

Fusion hindrance for asymmetric systems at extreme sub barrier energies

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Introduction

Recent measurements with medium-heavy nuclei accentuated phenomenon of fusion hindrance, observed as a steep change of slope in fusion excitation function and its logarithmic derivative ($L(E)$) with respect to the coupled channels (CC) calculation at deep sub-barrier energies [1]. At present there are two successful models to explain the deep sub-barrier fusion data - model suggested by Mišicu and Esbensen [2] based on sudden approximation using M3Y potential with repulsive core and a dynamical two-step model proposed by Ichikawa *et al.* [3] based on an adiabatic picture. Currently experimental studies at these low energies have been restricted mainly to the measurement of fusion cross sections of symmetric systems with the exception of $^{16}\text{O} + ^{204,208}\text{Pb}$ [4] and $^6\text{Li} + ^{198}\text{Pt}$ [5]. Unlike the sharp change in slope of $L(E)$ as observed in symmetric medium-heavy systems, a saturation in the slope of $L(E)$ was observed for asymmetric $^{16}\text{O} + ^{208}\text{Pb}$ system [4]. In case of very asymmetric system involving light weakly bound projectile $^6\text{Li} + ^{198}\text{Pt}$ [5] absence of fusion hindrance was reported. In the present work we investigate whether absence of fusion hindrance in $^6\text{Li} + ^{198}\text{Pt}$ system arises from the effect of weakly bound cluster structure or it is a property of the very asymmetric systems. For this purpose we selected one

weakly bound (^7Li) and a tightly bound (^{12}C) projectile on the same target. The data is analysed with standard coupled channels and the adiabatic model of fusion to understand the underlying mechanism at deep sub barrier energies.

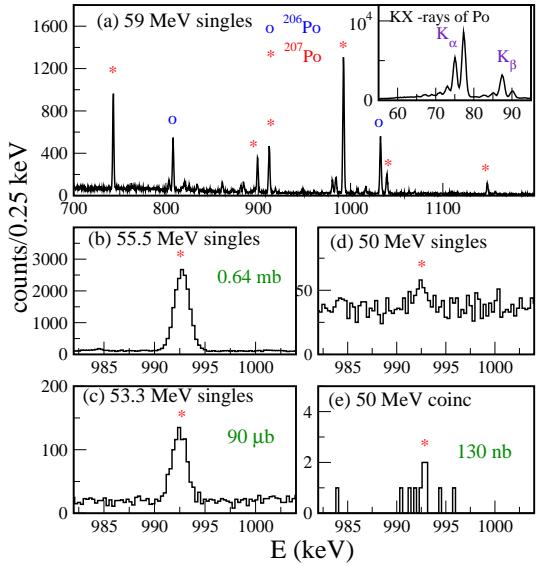


FIG. 1: Activation γ -ray spectra for the $^{12}\text{C} + ^{198}\text{Pt}$ system (a) inclusive spectrum at E_{lab} of 59 MeV. Inset shows X-ray region of the spectrum. The dominant γ -rays arising from the evaporation residues are labeled (b)-(d) inclusive spectra showing photo-peak at 992.3 keV, corresponding to the residue ^{207}Po (e) same as (d) but in coincidence with the $\text{K}_{\alpha}\text{X}$ rays shown in the inset of (a).

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Experimental Details

The experiment was performed at Pelletron Linac Facility-Mumbai, using beams of ^7Li (20 - 35 MeV) and ^{12}C (50 to 64 MeV) on a ^{198}Pt target with beam current in the range of 10 to 35 pA. The targets were self supporting rolled foils of enriched ^{198}Pt ($\sim 1.3 \text{ mg/cm}^2$ thick) followed by an Al catcher foil. Two efficiency calibrated HPGe detectors - one with Al window for detection of γ -rays and another with a Be window for detection of KX-rays were used for performing KX- γ -ray coincidence. The measurements were performed in a low background setup with a graded shielding. The evaporation residues were uniquely identified by means of their characteristic γ -ray energies and half-lives. The residues populated are $^{205-207}\text{Po}$ in case of $^{12}\text{C}+^{198}\text{Pt}$ and $^{200-202}\text{Tl}$ in case of $^7\text{Li}+^{198}\text{Pt}$. Typical inclusive γ -ray spectra resulting from the residues of $^{12}\text{C}+^{198}\text{Pt}$ are plotted in Fig. 1(a) - (d) at different beam energies. The γ -ray yields at lowest energies were extracted by gating on their KX-ray transitions. Due to the increased sensitivity of the KX- γ -ray coincidence method, cross-section down to a few nano-barns could be measured (Fig. 1(e)).

Analysis and Summary

The data were analysed using the standard coupled-channels (CC) calculations and the adiabatic model that simulates a smooth transition between the two-body and the adiabatic one-body states by damping gradually the off-diagonal part of the coupling potential [6]. The standard CC calculations for both the systems were performed including the quadrupole excitation in ^{198}Pt in the vibrational mode. Projectiles ^7Li and ^{12}C were coupled in the rotational mode. In case of $^7\text{Li}+^{198}\text{Pt}$, fusion hindrance was not observed as the CC calculations nicely reproduce the data for energies around and well below the barrier in the measured energy range. A change in slope in fusion excitation function and $L(E)$ as compared to CC calculations was clearly observed for $^{12}\text{C}+^{198}\text{Pt}$, at lowest energies indicating onset of fusion hindrance. In order to explain the fusion data at energies

deep below the barrier in case of $^{12}\text{C}+^{198}\text{Pt}$ system, calculations were performed using the adiabatic model of Ref. [6]. On inclusion of damping in the adiabatic framework an excellent agreement with the fusion and $L(E)$ data was observed as shown in Fig. 2. These results are relevant to fusion hindrance at deep sub-barrier energies with respect to observations from $^{6,7}\text{Li}+^{198}\text{Pt}$, $^{16}\text{O}+^{208}\text{Pb}$ and systems with different mass asymmetry.

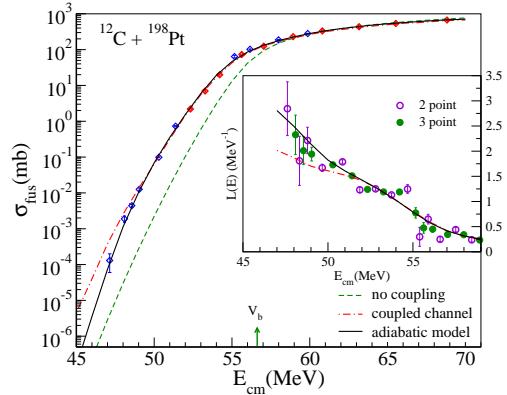


FIG. 2: Fusion excitation function and $L(E)$ for $^{12}\text{C}+^{198}\text{Pt}$. CC calculations with and without inclusion of coupling along with adiabatic model calculations are shown.

Acknowledgments

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