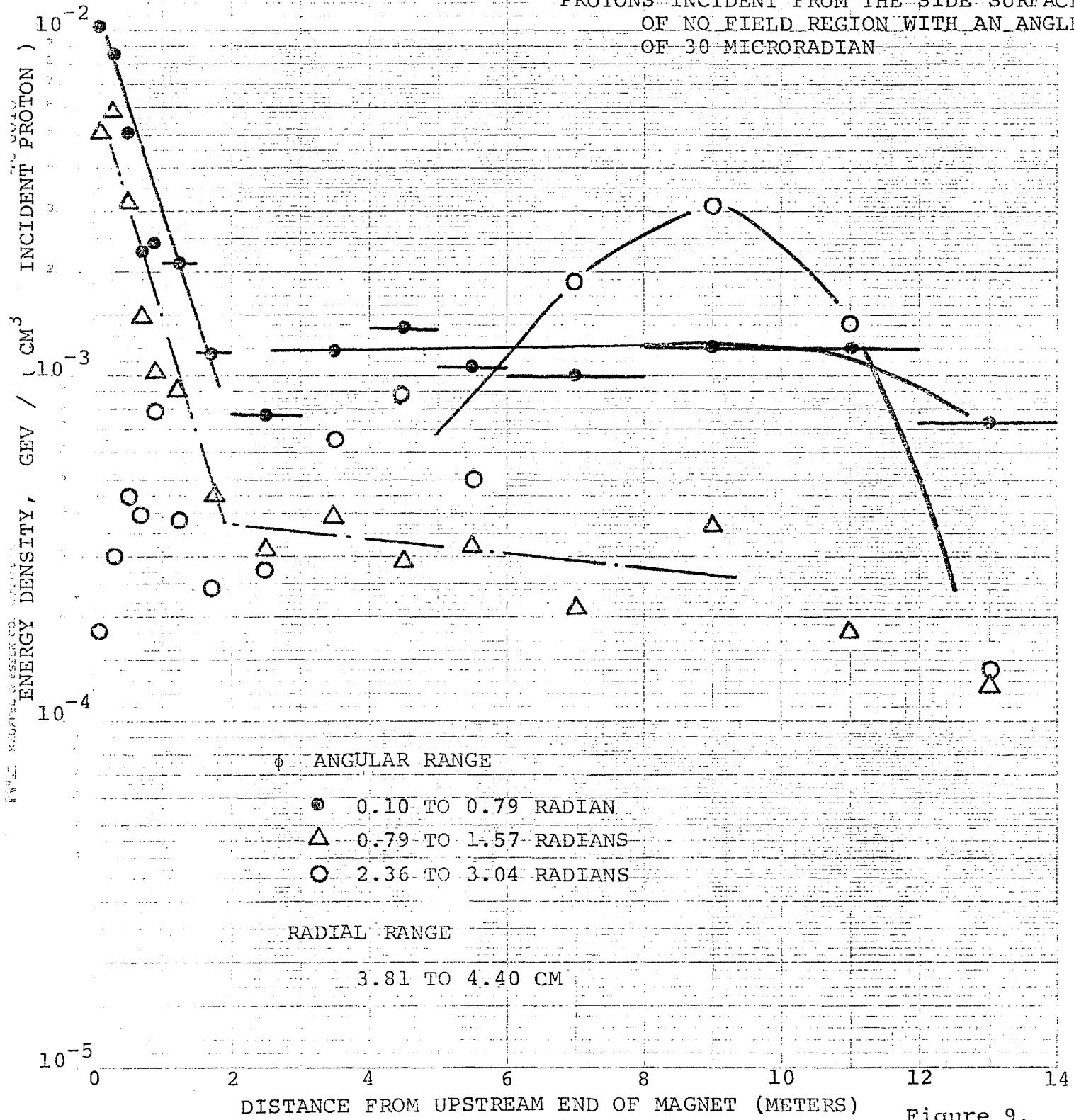


Figure 9.

Straight Hit on Vacuum Chamber by 1000 GeV Protons

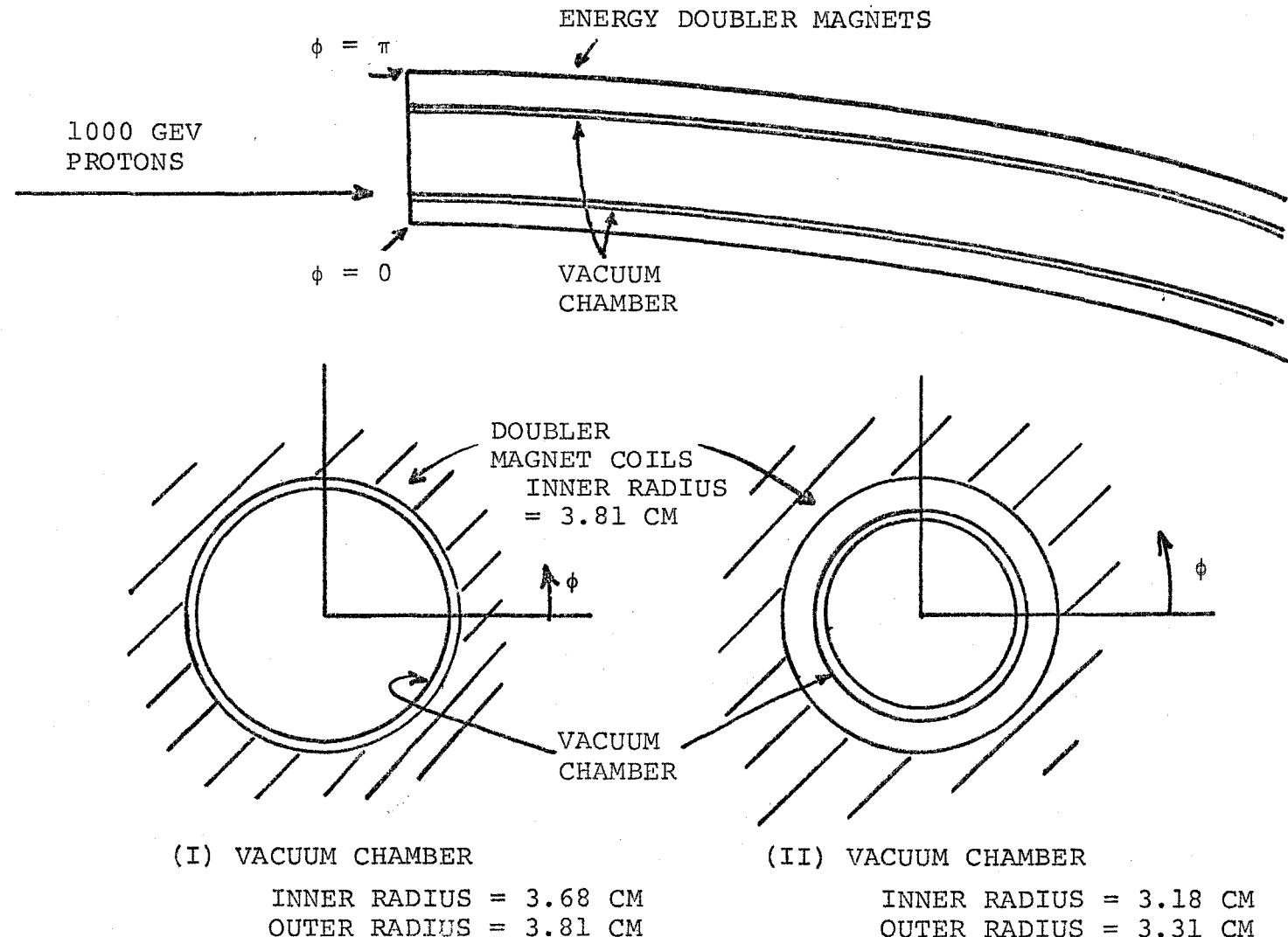
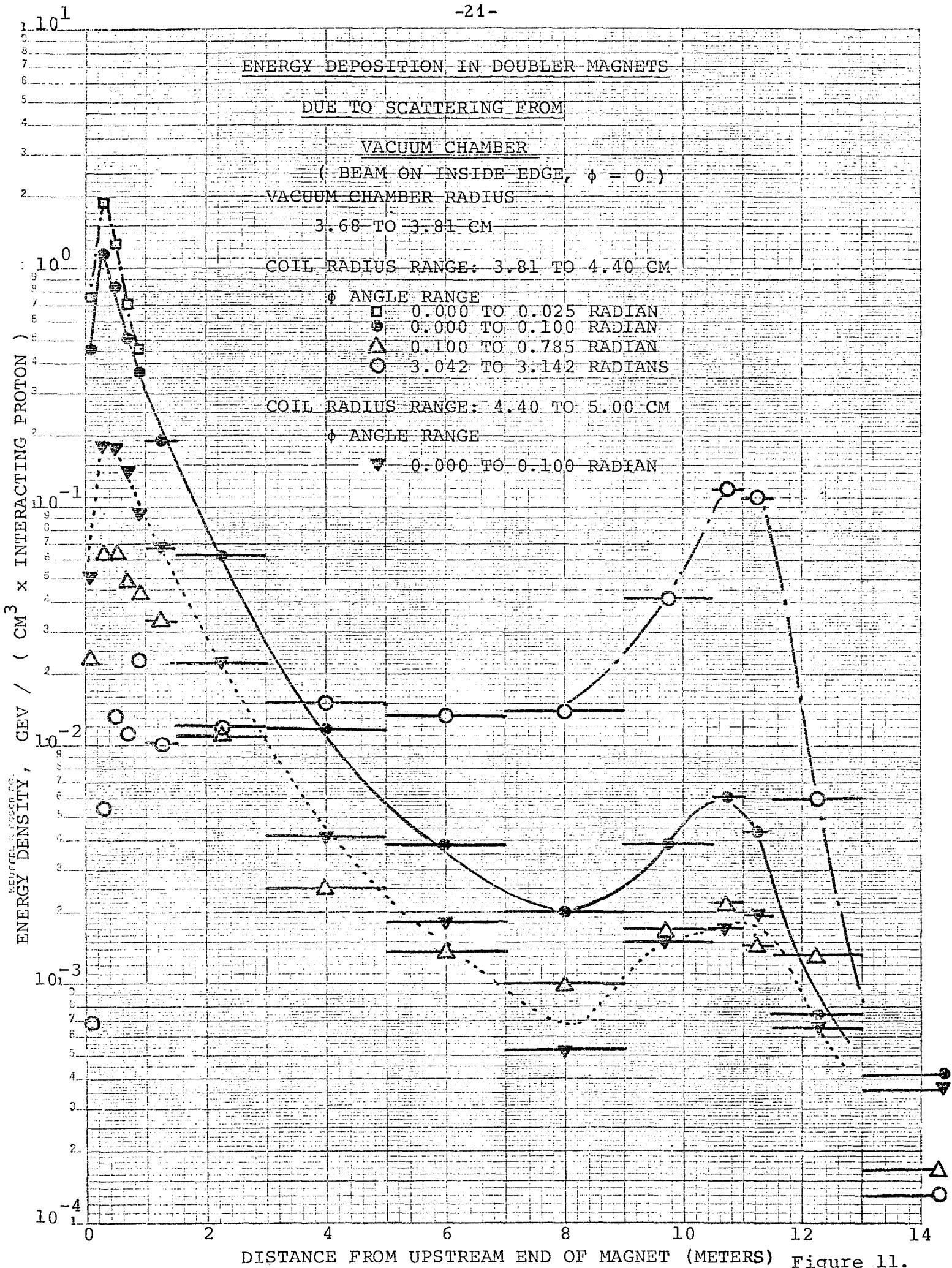


Figure 10.



ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

( BEAM ON INSIDE EDGE,  $\phi = 0$  )

VACUUM CHAMBER RADIUS:

3.18 TO 3.31 CM

COIL RADIUS RANGE: 3.81 TO 4.40 CM

$\phi$  ANGLE RANGE

□ 0.000 TO 0.025 RADIANS

● 0.000 TO 0.100 RADIANS

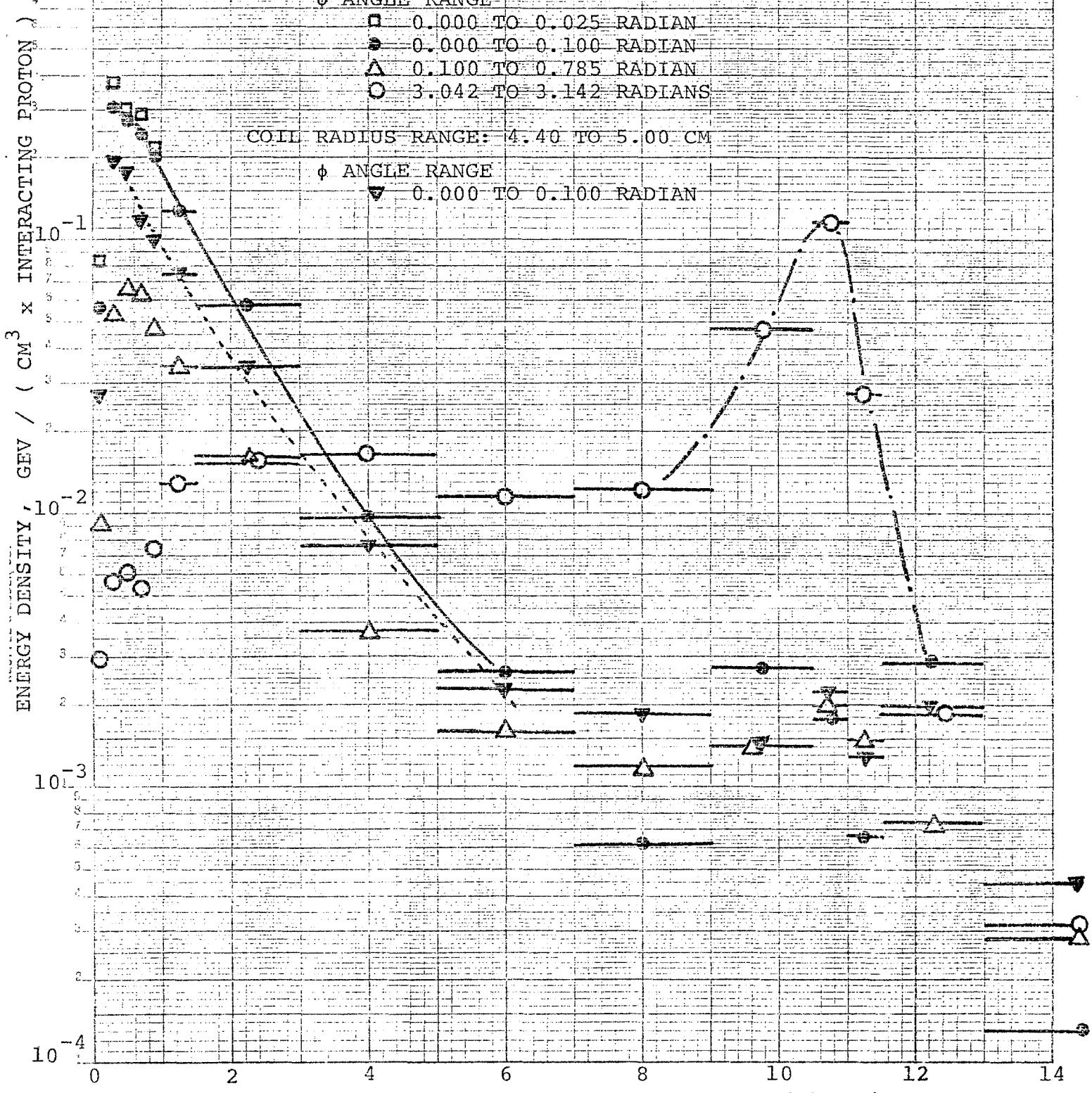
△ 0.100 TO 0.785 RADIANS

○ 3.042 TO 3.142 RADIANS

COIL RADIUS RANGE: 4.40 TO 5.00 CM

$\phi$  ANGLE RANGE

▽ 0.000 TO 0.100 RADIANS



DISTANCE FROM UPSTREAM END OF MAGNET (M)

Figure 12.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

( BEAM ON INSIDE EDGE,  $\phi = \pi$  )

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COIL RADIUS RANGE: 3.81 TO 4.40 CM

$\phi$  ANGLE RANGE

● 0.000 TO 0.100 RADIANS

△ 2.350 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIUS RANGE: 4.40 TO 5.00 CM

$\phi$  ANGLE RANGE

▼ 3.042 TO 3.142 RADIANS

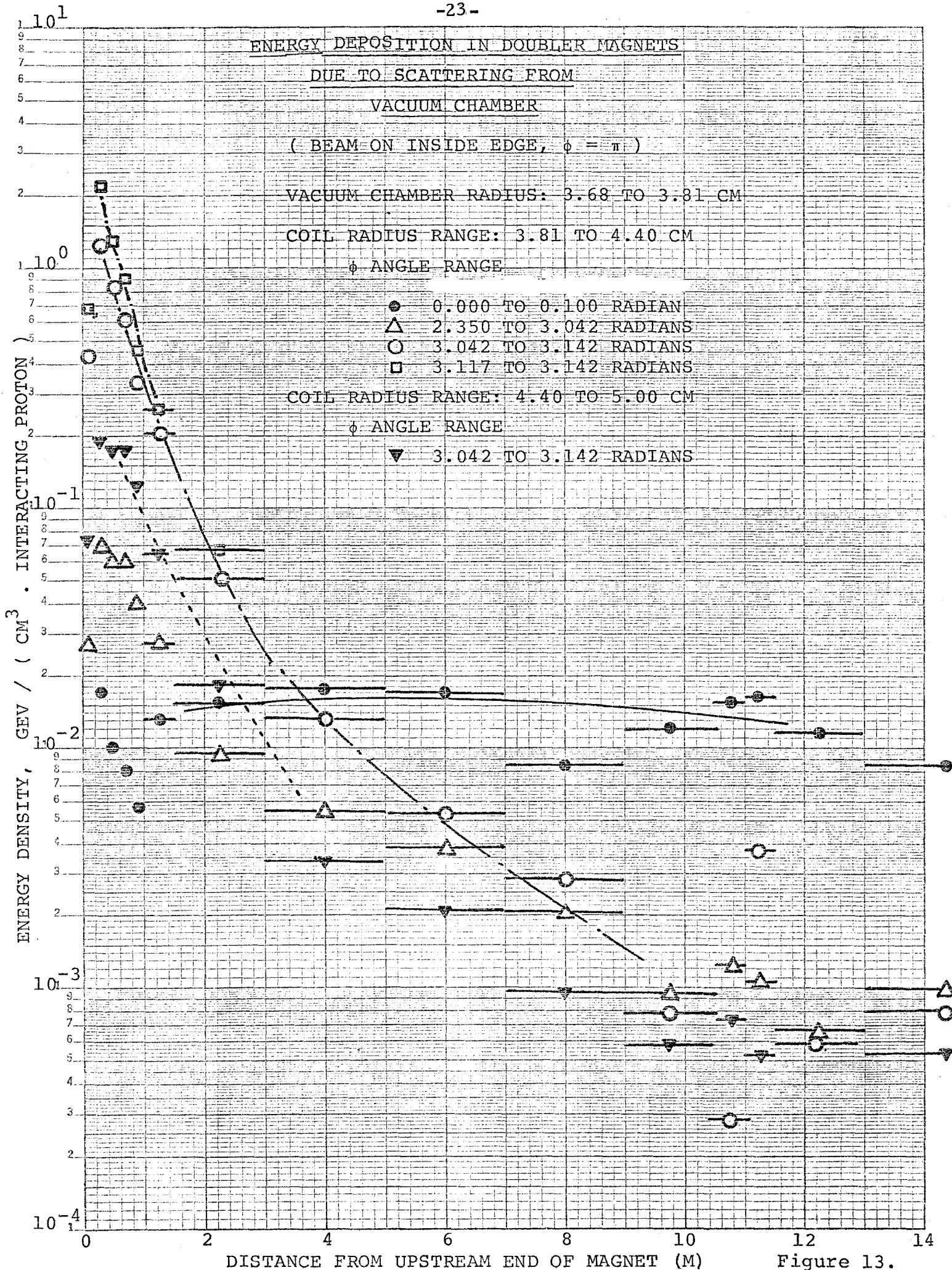


Figure 13.

