

Preliminary
DraftTransverse Cooling in Precooler

The merits and implications of introducing transverse cooling in the precoolers are briefly discussed.

I. Reference Case :

Characteristic parameters:

i) $\sigma = 5 \cdot 10^{-7}$ p accepted in $E_r = (E_v E_H)^{1/2} \approx 2.5 \pi \cdot 10^6$ ev, $\Delta E_r = \pm 0.02$

ii) $\frac{\Delta p}{p}$ reduced by a factor 7 in ΔE_r , $2.5 \leq \Delta E_r \leq 3.5$

iii) beam spot on target: $r_p = 0.15$ mm,

p divergence $X'_p \approx 16 \cdot 10^{-3}$

II. Claim from larger ε

with constant r_p , $\varepsilon \propto X'_p$. Available information indicates that $X'_p \gtrsim 50$ mrad does not bring substantial further gain.

Approximate γ yield by

$$\gamma \approx 2.5 \cdot \ln \frac{L \varepsilon}{2 r_p^2} \quad (\text{for } \beta_{0,p} \gg \beta_{0,p})$$

$$\text{No. electrons} = 3.4 \approx \frac{\gamma(3 \varepsilon_r)}{\gamma(\varepsilon_r)} \approx 5.1$$

$$\text{for } 60 \text{ mm} \geq L \geq 150 \text{ mm}$$

We use simply a factor 3 for the possible increase in yield.

This implies two transverse cooling systems (the use of skew quadrupoles may allow working with only one) each reducing ϵ_x, ϵ_y by a factor of 3.

To realize this gain of 3 all cooling times must not exceed that of the reference case.

In general:

$$\frac{1+\alpha}{1+\alpha/3} \leq G \leq 3$$

where $\tau_{unstack, deel} = \alpha \tau_{cool}$ in the reference case.

Example: if any of the cooling times increase to 6s

(instead of 2) and assuming $\tau_{unstack, deel} = 2s$,

we will skip one HX cycle and G is reduced to 1.5.

III. Rough Performance Estimate for Transverse Systems

The usual equation for transverse cooling is :

$$\frac{d\varepsilon}{dt} = -\frac{W}{N} \left[2g - g^2 [\Gamma + \eta] \right] \varepsilon \quad (1)$$

$$\Gamma = \text{mixing term} \cong \text{Max}(1, \left(h_{\text{max}} \frac{A_{\text{max}}^{\text{tr}}}{A_{\text{min}}^{\text{tr}}} \right) \ln \frac{A_{\text{max}}^{\text{tr}}}{A_{\text{min}}^{\text{tr}}}) = \Gamma(t)$$

$$\eta = \frac{\text{noise power}}{\text{signal power}} = \eta(t) = \eta(t=0) \frac{\varepsilon(0)}{\varepsilon(t)}$$

For simultaneous cooling in all 3 plane planes. $\Gamma(t)$ is controlled by the progress of longitudinal cooling. A general solution for (1) can immediately be written down but is not very transparent and optimal gains and cooling times are not obtained without numerical integrations for $T = \Gamma(t)$, $\eta = \eta(t)$, $T \propto \eta$.

We distinguish 2 cases:

a) cooling at $\sim 4 \text{ GeV}$

b) cooling at $\sim 200 \text{ MeV}$, applicable for decelerating precoolers

No. disease (i) first. ($\perp$ cooling at 4 GeV)

A system with $W \approx 500 \text{ MHz}$ ($250 \text{ MHz} < f < 750 \text{ MHz}$)

appeared rather hopeless. A second case with
 $\perp \text{ GHz} < f < 4 \text{ GHz}$ is discussed in the following.

For the mixing parameter we obtain:

$\frac{\Delta p}{p} [10^{-2}]$	Γ
$2 \cdot 10^{-2}$	1
10^{-2}	1
$5 \cdot 10^{-3}$	2
$1.5 \cdot 10^{-3}$	4

This case is then (primarily) noise limited, rather than
 mixing limited:

$$\frac{d\varepsilon}{dt} \approx -\frac{W}{N} [2g - g^2\eta] \varepsilon \quad (2)$$

We invoke the capability to continuously adjust the
 gain factor to its optimum value:

$$g_0 = \frac{1}{\eta} \quad \text{leading to}$$

$$\frac{d\varepsilon}{dt} \approx -\frac{W}{N} \frac{\varepsilon}{\eta} = -\frac{W}{N} \frac{\varepsilon^2}{\eta(0) \varepsilon(0)} \quad (3)$$

with the solution

$$\varepsilon(t)/\varepsilon(0) = \frac{t_0}{t+t_0}, \quad t_0 = \frac{N}{W} \eta(0) \quad (4)$$

Note: with constant gain $\varepsilon(t) \rightarrow \varepsilon(\infty) \neq 0$ exponentially,

to reach a given $\varepsilon(t)$ then typically takes 1.3 to 1.7 times longer than in the adjusted gain case.

For a reduction by a factor 3 we require $\Delta t = 2 t_0$

$$\Delta t = \frac{2N}{W} \eta(0) \rightarrow \eta(0) = \frac{\Delta t W}{2N}$$

with $W = 3 \text{ CHZ}$, $2N = 2.3 N_r = 3 \cdot 10^8$ we obtain

$\Delta t [\text{s}]$	$\eta(0)$
2	20
4	40
6	60

These values justify neglecting Γ in the original equation.

We must establish the feasibility of obtaining these

noise/signal ratios, discuss required bricker power and

state lattice requirements for these components.

We analyze η for a PU's following Höhl (CERN PS/SL 78-25):

$$\eta = \frac{\xi}{n} \frac{\pm^2}{I_{pu}^2}$$

where ξ represents signal power loss in the combining network, $\xi > 1$.

$$\eta(0) \approx \frac{8\xi}{n} \frac{10^{NF/10} kT}{S^2 R c^2 N f_0}$$

with $S = 0.5$, $NF = 1.5 \text{ dB}$, $R = 50 \Omega$, $T = 300$, $N = 1.5 \cdot 10^6$,

$f_0 = 0.6 \text{ MHz}$, $\xi \approx 2$:

$$\eta(0) \approx \frac{1}{n} \cdot 3.25 \cdot 10^3$$

And therefore:

$\Delta t. [s]$	$\eta(0)$	n
2	20	~ 160
4	40	~ 80
6	60	~ 54

Remarks:

- i) $S = 0.5$ is an ICE value (see Höhl's paper) and may not be applicable to GHz range PUs
- ii) the signal combining business should be explored somewhat more carefully
- iii) an actual PU designer is not available now,

a wall current monitor (difference between L-R, U-D)
might be a reasonable choice. The

iv) the gap spacing needs investigation; we assume
here ~ 10 cm. This may be overly conservative,
(the AA precooling assembly has gaps spaced ~ 2 cm)

Case b) low energy (200 keV + cooling)

In this case the work of the LBL experiment at
the ECR should be directly applicable. At present
experimental data are not available but we expect
that a reduction of transverse emittances by a
factor of 3 should be possible in ≤ 0.5 s.

IV Power Requirements

In the noise dominated regime the power depends on B_p, v, ϵ_0, t_0 and also the specific type of kicker chosen. For a kicker with a gap not much wider ~~than~~ (higher) than the beam we estimate

$$P_t \sim 3 \frac{(B_p v \epsilon_0)^2}{L^2 f_0 K_k} \frac{1}{t_0} \quad \text{for single kicker}$$

$$\sim \frac{3}{n} \frac{(B_p v \epsilon_0)^2}{L^2 f_0 K_k} \frac{1}{t_0} \quad \text{for } n \text{ kickers}$$

For the low energy case a single TW-kicker

($\beta_w = \beta_r$) is envisaged and power levels of a few kilowatts (200 to 400) are expected for

$$\Delta t = 0.4 \text{ s}, \quad f_0 = 0.6 \text{ MHz}, \quad \epsilon_0 = 3 \cdot \epsilon_{0, \text{ref.}} \quad (A_0 = 12 \text{ cm})$$

For the high energy case, using a kicker array of ~ 200 simple plates $\sim 2.5 \text{ cm}$ long in beam direction $P_t \sim 1 \text{ to } 2 \text{ kW}$ is estimated.