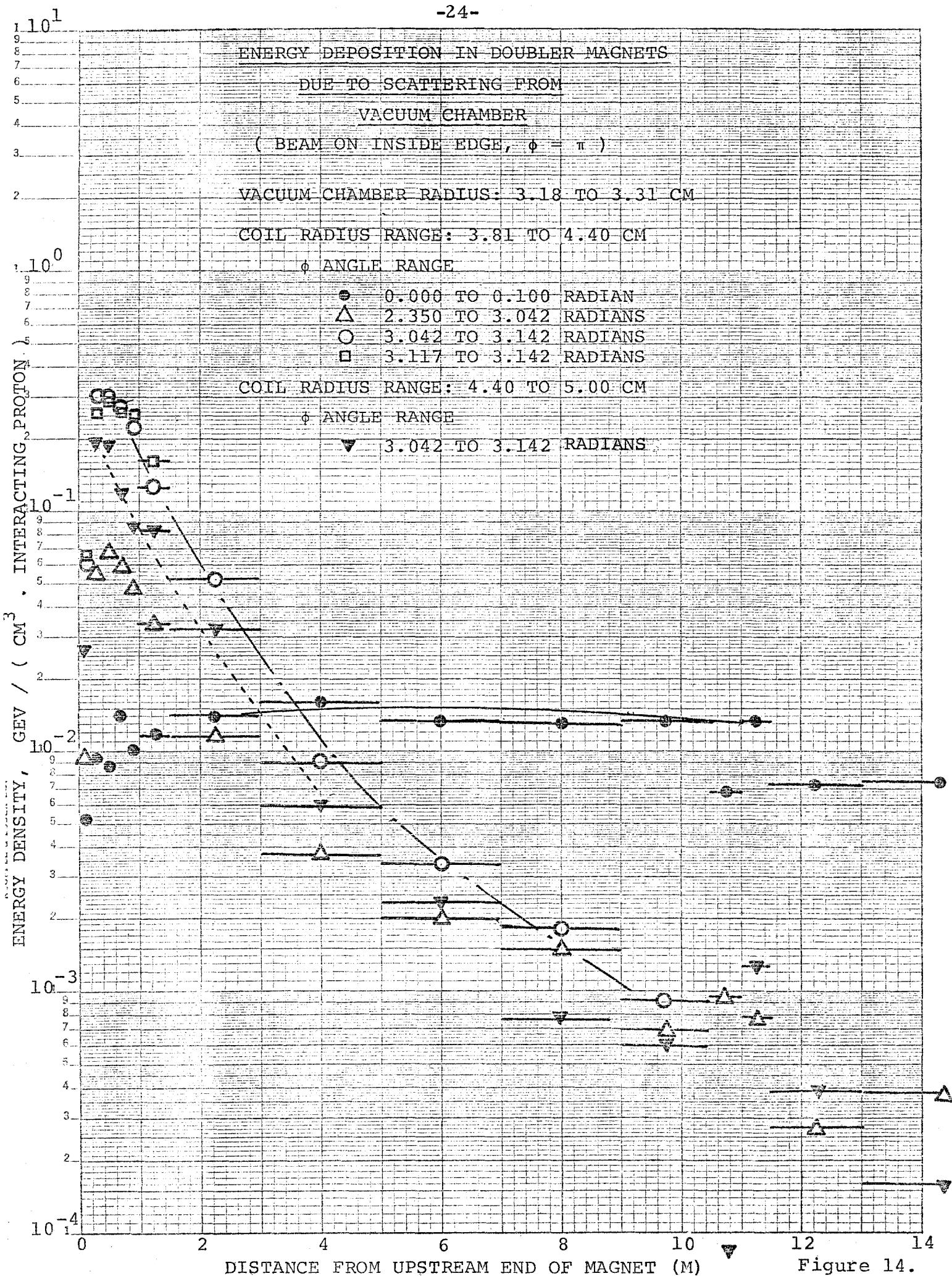
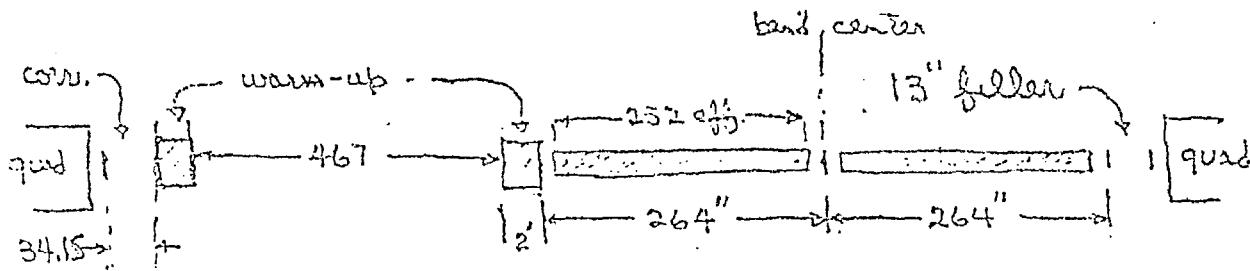
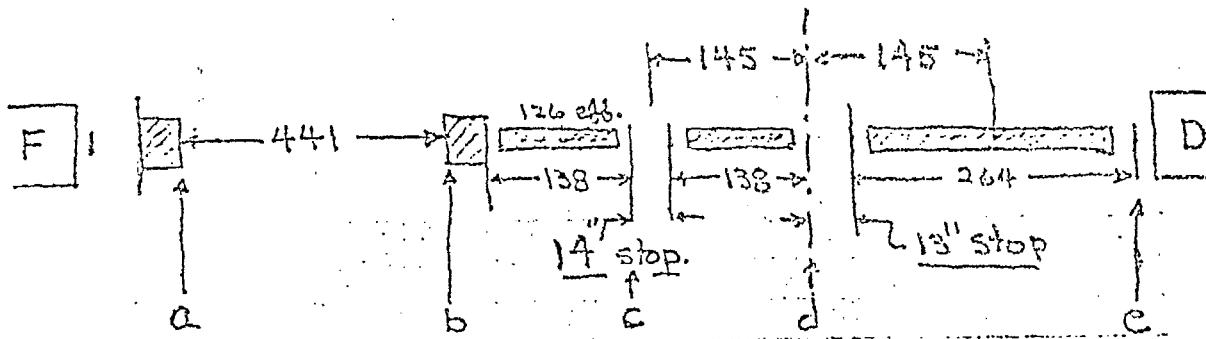


Figure 14.



Medium straight section as designed.



Medium straight - modified with 2 half-bends and using  
filler space to create spaces for absorbers.

Figure 15.

From outside, 1.45 to 2.35 mrad hit 3rd mag.  
center, .80 to 1.46 mrad " "

1.45 mrad.  
+ .769 "

beam  
 $\pm .8$ " at  $\beta_0$

- .769 "

note: no neutrals hit magnets

from this edge, (the low mom. edge).

441" COLLIMATOR.

beam  $\pm .61$   
ap.  $\pm .65$

center - .28

beam  $\pm .55$

ap.  $\pm .60$

center - 1.18  
beam  $\pm .51$   
ap.  $\pm .56$

special 4" magnets  
wall at  $\pm 1.8$ "

regular magnet

-100" -

Figure 16.

SUPPLEMENT TO UPC 30

We have discussed some of the typical cases in which the proton beam strikes the vacuum chamber of the Doubler magnets on the inside edge. In this supplement we report various cases in which protons strike different parts of the vacuum chamber with various incident angles. The vacuum chamber radius is either from 3.18 cm to 3.31 cm or from 3.68 cm to 3.81 cm. The incident beam angle at which protons strike the vacuum chamber with respect to the perfect beam orbit direction is either  $\pm 0.7$  mrad, or zero (parallel incidence). The positive (negative) angle corresponds to a beam moving outward (inward) in the accelerator radius direction. The incident angle of  $\pm 0.7$  mrad is roughly the maximum beam divergency angle after defocussing quadrupole magnets in the accelerator. The incident angle of 0 mrad corresponds to parallel beam particles along the vacuum chamber side wall. In the parallel case, when protons hit the edge of the vacuum chamber or plug, the probability that some of the protons are scattered out due to Coulomb scattering before nuclear interaction was crudely taken into account.

Figures S1 through S21 show energy density distributions in the Doubler magnet coils as functions of azimuthal angle and coil radius. Table I summarizes some of the interesting properties for all the cases.

The maximum energy deposition at the upstream end and in the shallow radial region of coils (3.81 cm to 4.4 cm) is very sensitive to the vacuum chamber radius. The maximum energy deposition is reduced by a factor of about 5 by changing the chamber inner radius from 3.68 cm to 3.18 cm in all the cases discussed here. The second peaks due to neutral pion production depend strongly upon beam conditions because the vacuum chamber material can absorb

photons from the neutral pion decay.

The energy deposition in the large radial region of coils (4.4 cm to 5.0 cm) is substantially smaller than that in the shallow region. This is due to shielding by the coil material itself.

The vacuum chamber plug suppresses the energy deposition level to 0.01 GeV/(cm<sup>3</sup> interacting proton) or less except for in the upstream end region.

TABLE I

Vacuum Chamber Inner Radius (cm)	Beam Position on Chamber Radius	$\phi$ (rad)	$\theta_{\text{Beam}}$ (mrad)	Maximum Energy Deposition, $\text{GeV}/(\text{cm}^3 \cdot \text{Int. Protons})$				
				Coil Radius: 3.81-4.40 cm		4.40-5.00 cm		
				Upstream ( $\phi$ )	2nd peak* ( $\phi$ )	Upstream ( $\phi$ )	2nd peak* ( $\phi$ )	
S1	3.18	Inside	0	0	0.38 (0)	0.12 ( $\pi$ )	0.19 (0)	0.05 ( $\pi$ )
S2	3.18	Inside	0	+0.7	0.56 (0)	0.03 ( $\pi$ )	0.24 (0)	0.01 ( $\pi$ )
S3	3.18	Middle	0	0	0.85 (0)	0.003 ( $\pi$ )	0.38 (0)	0.002 ( $\pi$ )
S4	3.18	Middle	0	+0.7	0.88 (0)	0.001 ( $\pi$ )	0.41 (0)	0.001 ( $\pi$ )
S5	3.18	Outside	0	0	0.67 (0)	0.01 (0)	0.28 (0)	0.003 (0)
S6	3.68	Inside	0	0	1.9 (0)	0.12 ( $\pi$ )	0.19 (0)	0.05 ( $\pi$ )
S7	3.68	Inside	0	+0.7	2.6 (0)	0.06 ( $\pi$ )	0.30 (0)	0.02 ( $\pi$ )
S8	3.68	Middle	0	+0.7	5.4 (0)	0.001 ( $\pi$ )	0.60 (0)	0.001 ( $\pi$ )
S9	3.18	Inside	$\pi$	0	0.30 ( $\pi$ )	0.015 (0)	0.19 ( $\pi$ )	0.01 (0)
S10	3.18	Inside	$\pi$	-0.7	0.52 ( $\pi$ )	0.006 (0)	0.25 ( $\pi$ )	0.004 (0)
S11	3.18	Middle	$\pi$	0	0.82 ( $\pi$ )	0.001 (0)	0.40 ( $\pi$ )	0.001 (0)
S12	3.18	Middle	$\pi$	-0.7	0.80 ( $\pi$ )	0.001 ( $\pi$ )	0.46 ( $\pi$ )	0.0005 (0)
S13	3.18	Outside	$\pi$	0	0.61 ( $\pi$ )	0.004 ( $\pi$ )	0.25 ( $\pi$ )	0.001 (?)
S14	3.68	Inside	$\pi$	0	2.2 ( $\pi$ )	0.02 (0)	0.19 ( $\pi$ )	0.008 (0)
S15	3.68	Inside	$\pi$	-0.7	2.9 ( $\pi$ )	0.006 (0)	0.30 ( $\pi$ )	0.005 (0)
S16	3.68	Middle	$\pi$	-0.7	5.7 ( $\pi$ )	0.0005 (?)	0.62 ( $\pi$ )	0.0005 (0)
S17	3.68	Outside	$\pi$	-0.7	9.0 ( $\pi$ )	0.0005 (?)	0.62 ( $\pi$ )	0.0005 (?)
S18	Plug**	Inside	0	0	0.09 (0)	0.01 ( $\pi$ )	0.04 (?)	0.005 ( $\pi$ )
S19	Plug	Inside	0	+0.7	0.15 (0)	0.004 ( $\pi$ )	0.07 (?)	0.001 (?)
S20	Plug	Inside	$\pi$	0	0.10 ( $\pi$ )	0.005 (0)	0.05 ( $\pi$ )	0.002 (0)
S21	Plug	Inside	$\pi$	-0.7	0.12 ( $\pi$ )	0.002 (0)	0.06 ( $\pi$ )	0.001 (?)

\* when no apparent second peak appears, the maximum energy deposition at about 10m from the upstream end of the Doubler magnet is quoted.

\*\* the aperture of the vacuum chamber plug hole is 5cm (horizontal) x 3cm (vertical). The space between the coils and the plug aperture is filled with iron.

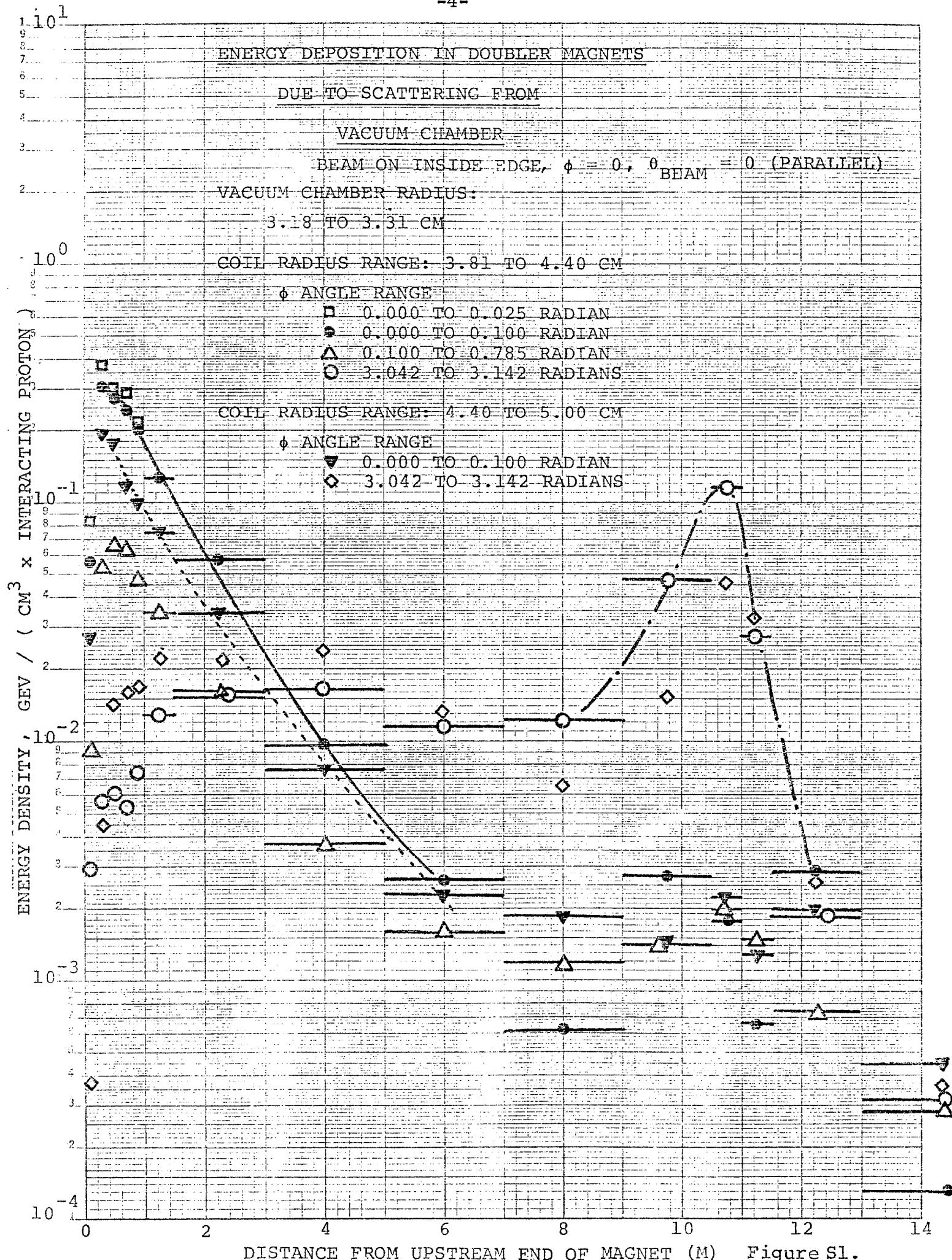


Figure S1.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON INSIDE EDGE

$\phi = 0, \theta = 0.7$  MRAD

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

$\square$  0.000 TO 0.025 RADIANS

$\bullet$  0.000 TO 0.100 RADIANS

$\Delta$  0.100 TO 0.785 RADIANS

$\circ$  3.042 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

$\nabla$  0.000 TO 0.100 RADIANS

$\diamond$  3.042 TO 3.142 RADIANS

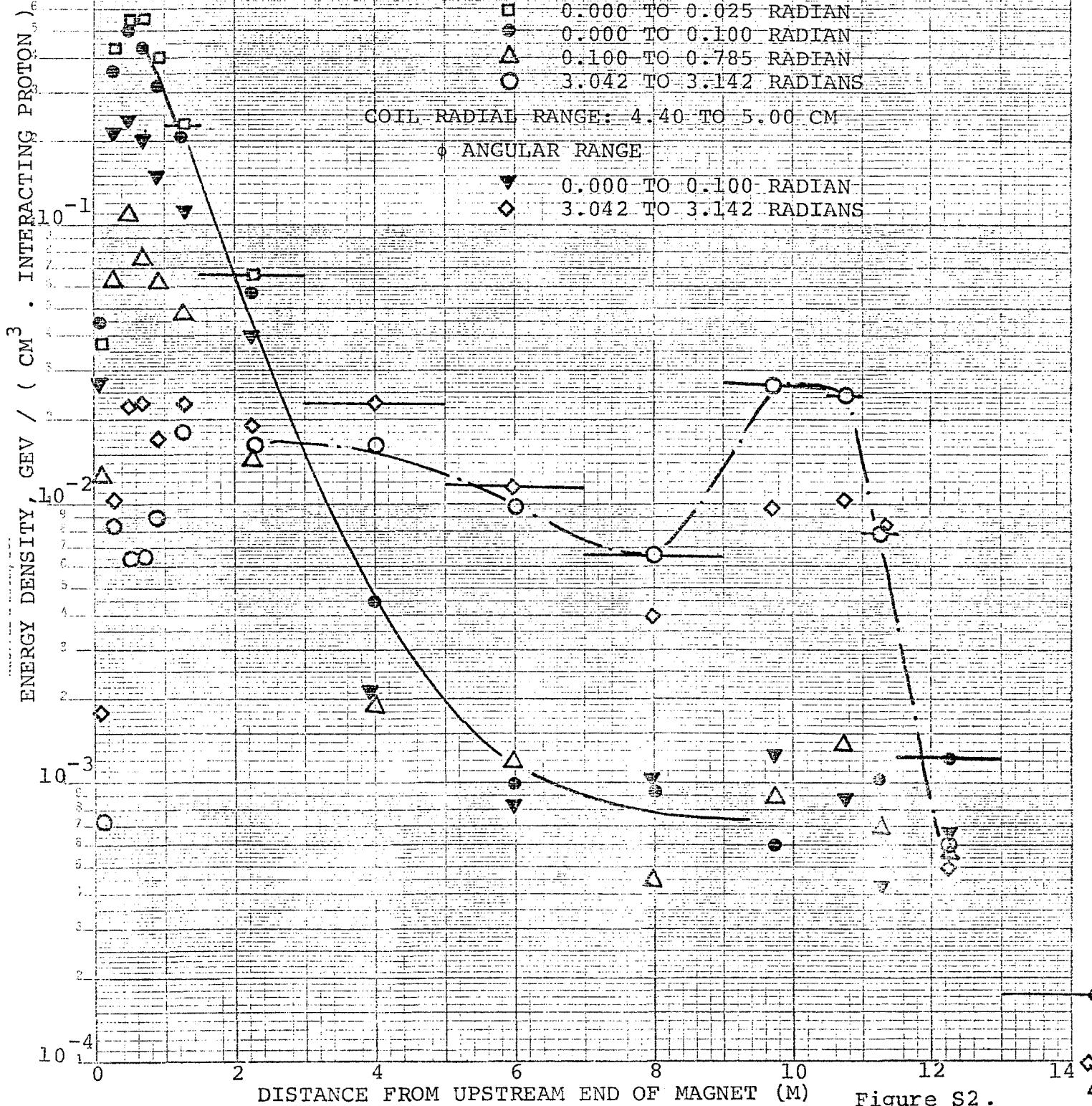


Figure S2.