V. The Three Requirements

a) Feasibility for all systems:

Sufficient path differences (beam - signal) for overcomes amplifier delays.

High energy: assume approximately circular orbit

$$S_B - S_S \approx C \cdot \Delta \varphi_{amp}$$

$$\theta - 2 \sin \frac{\theta}{2} \approx \frac{r_c}{\Delta \varphi_{amp}}$$

$$\text{typical } \Delta \varphi_{amp} \sim 5 \text{ ns}$$

with $R = 7.5 \text{ m}$ we need an azimuthal

separation θ between P_u and K of $\geq 0.55 \text{ sr}$

In the limit of azimuthal $W-K$ mixing

the K beam should be as close as possible to the

pick-up (but of course satisfy the above

constraints). In the longitudinal case the correction

may be applied on full turn delayed ($\leftarrow$ transverse)

which requires these constraints.

low energy (200 keV)

In this case a safe K-PU separation is:

$$\Delta s_{\min} = \frac{\beta}{1-\beta} \cdot \Delta t \cdot c \cong 20 \text{ m}$$

b) transverse cooling systems

high & low energy:

i) phase advance. $\Delta\psi_{pu,k} = \frac{2k+1}{2} \pi$, $k=0,1,2, \dots$

ii) dispersion $\eta = 0$ at horizontal pick-up

high energy

i) PU-straight: $L \lesssim 20 \text{ m}$, "high β ", $\int_{L_{pu}} \frac{ds}{\beta(s)} \lesssim 10^\circ$

try to "fill" PU of $\sim 50 \text{ mm}$ (roughly)

horizontal & vertical PU's may possibly be combined

requiring then only one such straight

ii) K-straights: 2 (vertical & horizontal)

$L \approx 10 \text{ m}$, probably best to keep smallest

, satisfying $\int_{L_K} \frac{ds}{\beta(s)} \lesssim 30^\circ$

low energy

i) PU: 2 (vertical and horizontal)

each $\sim 5m$ long, $\int_{P_u} \frac{ds}{\beta(s)} \lesssim 10^\circ$

both pick-ups may be placed in same straight

(provided suitable kicker locations with correct phase advance are provided)

ii) K-straights:

2 (vertical & horizontal)

each $\sim 2m$ long, again ideally small β ,

but $\int_{L_k} \frac{ds}{\beta(s)} \lesssim 30^\circ$

c) longitudinal system

The same system may be used at high and low energy [except notch filter]

i) dispersion: $\eta = 0$ at kicker

ii) Length $\left. \begin{array}{l} L_{pu} \lesssim 20m \\ L_k \gtrsim 20m \end{array} \right\} \begin{array}{l} \text{as proposed seems} \\ \text{reasonable.} \end{array}$

iii) β -function: allow reasonable PL_{K-1D}

if decelerating with increased acceptance:

try to approach $\beta_{max,pu} \approx \beta_{max,K} \rightarrow L_{pu}, L_K$

VI. Tentative Conclusions

in $\bar{p}$ accumulation

- i) a total gain of 1.5 to 3₁ may be feasible using stochastic, transverse cooling
- ii) the estimates of PL-sensitivity, questions of electronics and trigger design of a ~ 4 GHz high energy cooling system contain substantial uncertainties
- iii) the proposed PL-gap spacings (leading to long straight sections) may be conservative (compare AA-ring); this should (if necessary experimentally) be explored
- iv) transverse cooling at low energy (i.e. in a decelerating precoolers) appears more promising, since we have at least bunch lots of PL's, will have experimental data soon, estimate cooling times < 0.5 s and require much less straight section length.