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON MIDDLE OM CHAMBER WIDTH

$\phi = 0, \theta_{BEAM} = 0$  (PARALLEL)

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

□ 0.000 TO 0.025 RADIANS

● 0.000 TO 0.100 RADIANS

△ 0.100 TO 0.785 RADIANS

○ 3.042 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 0.000 TO 0.100 RADIANS

◇ 3.042 TO 3.142 RADIANS

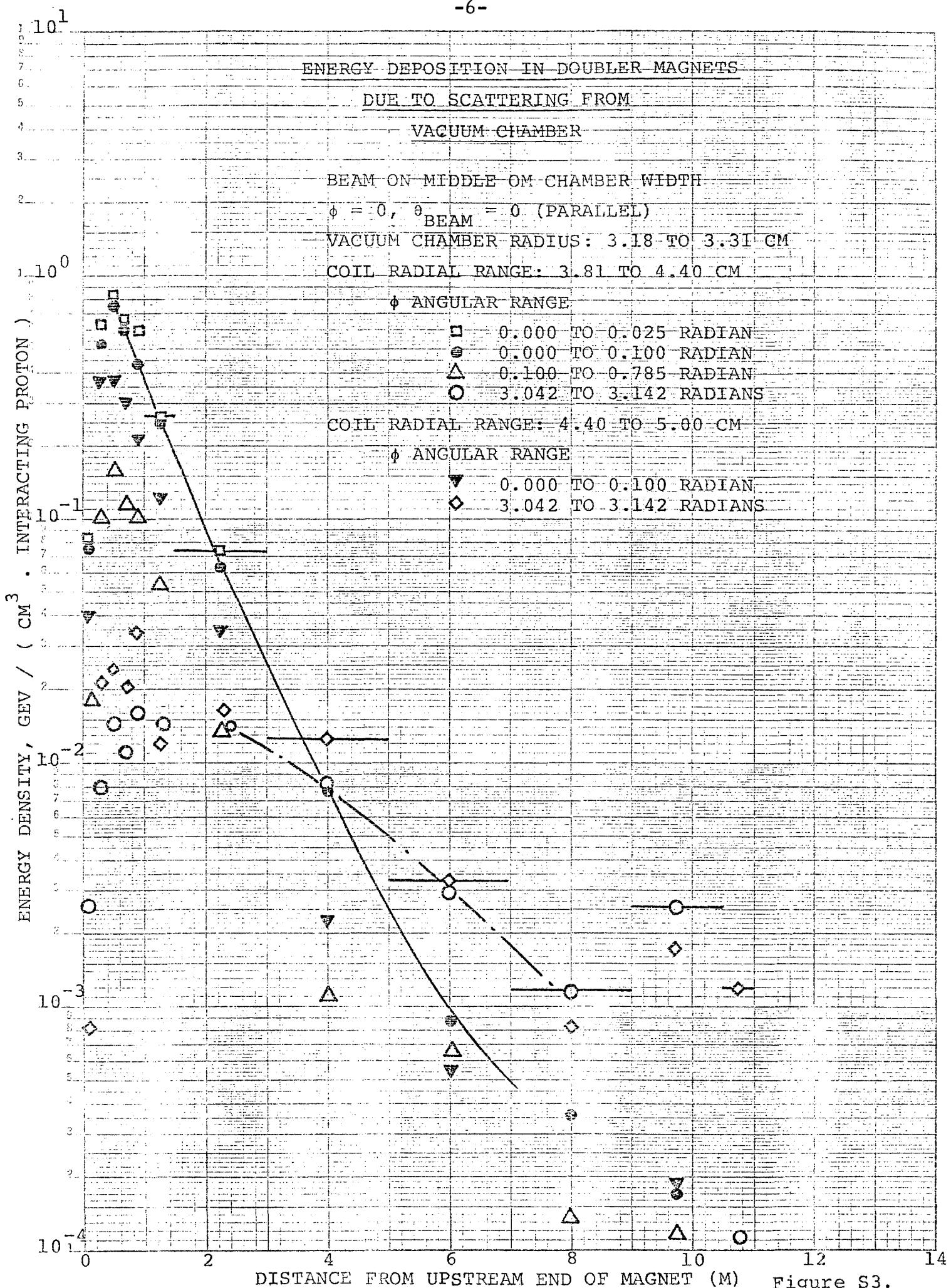


Figure S3.

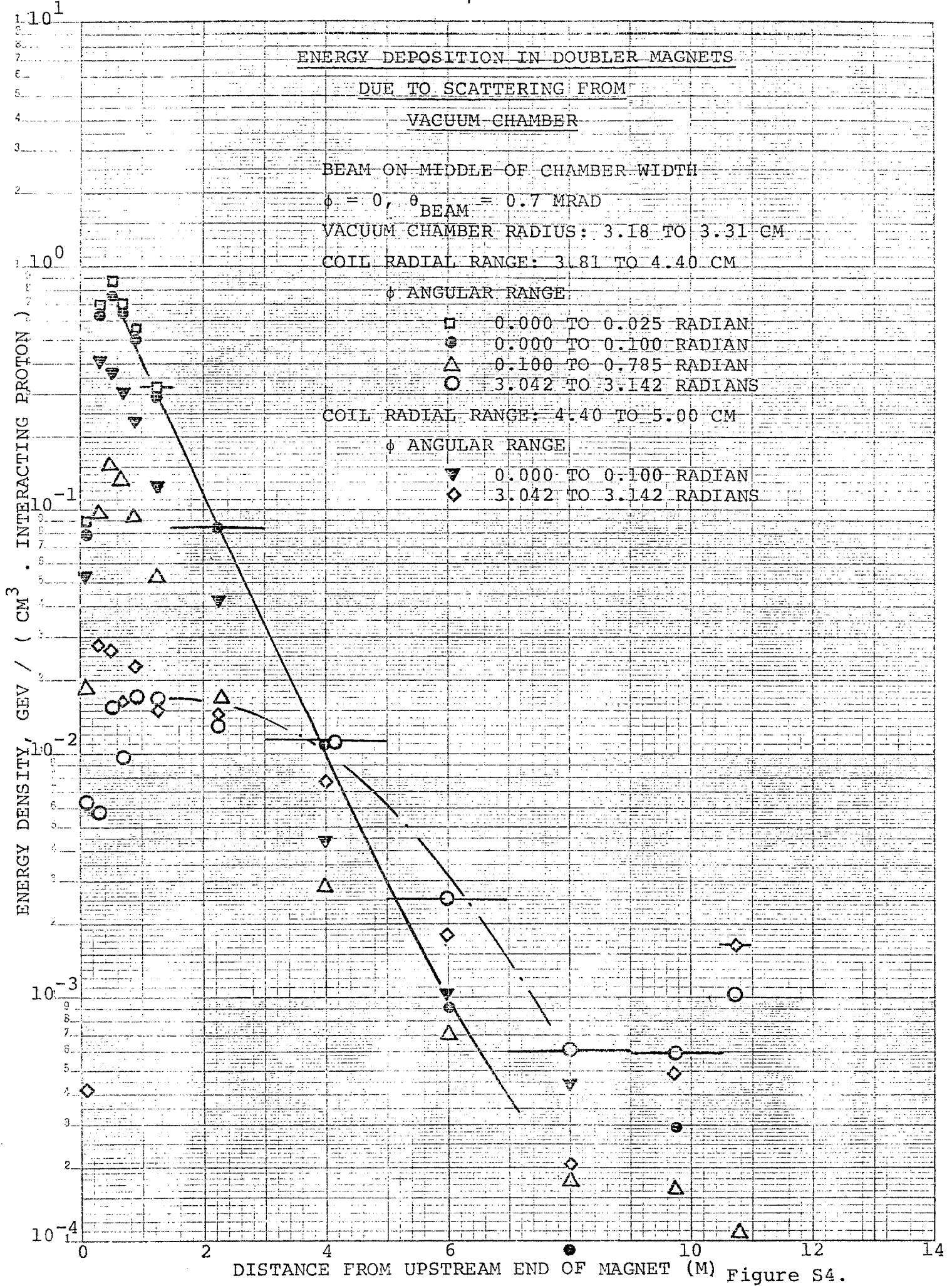


Figure S4.

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON OUTSIDE EDGE

$\phi = 0, \theta_{BEAM} = 0$  (PARALLEL)

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

□ 0.000 TO 0.025 RADIANS

● 0.000 TO 0.100 RADIANS

△ 0.100 TO 0.785 RADIANS

○ 3.042 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 0.000 TO 0.100 RADIANS

◇ 3.042 TO 3.142 RADIANS

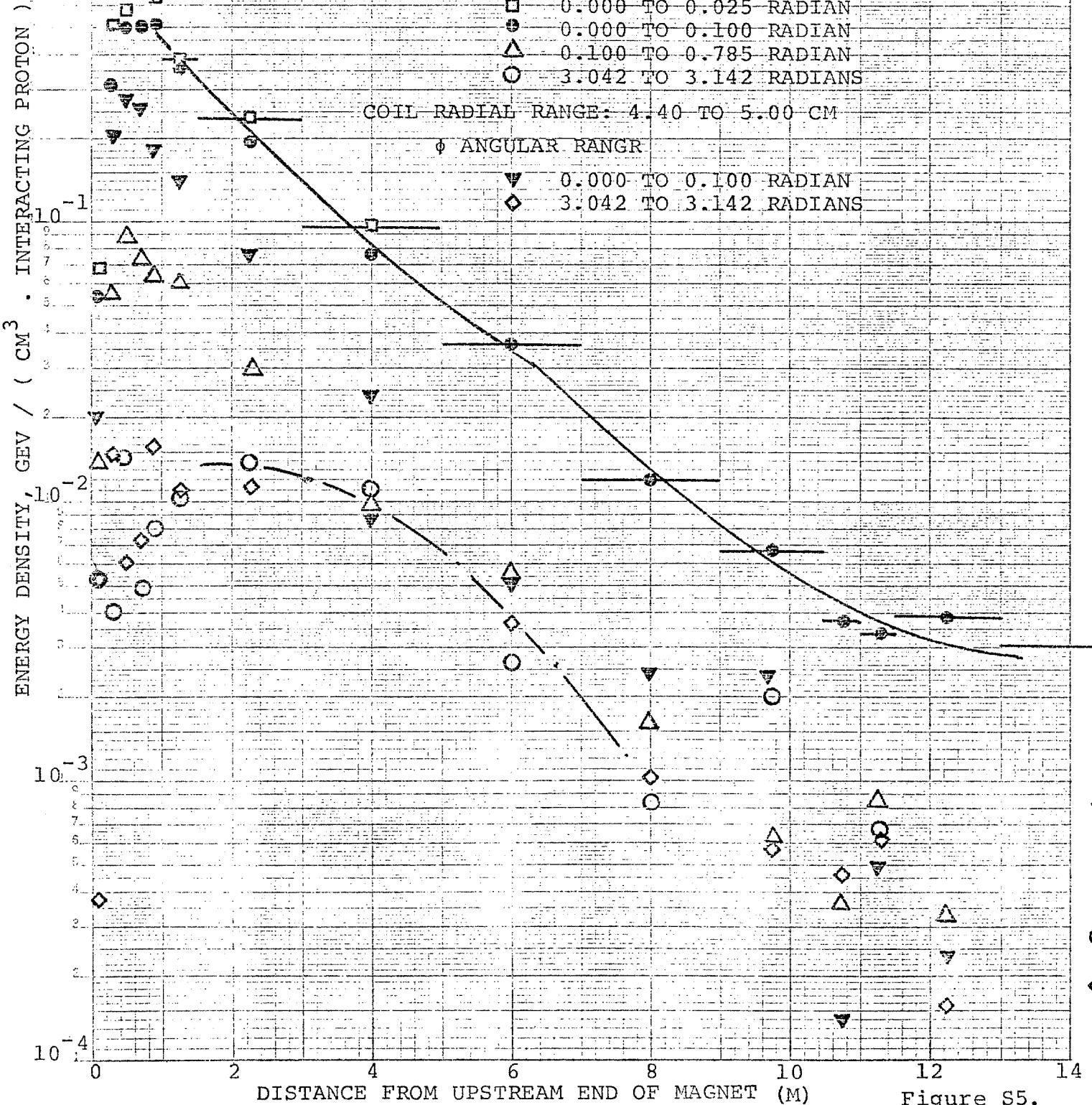
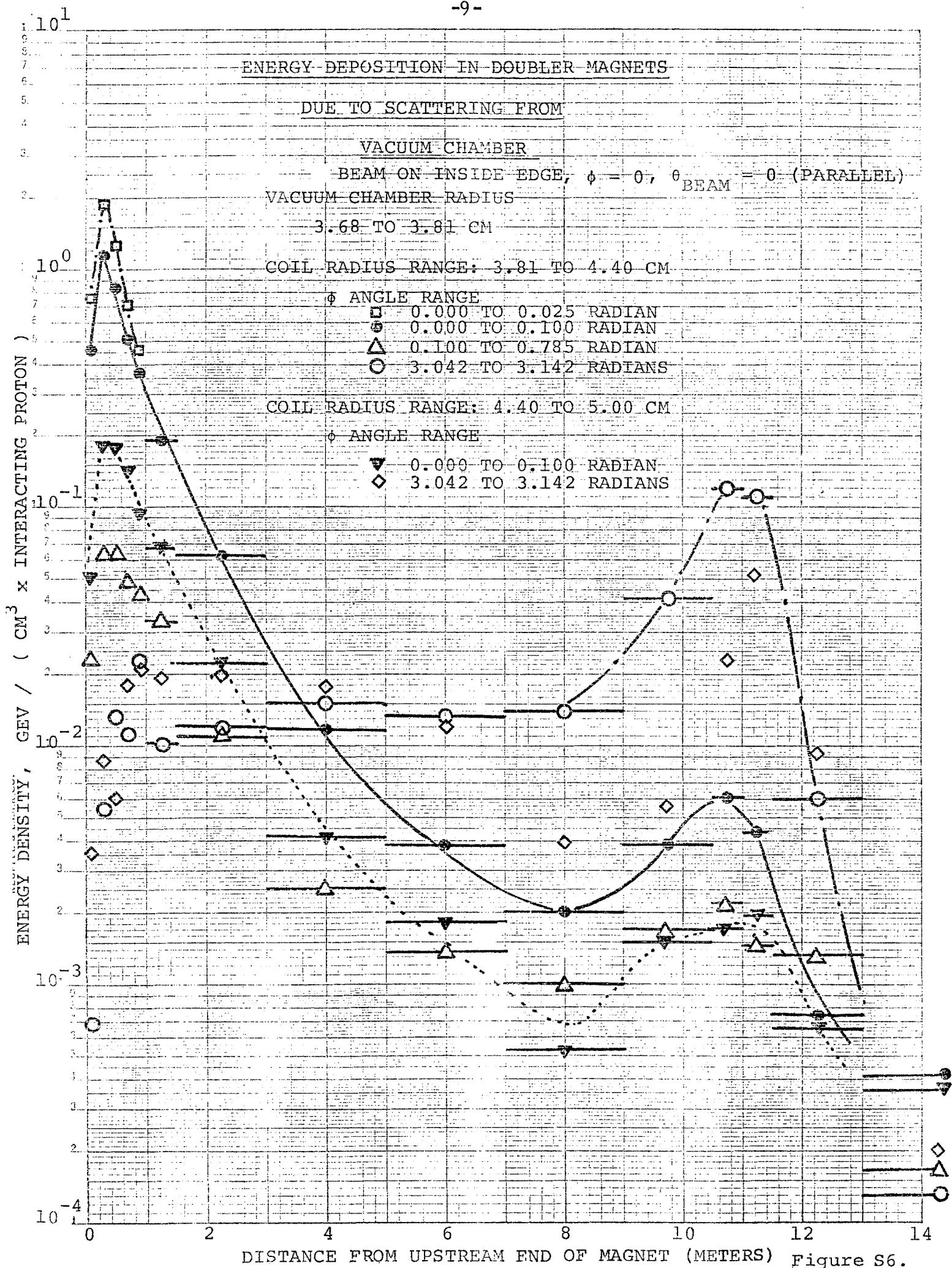


Figure S5.



ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON INSIDE EDGE

$\phi = 0$ , BEAM = 0.7 MRAD

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

□ 0.000 TO 0.025 RADIANS

● 0.000 TO 0.100 RADIANS

△ 0.100 TO 0.785 RADIANS

○ 3.042 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 0.000 TO 0.100 RADIANS

◇ 3.042 TO 3.142 RADIANS

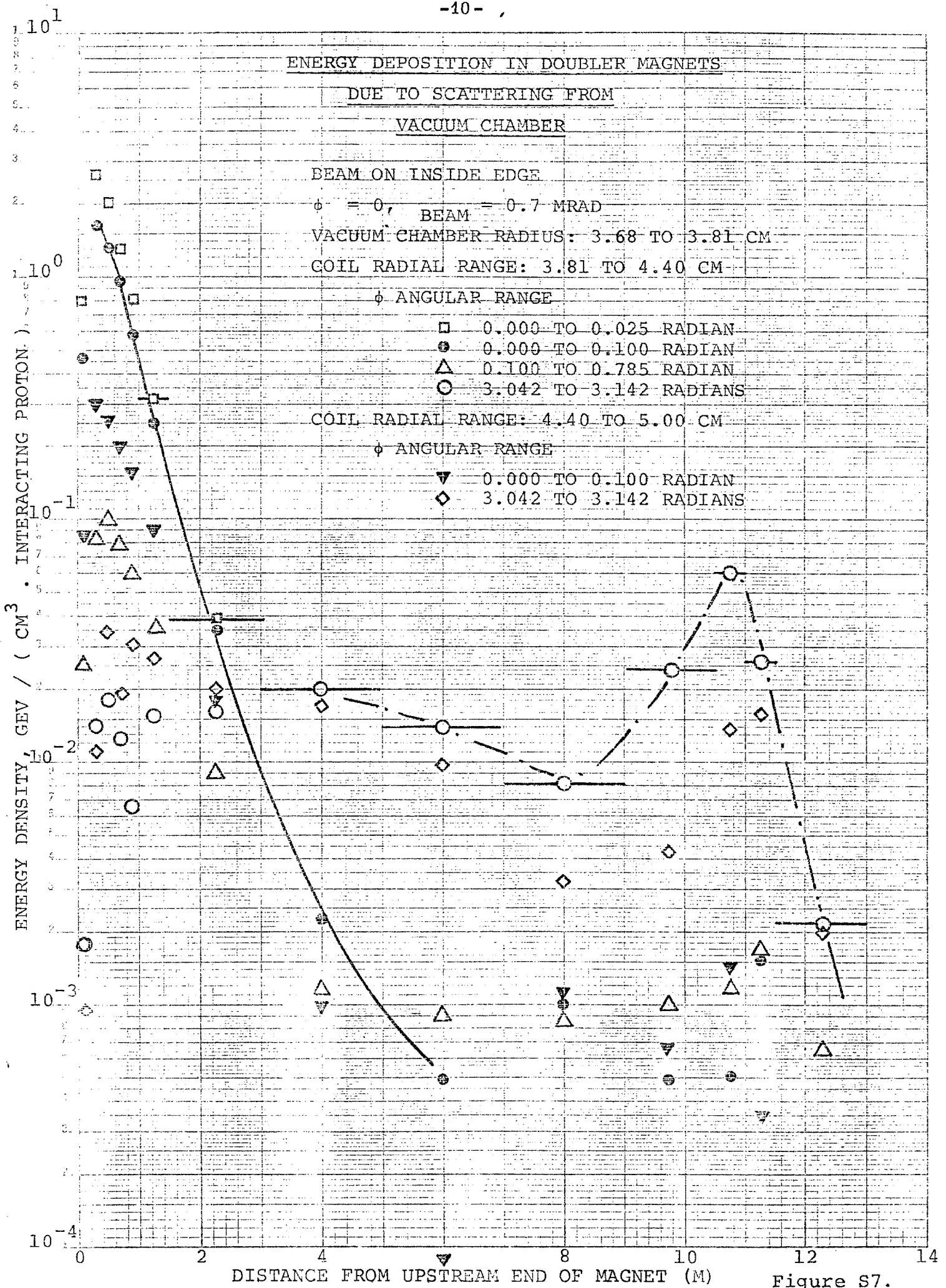


Figure S7.

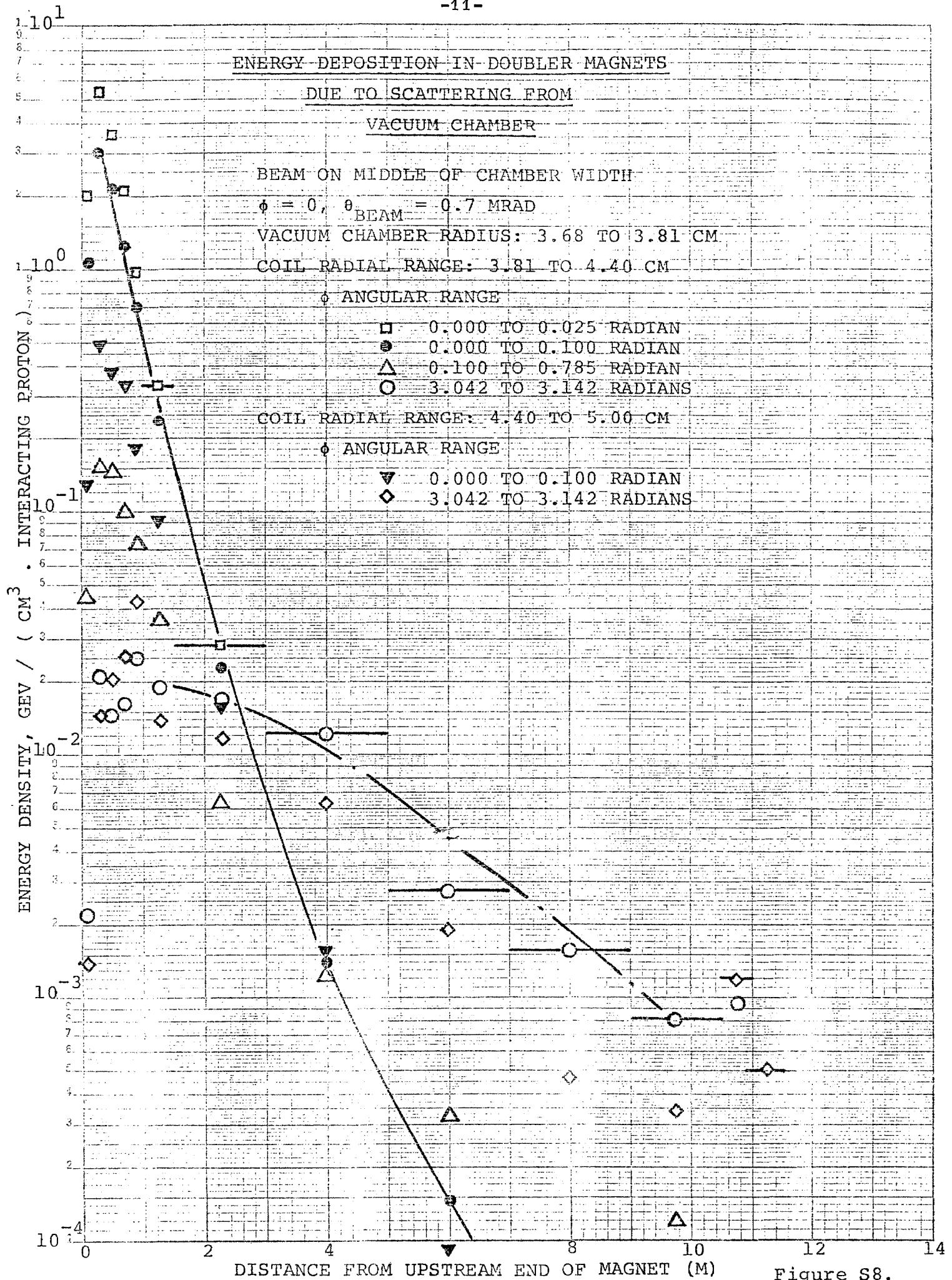


Figure S8.

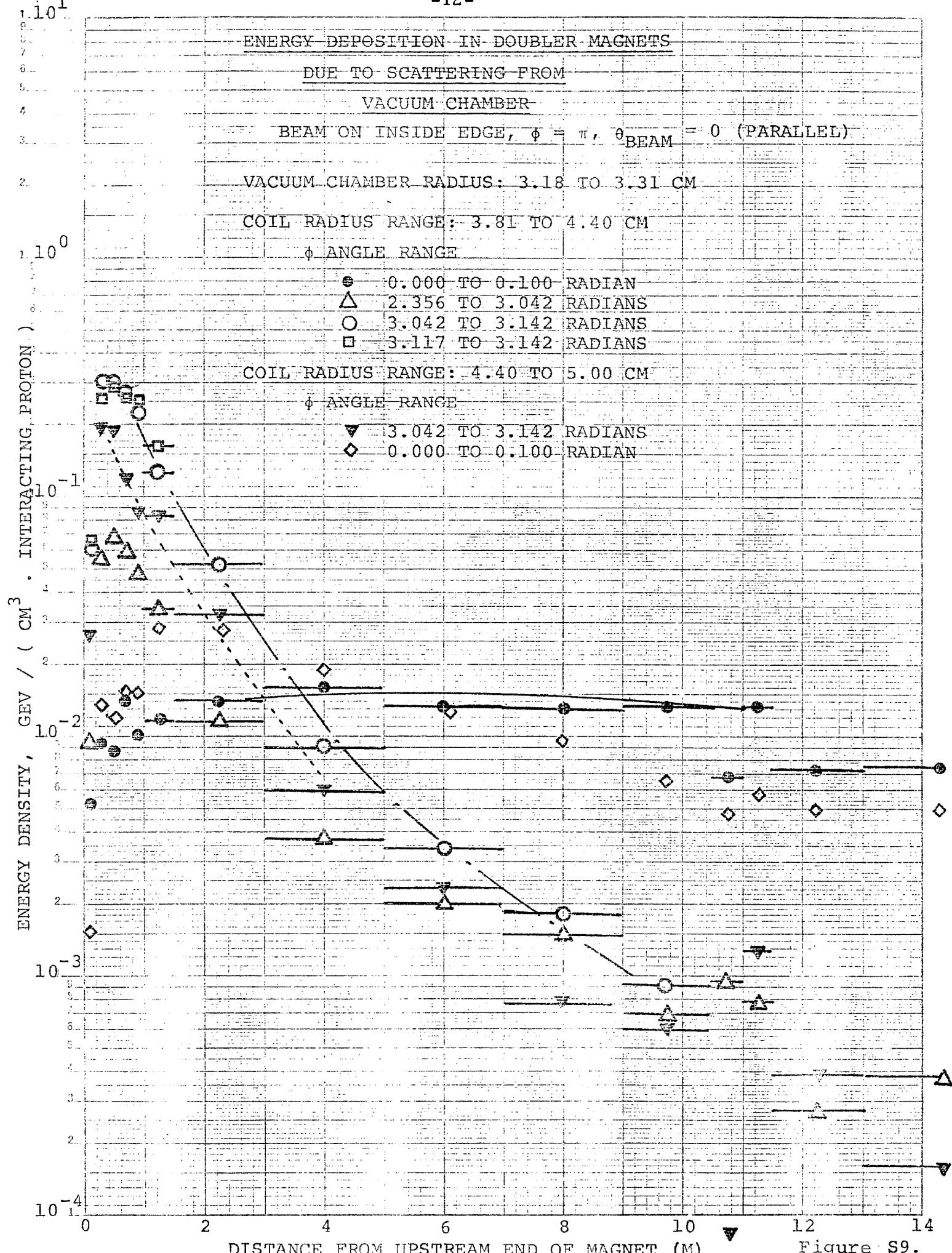


Figure S9.

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON INSIDE EDGE

$\phi = \pi$ ,  $\theta_{\text{BEAM}} = -0.7$  MRAD

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

■ 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

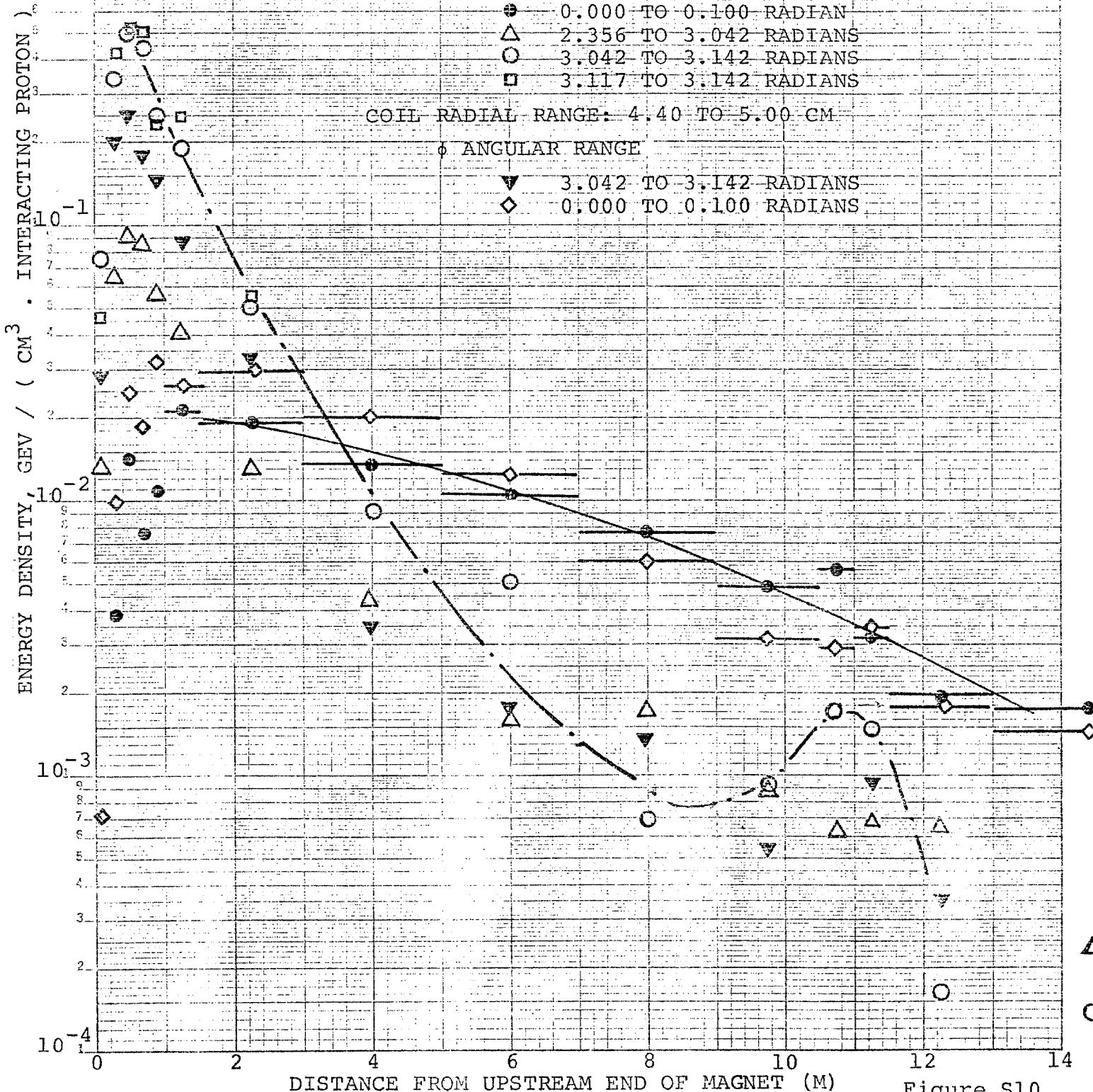


Figure S10.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

BEAM ON MIDDLE OF CHAMBER WIDTH

$\phi = \pi$ ,  $\theta = 0$  (PARALLEL)

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

0.000 TO 0.100 RADIANS

2.356 TO 3.042 RADIANS

3.042 TO 3.142 RADIANS

3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

3.042 TO 3.142 RADIANS

0.000 TO 0.100 RADIANS

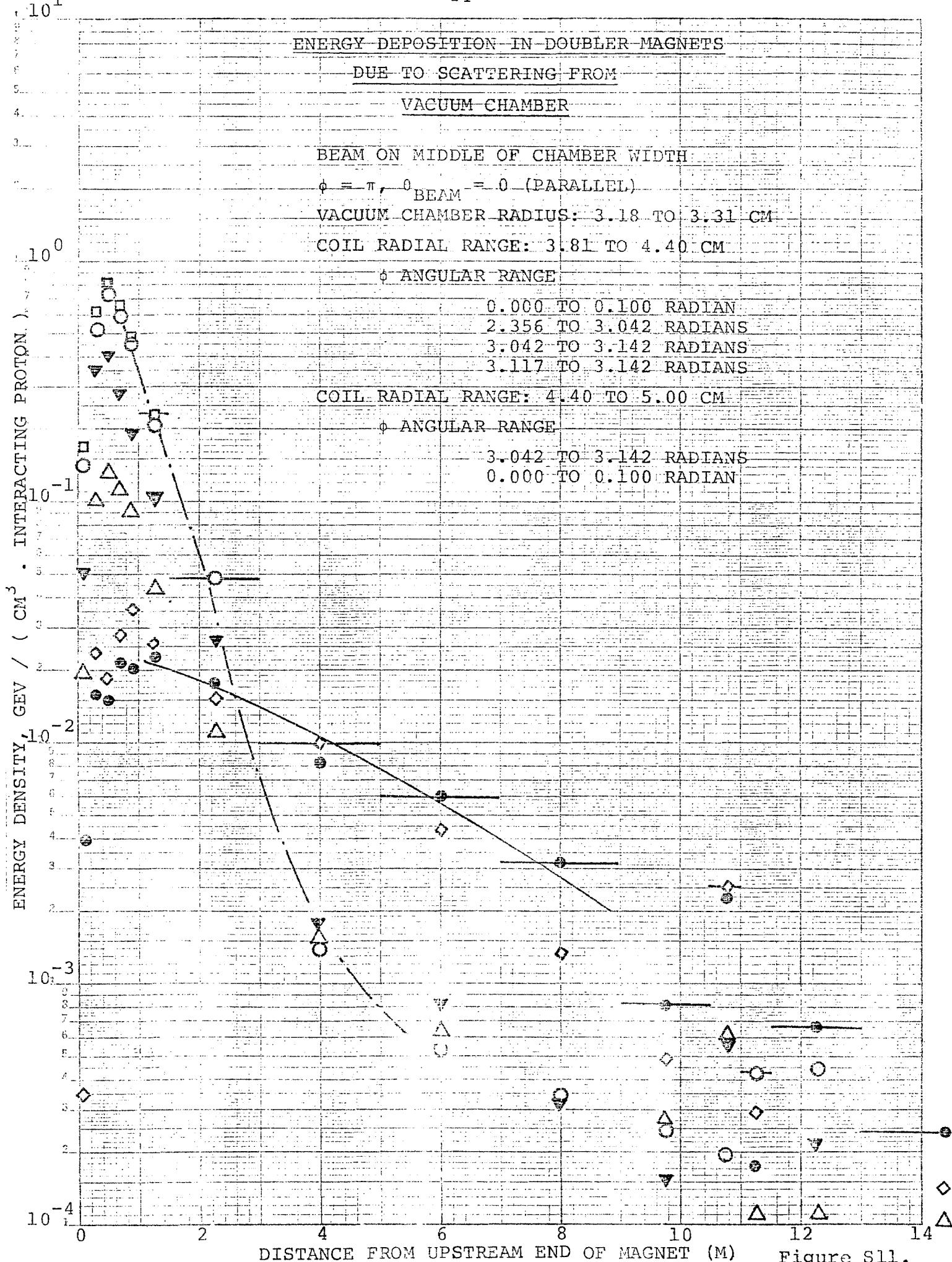


Figure S11.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

BEAM ON MIDDLE OF CHAMBER WIDTH

$\phi = \pi$ ,  $\theta = -0.7$  MRAD

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▽ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

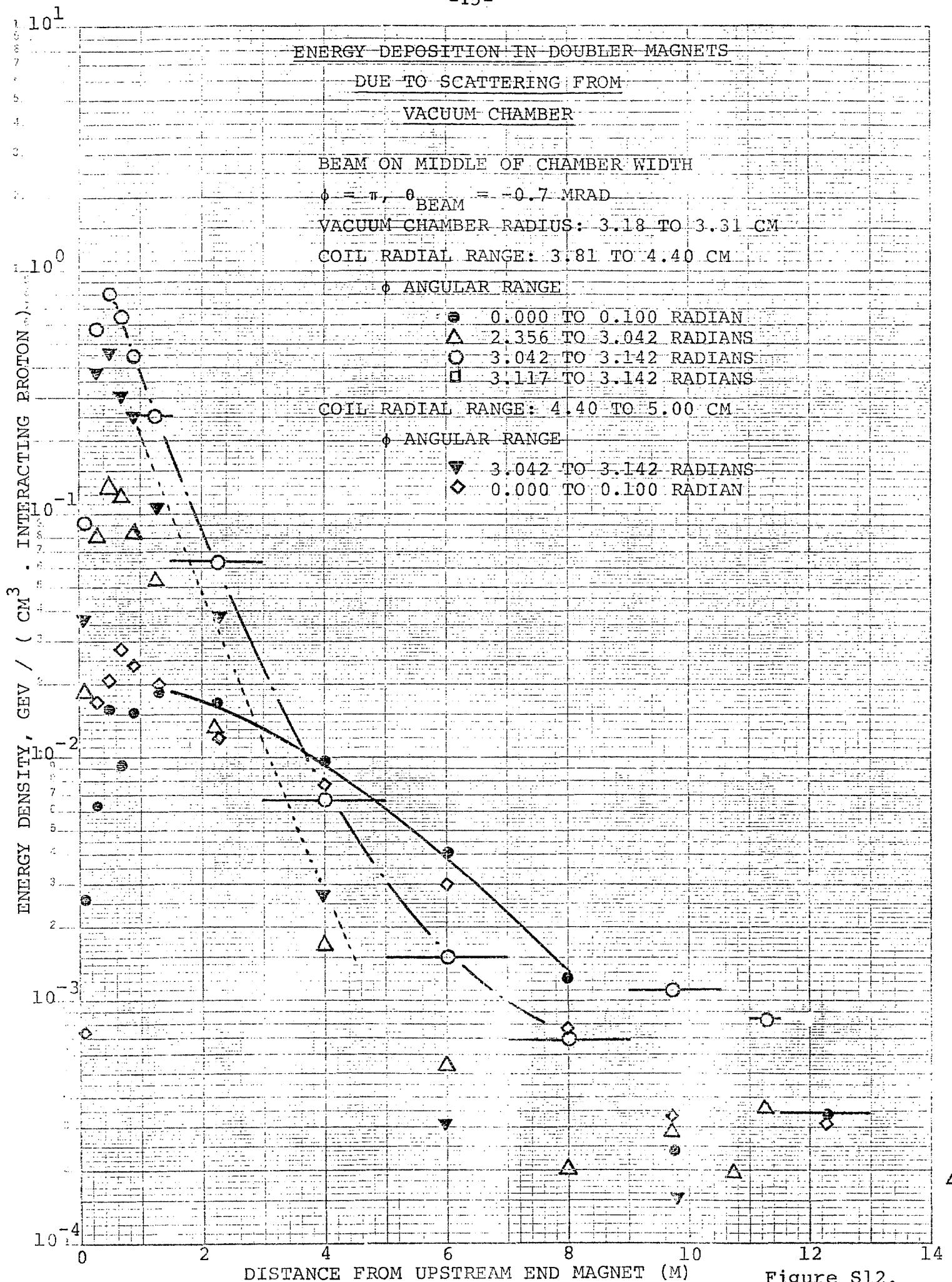


Figure S12.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

BEAM ON OUTSIDE EDGE

$\phi = \pi, \theta = 0$  (PARALLEL)

VACUUM CHAMBER RADIUS: 3.18 TO 3.31 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▽ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

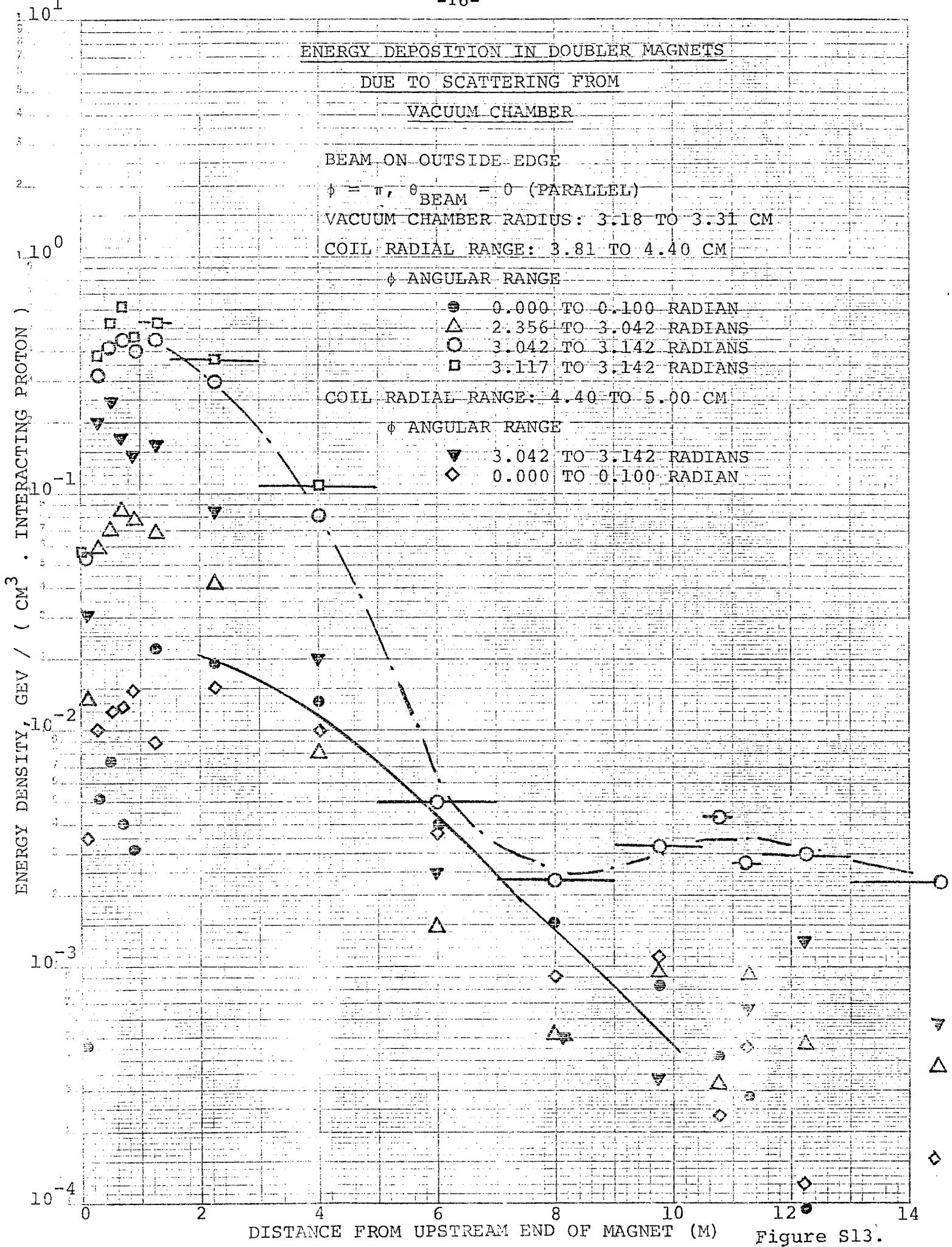


Figure S13.

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON INSIDE EDGE,  $\phi = \pi$ ,  $\theta_{\text{BEAM}} = 0$  (PARALLEL)

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COIL RADIUS RANGE: 3.81 TO 4.40 CM

$\phi$  ANGLE RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIUS RANGE: 4.40 TO 5.00 CM

$\phi$  ANGLE RANGE

▼ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

10<sup>1</sup>

10<sup>0</sup>

10<sup>-1</sup>

10<sup>-2</sup>

10<sup>-3</sup>

10<sup>-4</sup>

ENERGY DENSITY, GEV / ( CM<sup>3</sup> · INTERACTING PROTON )

0

2

4

6

8

10

12

DISTANCE FROM UPSTREAM END OF MAGNET (M)

Figure S14.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

BEAM ON INSIDE EDGE

$\phi = \pi$ ,  $\theta_{\text{BEAM}} = -0.7$  MRAD

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

ENERGY DENSITY, GEV / ( CM<sup>3</sup> · INTERACTING PROTON )

10<sup>1</sup>

10<sup>0</sup>

10<sup>-1</sup>

10<sup>-2</sup>

10<sup>-3</sup>

10<sup>-4</sup>

0

DISTANCE FROM UPSTREAM END OF MAGNET (M)

2 4 6 8 10 12 14

Figure S15.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER

BEAM ON MIDDLE OM CHAMBER WIDTH

$\phi = \pi/0$  BEAM = -0.7 MRAD

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIAN

△ 2.356 TO 3.042 RADIAN

○ 3.042 TO 3.142 RADIAN

□ 3.117 TO 3.142 RADIAN

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 3.042 TO 3.142 RADIAN

◇ 0.000 TO 0.100 RADIAN

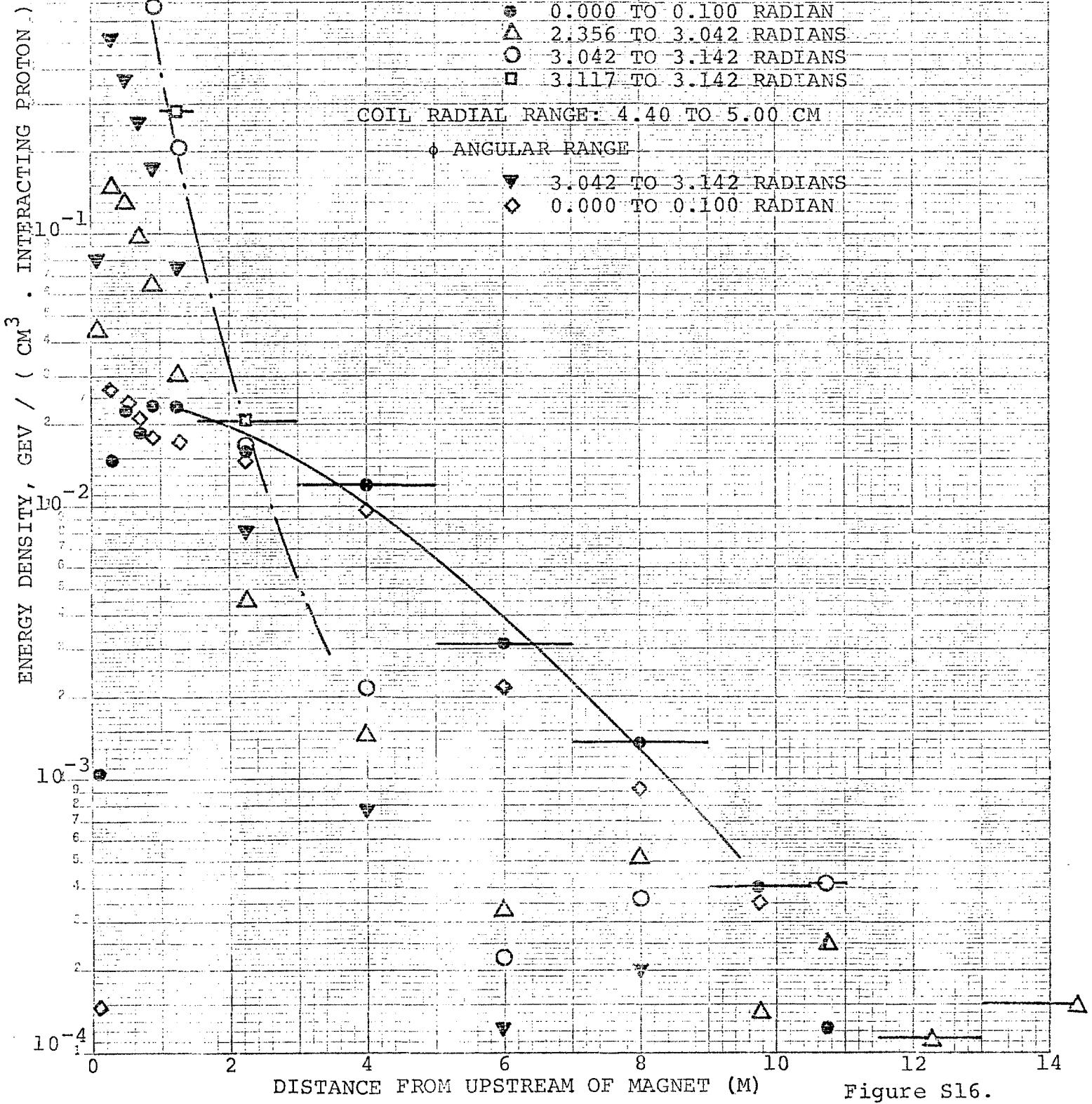


Figure S16.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM  
VACUUM CHAMBER

BEAM ON OUTSIDE EDGE

$\phi = \pi$ ,  $\theta_{\text{BEAM}} = 0.7 \text{ MRAD}$

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▽ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

ENERGY DENSITY, GEV / ( CM<sup>3</sup> · INTERACTING PROTON )

10<sup>1</sup>

10<sup>0</sup>

10<sup>-1</sup>

10<sup>-2</sup>

10<sup>-3</sup>

10<sup>-4</sup>

10<sup>3</sup>

10<sup>1</sup>

10<sup>0</sup>

10<sup>-1</sup>

10<sup>-2</sup>

10<sup>-3</sup>

10<sup>-4</sup>

0 2 4 6 8 10 12 14

DISTANCE FROM UPSTREAM END OF MAGNET (M)

Figure S17.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER PLUG

PLUG: 5 CM (H) X 3 CM (V) APERTURE

BEAM ON INSIDE EDGE,  $\phi = 0$ ,  $\theta = 0$  BEAM

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

□ 0.000 TO 0.025 RADIANS

● 0.000 TO 0.100 RADIANS

△ 0.100 TO 0.785 RADIANS

○ 3.042 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 0.000 TO 0.100 RADIANS

◇ 3.042 TO 3.142 RADIANS

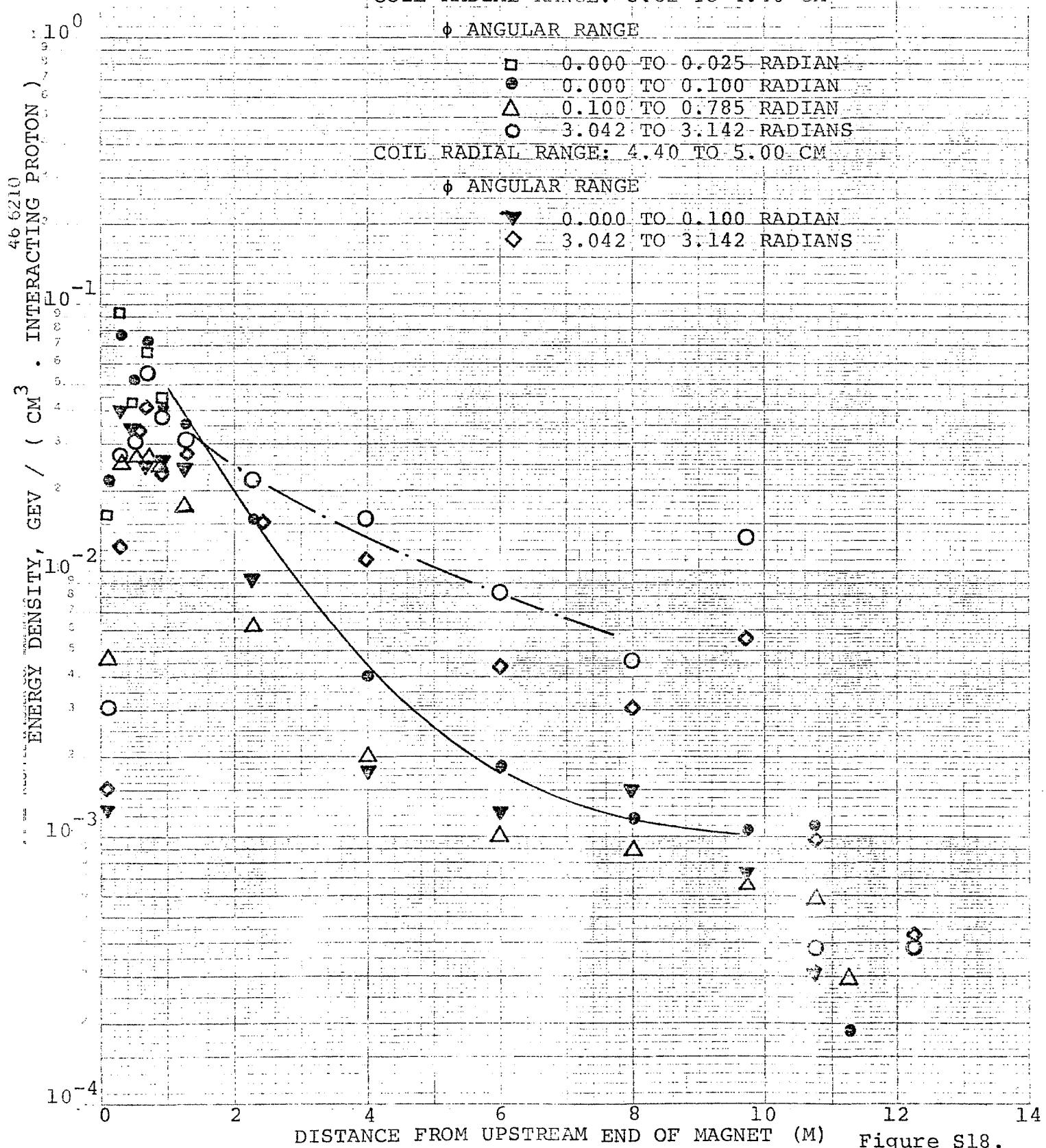


Figure S18.

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER PLUG

PLUG: 5 CM (H) X 3 CM (V) APERTURE

BEAM ON INSIDE EDGE,  $\phi = 0$ ,  $\theta_{BEAM} = 0.7$  MRAD

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

□ 0.000 TO 0.025 RADIANS

◐ 0.000 TO 0.100 RADIANS

△ 0.100 TO 0.785 RADIANS

○ 3.042 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 0.000 TO 0.100 RADIANS

◇ 3.042 TO 3.142 RADIANS

10<sup>1</sup>

10<sup>0</sup>

10<sup>-1</sup>

10<sup>-2</sup>

10<sup>-3</sup>

10<sup>-4</sup>

ENERGY DENSITY, GEV / ( CM<sup>3</sup> . INTERACTING PROTON )

◊

DISTANCE FROM UPSTREAM END OF MAGNET (M)

0 2 4 6 8 10 12 14

Figure S19.

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

VACUUM CHAMBER PLUG

PLUG: 5 CM (H)  $\times$  3 CM (V) APERTURE

BEAM ON INSIDE EDGE,  $\phi = \pi$ ,  $\theta_{BEAM} = 0$

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▼ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

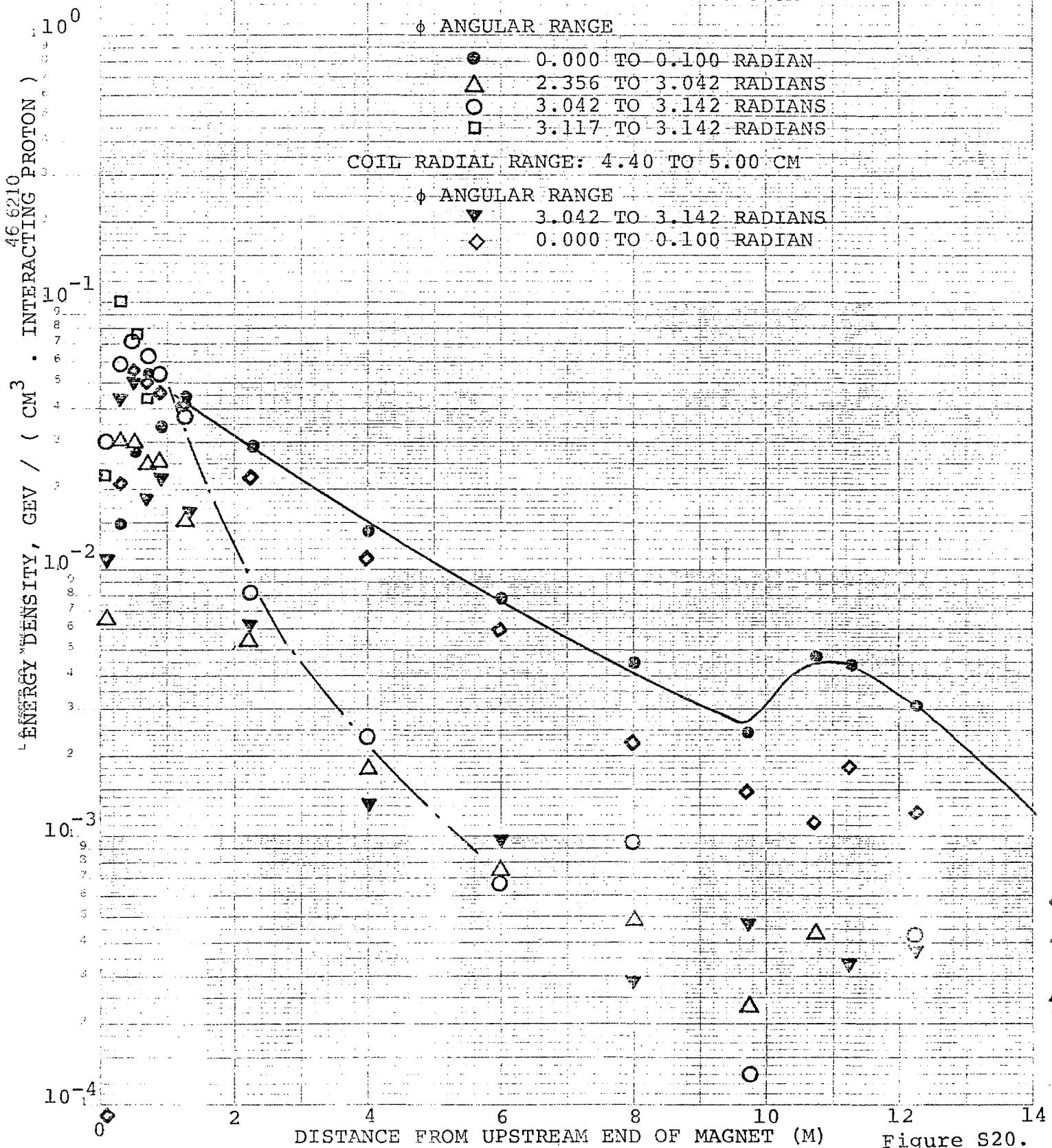


Figure S20.

ENERGY DEPOSITION IN DOUBLER MAGNETS  
DUE TO SCATTERING FROM  
VACUUM CHAMBER PLUG

PLUG: 5 CM (H) x 3 CM (V) APERTURE

BEAM ON INSIDE EDGE,  $\phi = \pi$ ,  $\theta_{\text{BEAM}} = -0.7 \text{ MRAD}$

COIL RADIAL RANGE: 3.81 TO 4.40 CM

$\phi$  ANGULAR RANGE

● 0.000 TO 0.100 RADIANS

△ 2.356 TO 3.042 RADIANS

○ 3.042 TO 3.142 RADIANS

□ 3.117 TO 3.142 RADIANS

COIL RADIAL RANGE: 4.40 TO 5.00 CM

$\phi$  ANGULAR RANGE

▽ 3.042 TO 3.142 RADIANS

◇ 0.000 TO 0.100 RADIANS

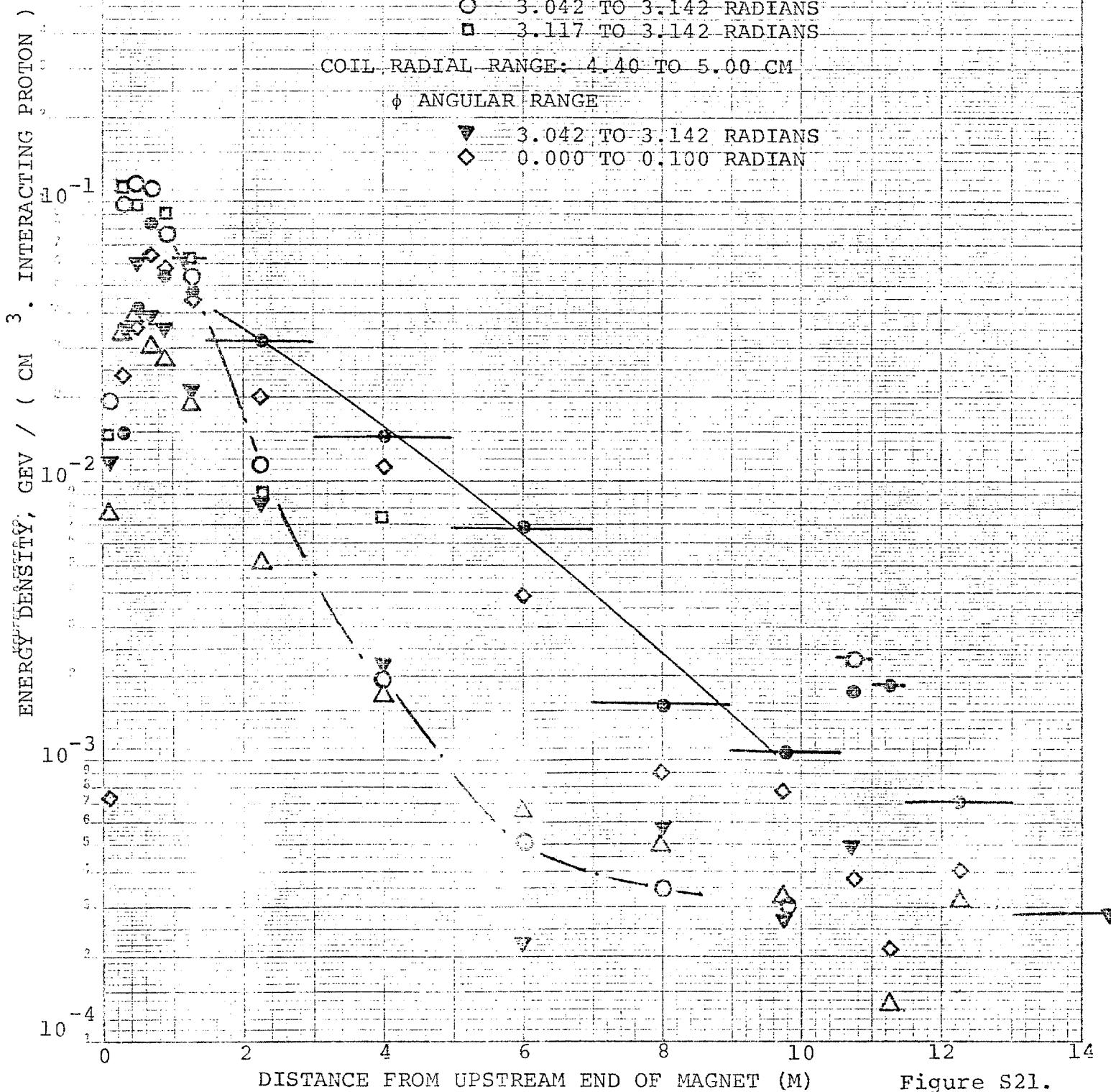


Figure S21.

SUPPLEMENT TO UPC 40

Figures S22 through S25 show energy density distributions in the Doubler magnet coils due to scattering from a beam scraper at a medium straight section. The iron beam scraper with an aperture of 4.5 cm(horizontal) by 2 cm(vertical) and 2 m in length is placed 12 m upstream of the Doubler magnets. The incident proton energy is 1000 GeV. The beam strikes an inside edge of the scraper with an incident angle of 0 mrad. The  $\phi$  angle is  $\pi$  in Figures S22, S23, and S24 and 0 in Figure S25. In Figure S22 beam collimator and vacuum chamber plug are not present. The vacuum chamber radius is 3.68 to 3.81 cm. The vacuum chamber plug with a 5 cm(H) by 3 cm(V) aperture is added in Figure S23. In Figures S24 and S25 a 2 m long iron collimator with a 5 cm(H) by 3 cm(V) aperture is placed immediately upstream of the Doubler magnets. The plug and collimator reduce the peak energy density to about  $6 \times 10^{-3}$  GeV/(cm<sup>3</sup>.interacting proton) from  $4 \times 10^{-2}$  GeV/(cm<sup>3</sup>.interacting proton).

Figures S26 through S33 show energy density distributions due to scattering from the extraction wire septum for various configurations of the beam bump arrangement at the long straight section. Four B-2 type bending magnets are used for the beam bump of 3 mrad as shown in Figure 6. The incident proton energy is 1000 GeV. A string of four quadrupole magnets which are to be placed upstream of the dipole string are omitted in the present study. Figure S26 is essentially the same as Figure 3 in which beam bump, vacuum chamber plug and collimator are not present. A small discrepancy at the upstream end is due to different

binnings in the two cases.

Table II summarizes the conditions and peak energy densities.

Table II. Peak energy densities due to scattering from the extraction electrostatic wire septum for various configurations shown in Figures S26 through S33.

Figure	Plug	Beam Bump	Collimator	Peak Energy Density GeV/(cm <sup>3</sup> · incid. prot.)	
				First	Second
S26	No	No	No	$4 \times 10^{-3}$	$2 \times 10^{-3}$
S27	Yes	No	Yes	$9 \times 10^{-4}$	$6 \times 10^{-4}$
S28	No	4M, H-In	-	-	$9 \times 10^{-4}$
S29	No	3M, H-In	-	-	$4 \times 10^{-4}$
S30	Yes	3M, H-In	-	$2 \times 10^{-4}$	$5 \times 10^{-5}$
S31	Yes	3M, H-Out	-	$2 \times 10^{-4}$	$6 \times 10^{-5}$
S32	No	3M, V	-	-	$8 \times 10^{-5}$
S33	Yes	3M, V	-	$5 \times 10^{-5}$	$4 \times 10^{-5}$

The 2 m long iron collimator with an aperture of 5 cm(H) by 3 cm(V) is used only in Figure S27. The vacuum chamber plug is made of iron and has an aperture of 5 cm(H) by 3 cm(V). The beam bumps, 4M and 3M correspond to the arrangements in which the electrostatic wire septum is placed upstream and downstream of the first bump magnet, respectively. The H-In and H-Out are the horizontal inward and outward bumps and the V is the vertical beam bump. The horizontal inward bump arrangement is slightly better than the horizontal outward bump when the vacuum chamber plug is not used, but

they are essentially the same when the plug is used. In general, the vertical bump gives much lower peak energy density in the Doubler magnet coils than the horizontal bump because the narrow vertical aperture of the bump magnets is more efficient to absorb off-momentum secondary particles scattered from septum wires.

The peak energy density of  $5 \times 10^{-5}$  GeV/(cm<sup>3</sup>.incident proton) achieved by the vertical bump (Figure S33) gives the limits for extraction due to scattering from the electrostatic wire septum of  $4.5 \times 10^{14}$  protons/sec for slow extraction and  $1.2 \times 10^{14}$  protons/pulse of 1 msec. The extraction inefficiency was assumed to be 2.5 %, i.e. 2.5 % of protons strike septum wires.

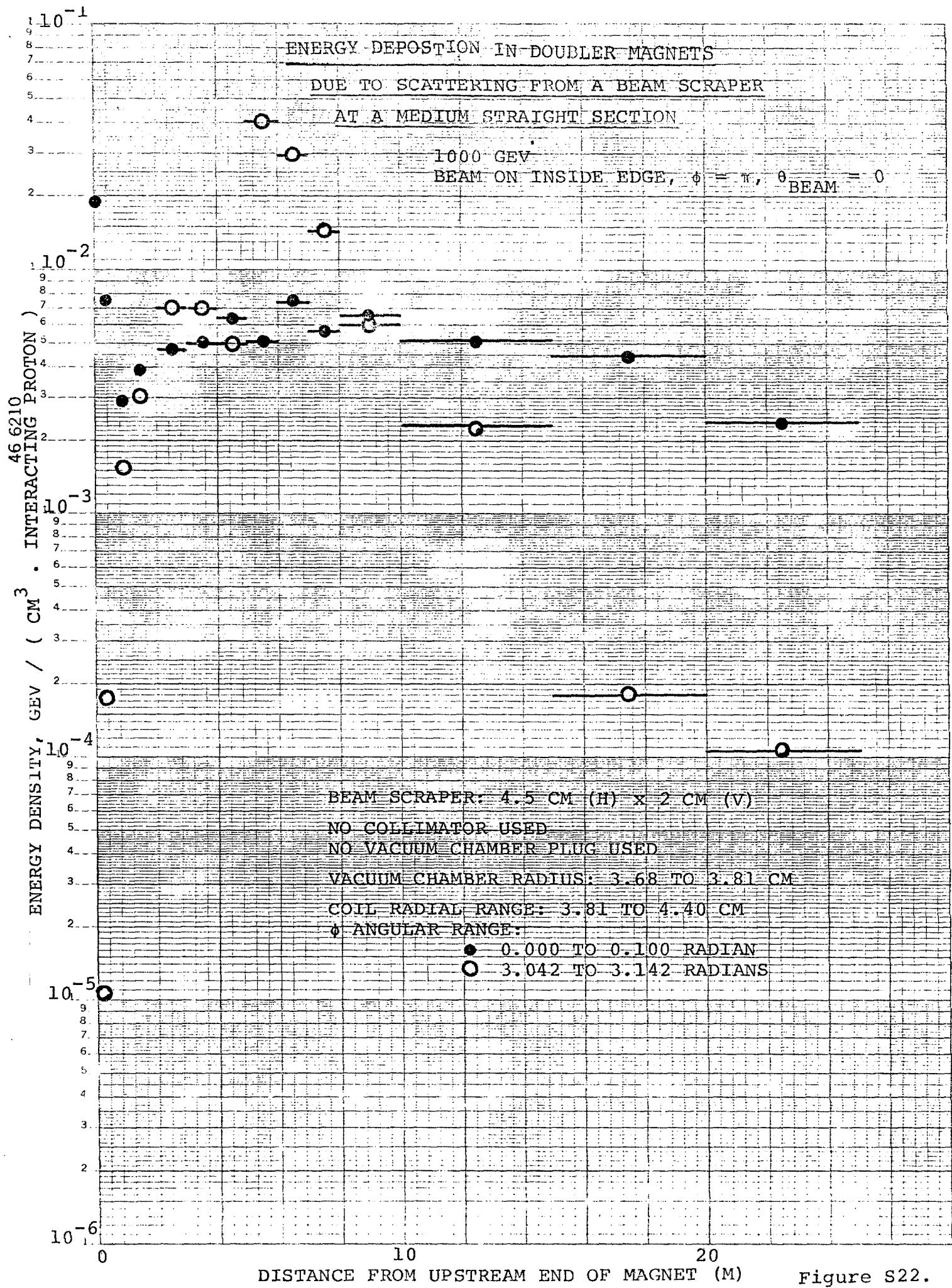


Figure S22.

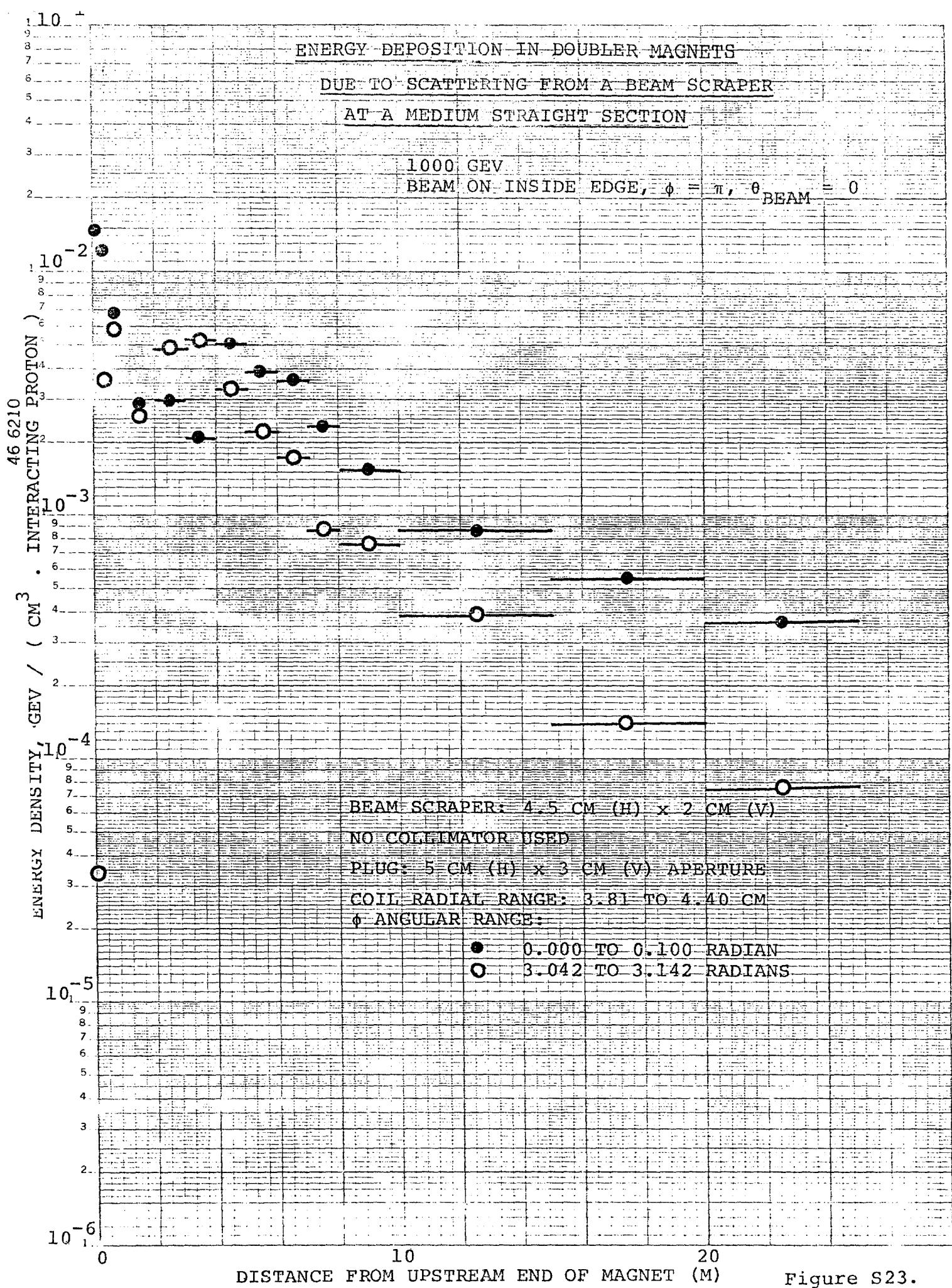


Figure S23.

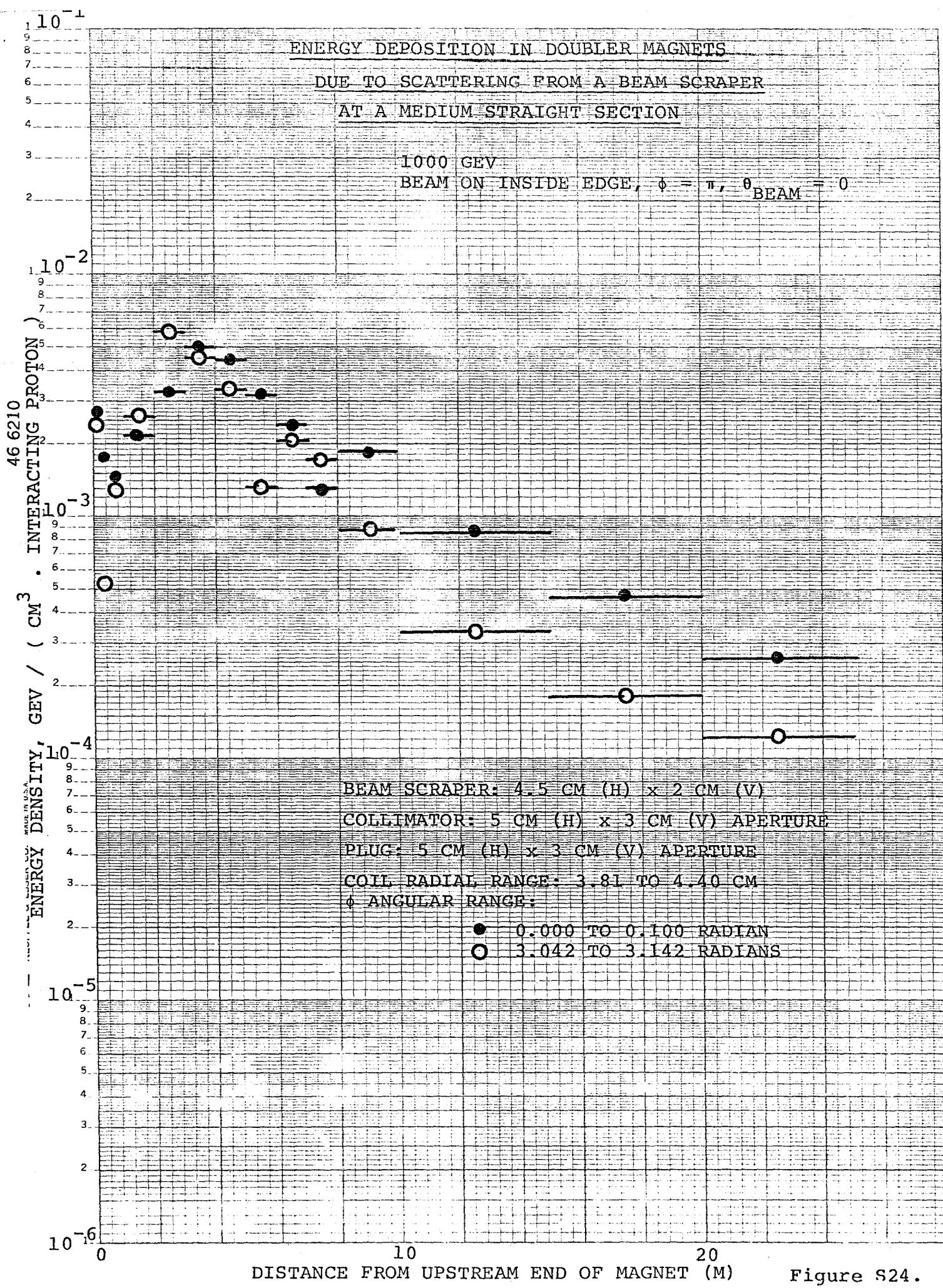


Figure S24.

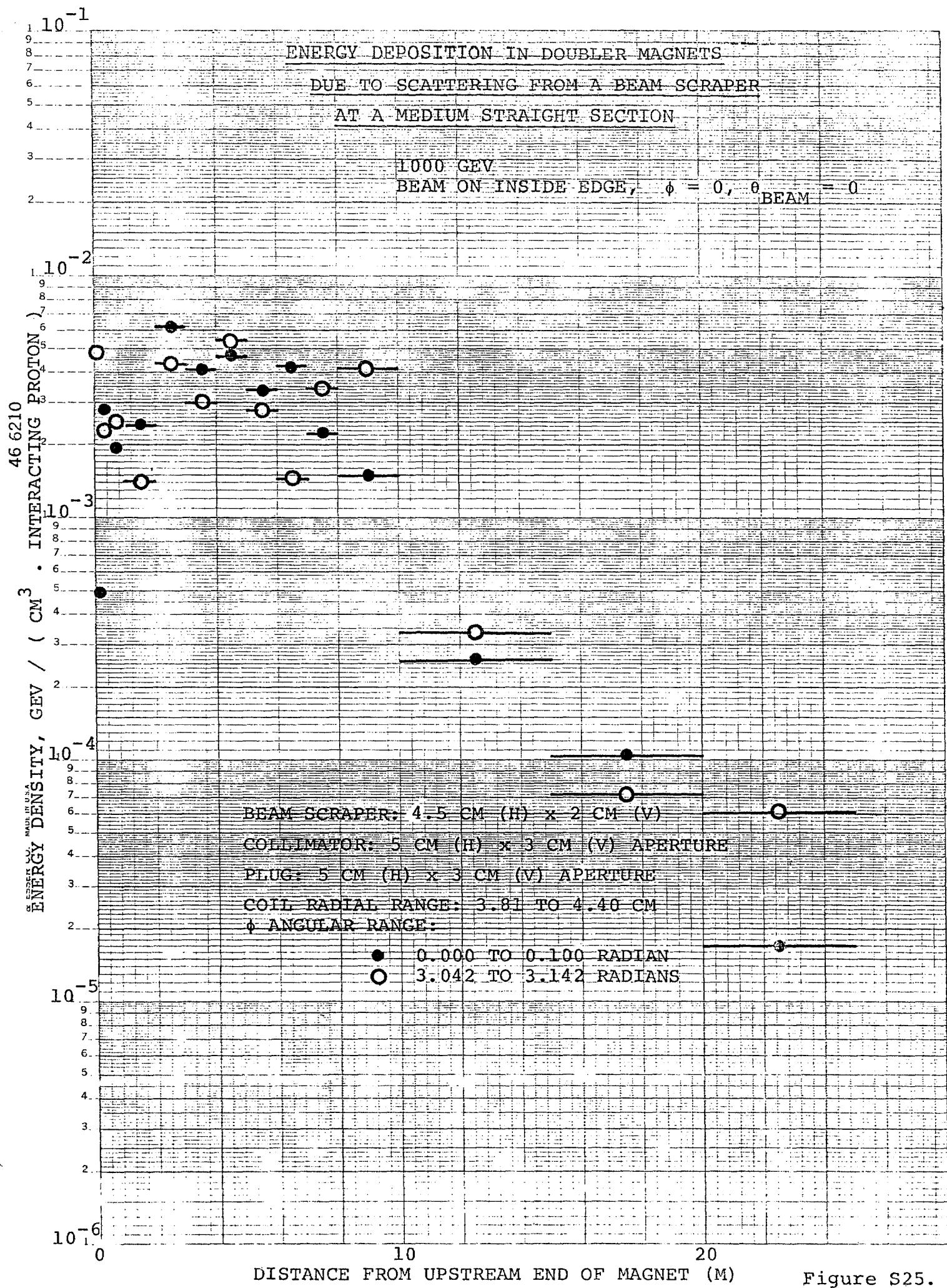


Figure S25.

1  $10^{-2}$   
8  
7  
6  
5

ENERGY DEPOSITION IN DOUBLER MAGNETS

DUE TO SCATTERING FROM

ES WIRE SEPTUM

1000 GEV

NO BEAM BUMP USED

NO PLUG USED

NO COLLIMATOR USED

46 6210  
INCIDENT PROTON

$10^{-4}$

$10^{-3}$

$10^{-2}$

$10^{-1}$

$10^0$

$10^1$

$10^2$

$10^3$

$10^4$

$10^5$

$10^6$

$10^7$

$10^8$

$10^9$

$10^{10}$

$10^{11}$

$10^{12}$

$10^{13}$

$10^{14}$

$10^{15}$

$10^{16}$

$10^{17}$

$10^{18}$

$10^{19}$

$10^{20}$

$10^{21}$

$10^{22}$

$10^{23}$

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$10^{26}$

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$10^{164}$

$10^{165}$

$10^{166}$

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$10^{169}$

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$10^{171}$

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$10^{192}$

$10^{193}$

$10^{194}$

$10^{195}$

$10^{196}$

$10^{197}$

$10^{198}$

$10^{199}$

$10^{200}$

SEMI-LOGARITHMIC 5 CYCLES X 70 DIVISIONS  
KEUFFEL & ESSER CO. MADE IN U.S.A.  
ENERGY DENSITY, GEV / ( CM<sup>3</sup> )

VACUUM CHAMBER RADIUS: 3.68 TO 3.81 CM

COLL RADIAL RANGE: 3.81 TO 4.40 CM

Φ ANGULAR RANGE: 0.000 TO 0.100 RADIANS

0.000 TO 0.100 RADIANS  
3.042 TO 3.142 RADFANS

10<sup>-7</sup>

10<sup>-6</sup>

10<sup>-5</sup>

10<sup>-4</sup>

10<sup>-3</sup>

10<sup>-2</sup>

10<sup>-1</sup>

10<sup>0</sup>

10<sup>1</sup>

10<sup>2</sup>

10<sup>3</sup>

10<sup>4</sup>

10<sup>5</sup>

10<sup>6</sup>

10<sup>7</sup>

10<sup>8</sup>

10<sup>9</sup>

10<sup>10</sup>

10<sup>11</sup>

10<sup>12</sup>

10<sup>13</sup>

10<sup>14</sup>

10<sup>15</sup>

10<sup>16</sup>

10<sup>17</sup>

10<sup>18</sup>

10<sup>19</sup>

10<sup>20</sup>

10<sup>-7</sup>

10<sup>-6</sup>

10<sup>-5</sup>

10<sup>-4</sup>

10<sup>-3</sup>

10<sup>-2</sup>

10<sup>-1</sup>

10<sup>0</sup>

10<sup>1</sup>

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10<sup>4</sup>

10<sup>5</sup>

10<sup>6</sup>

10

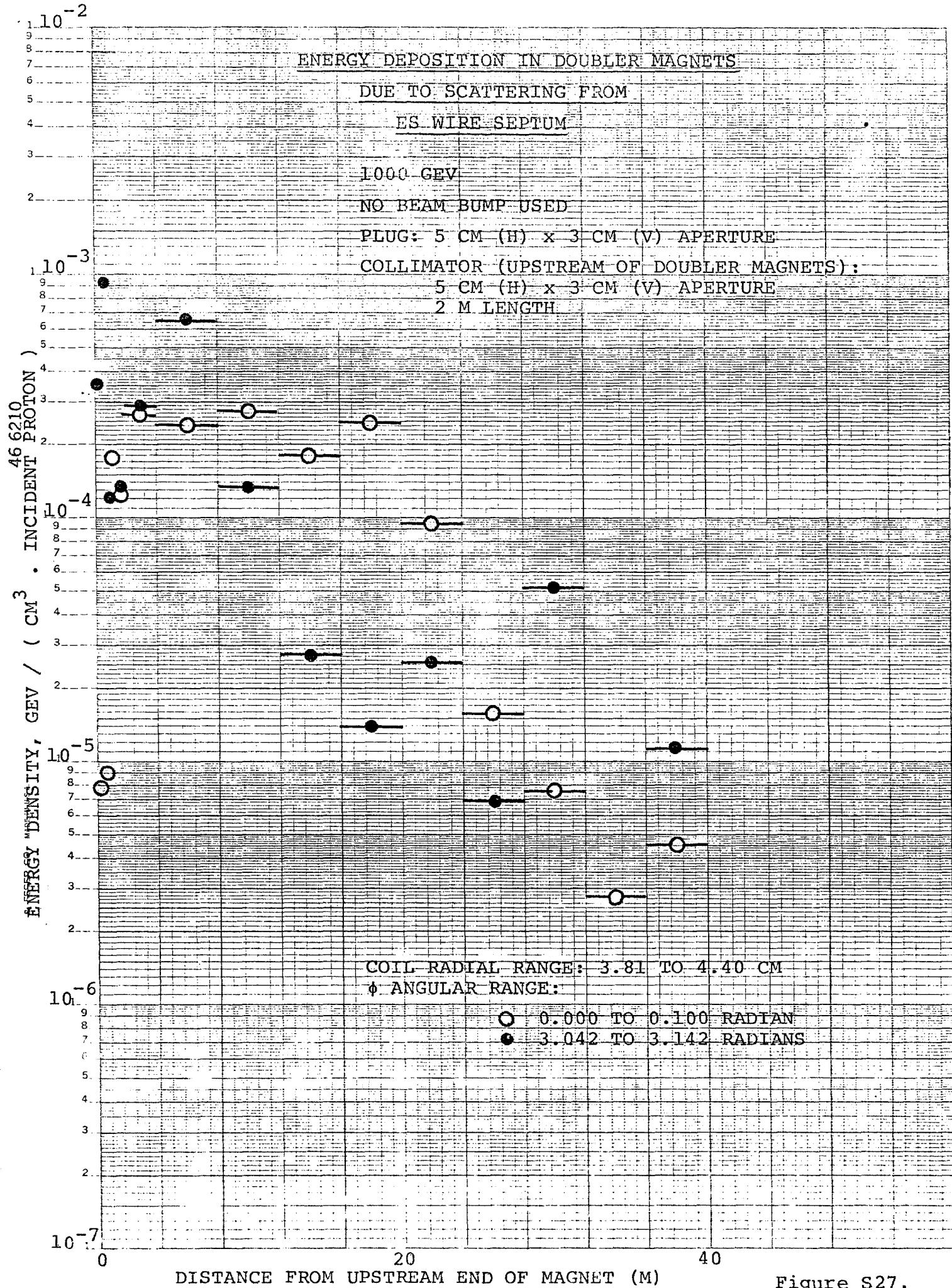


Figure S27.

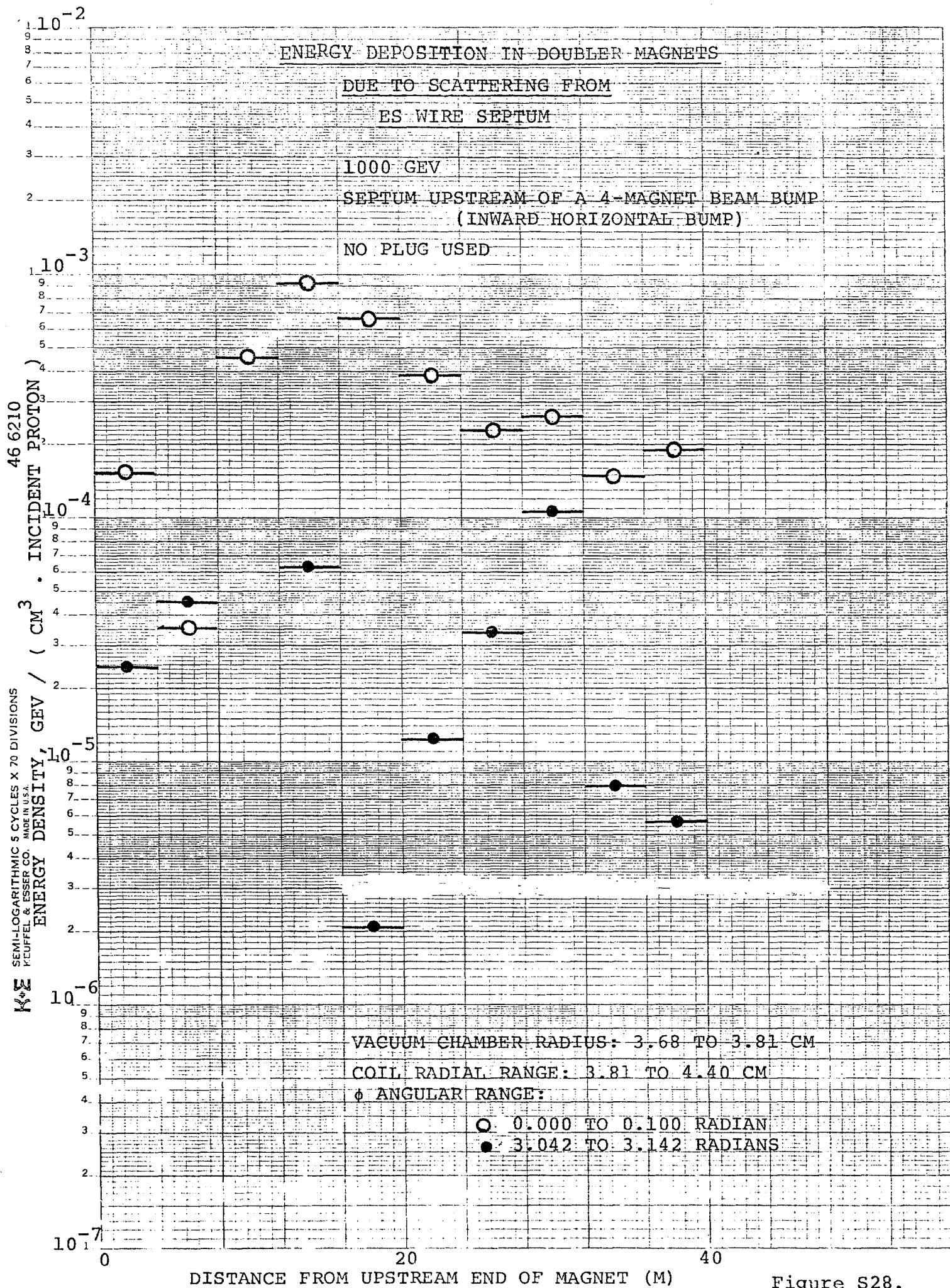


Figure S28.

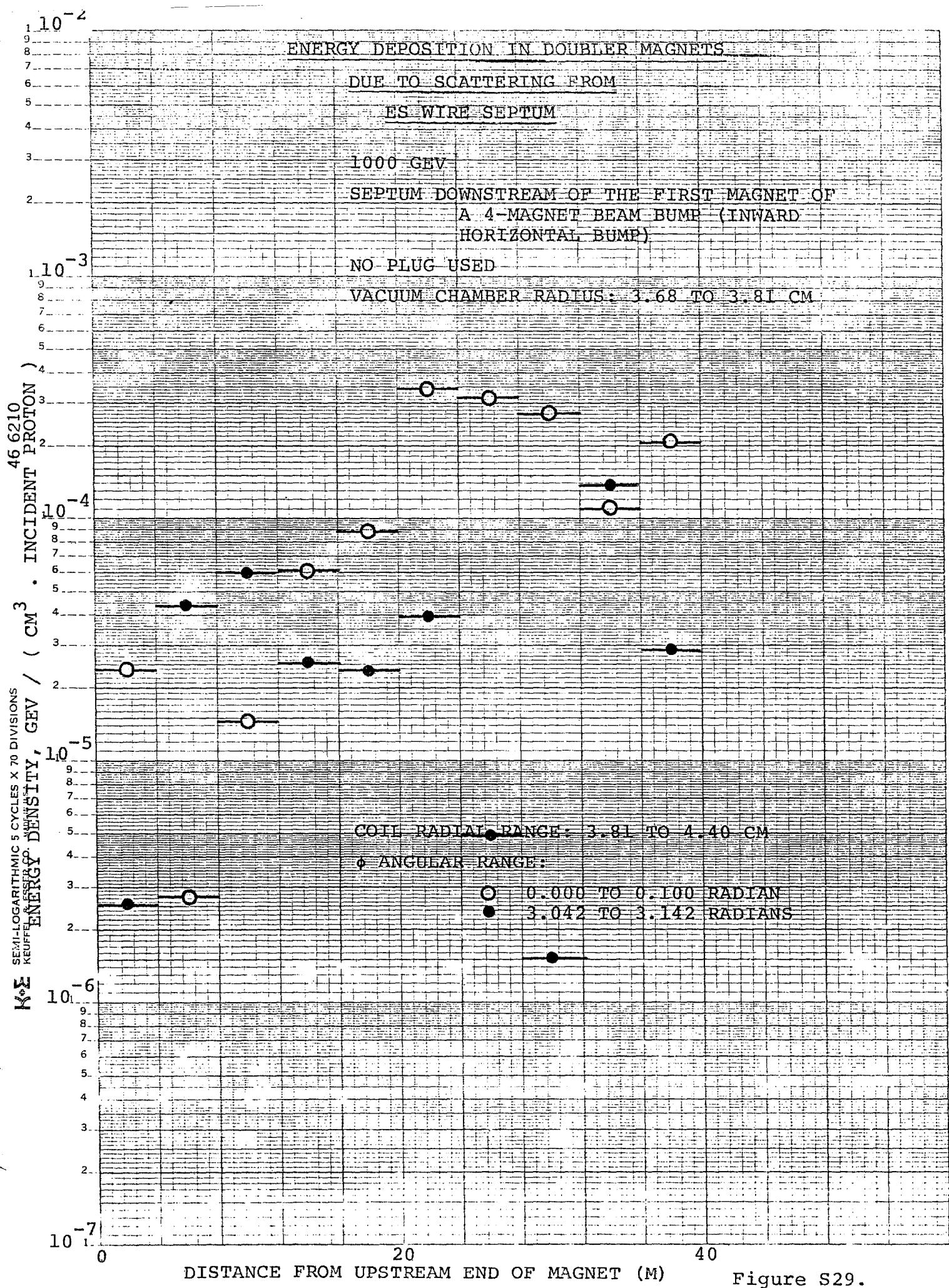


Figure S29.

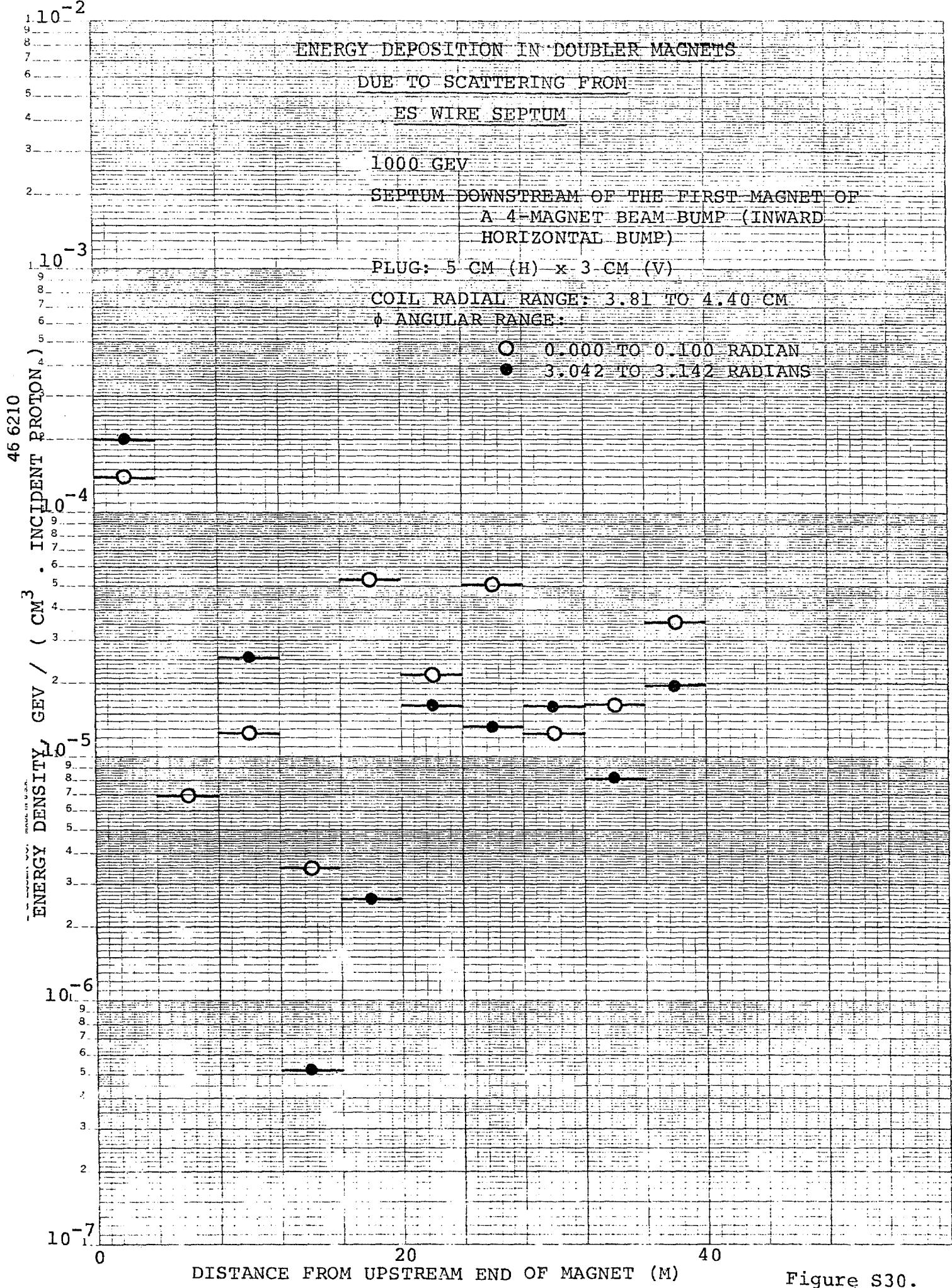


Figure S30.

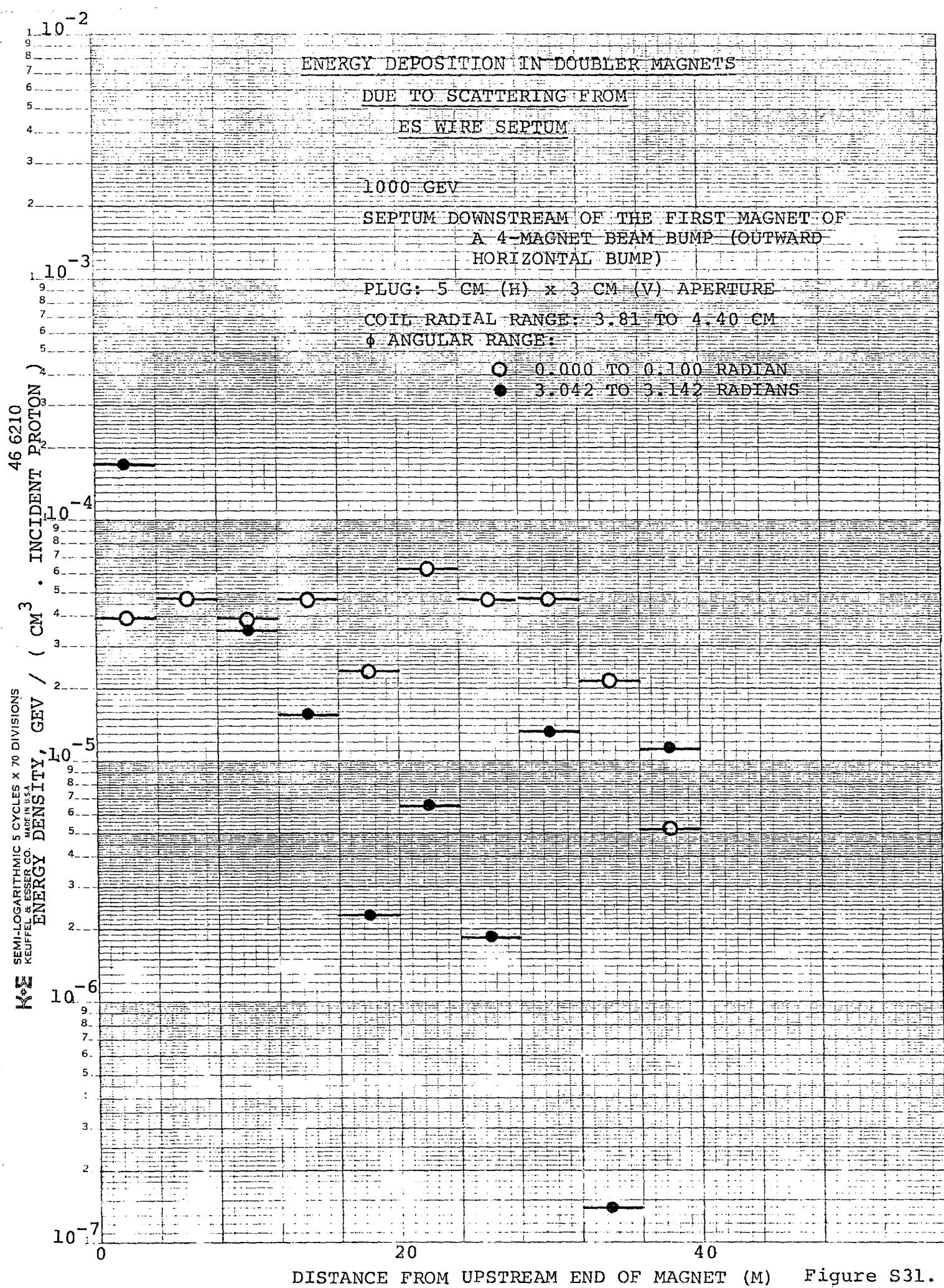


Figure S31.

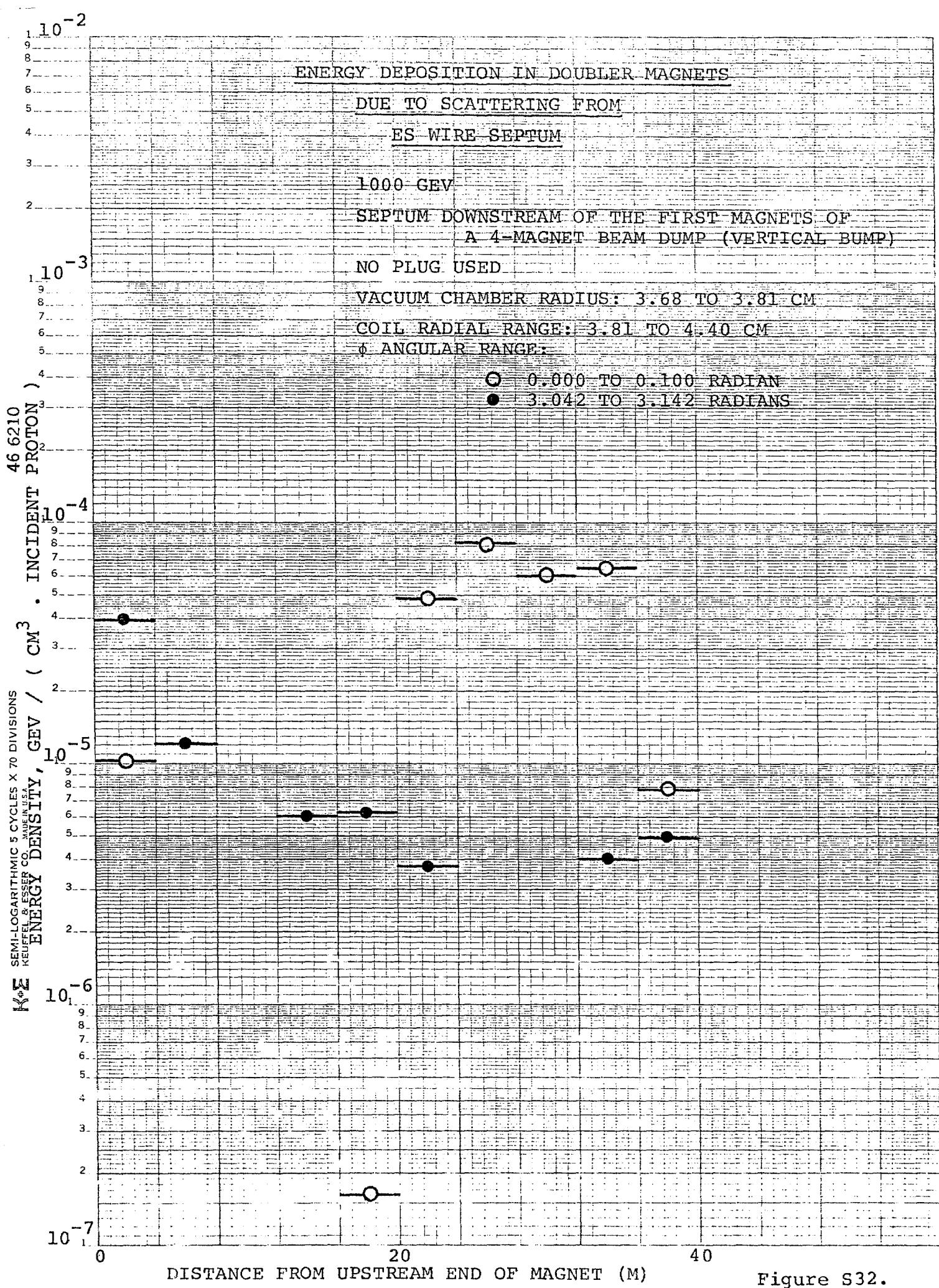


Figure S32.

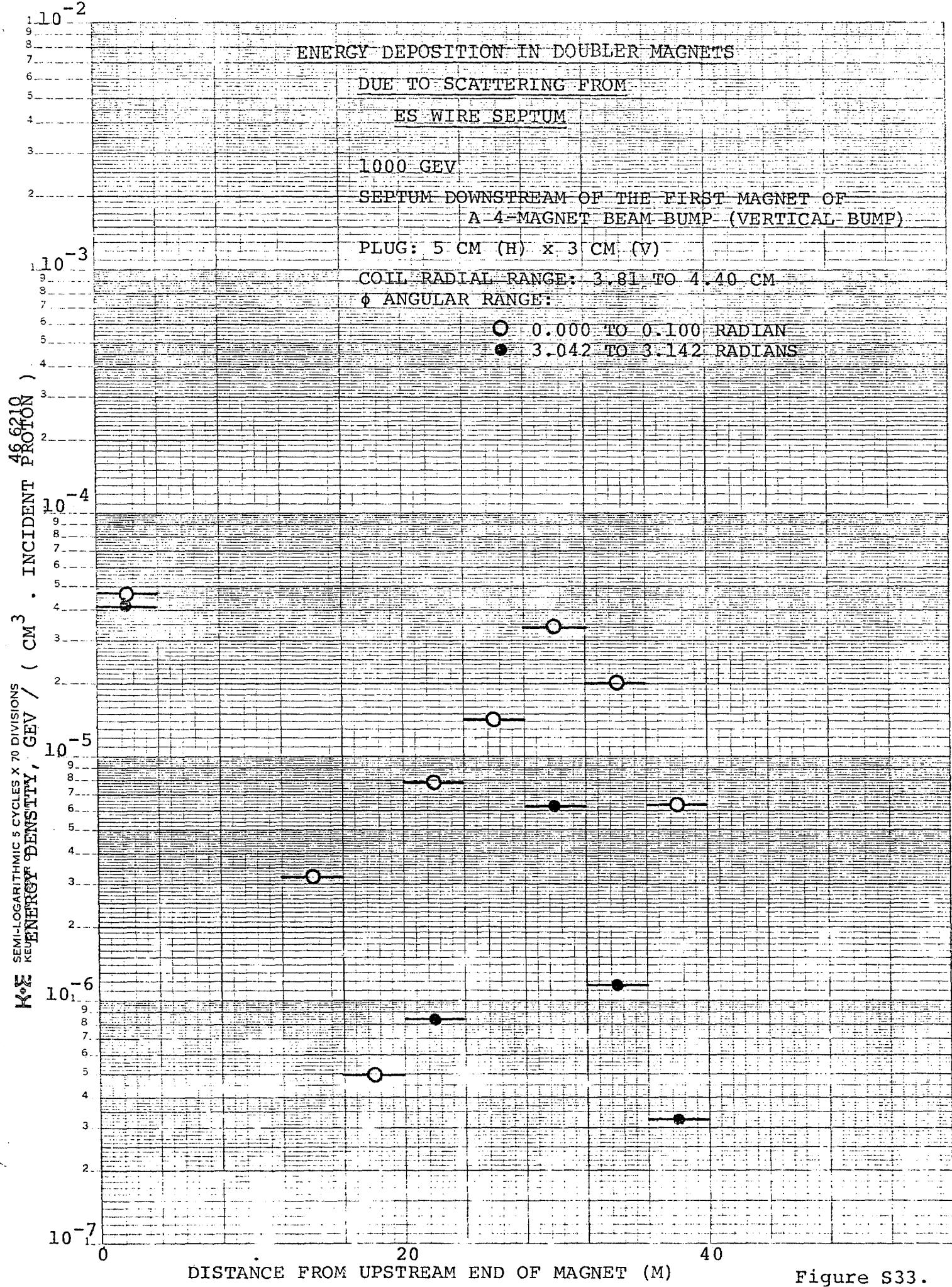


Figure S33.