

## Shape evolution of nuclei in $A \sim 55$ region

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### Introduction

The nuclear structure of  $A \sim 55$  region are governed by  $1f_{7/2}$  and intruder  $g_{9/2}$  orbitals across  $N = Z = 28$  shell closure. The nature of this shell closure, which is the first shell closure originated due to the spin-orbit coupling term, is in debate from various theoretical and experimental findings. Further, the presence of shape driving orbitals like  $1f_{7/2}$  and  $1g_{9/2}$  can induce deformation in the system. There are limited informations regarding the observation of different nuclear shapes in this mass region with nuclei around doubly magic  $^{56}\text{Ni}$  ( $Z = 28, N = 28$ ).

In this work, we have calculated the deformations for even-even nuclei from  $Z = 22$  (Ti) to  $Z = 28$  (Ni) for different  $N$  to get an idea regarding the evolution of nuclear shapes with change in proton and neutron Fermi surfaces in this mass region. The transition rate and quadrupole moment,  $B(E2)$  and  $Q_t$ , have been deduced from the deformation parameters, obtained from the calculation, and compared with the experimental ones. The experimental  $B(E2)$  and  $Q_t$  values were deduced from the measured lifetimes ( $\tau$ ) reported for the  $1^{st}$  ( $2^+$ ) and using the equations given in [1]. These calculations have been performed for lower rotational frequencies corresponding to spin value of  $2\hbar$ .

### Method of the calculation

The shape of a nucleus has been obtained from the Total Routhian Surface (TRS) calculations. In this method, the potential energies in the body-fixed frame are calculated in the Nilsson-Strutinsky formalism [2]. A deformed Woods-Saxon potential with universal param-

eters is used to calculate the single particle energies. The nuclear deformation is defined by the  $\beta_2$ ,  $\beta_4$  and  $\gamma$  parameters. In the Lund convention, used here,  $\gamma = 0^\circ$  ( $-60^\circ$ ) is for axially deformed prolate (oblate) shape and any other value of  $\gamma$  indicate triaxial shape. The detailed technical procedure has been given in [3, 4]. The total Routhians are obtained at different values of  $\beta_2, \beta_4, \gamma$  and are plotted in the  $\beta_2, \gamma$  surface after minimizing on  $\beta_4$ . The calculations are done for different quasi-particle configurations and at several rotational frequencies ( $\hbar\omega$ ). The shape of a nucleus for a particular  $\hbar\omega$  at a particular configuration is obtained from the values of  $\beta_2$  and  $\gamma$  corresponding to the minimum of the TRS at that  $\hbar\omega$  for that configuration.

### Results and Discussions

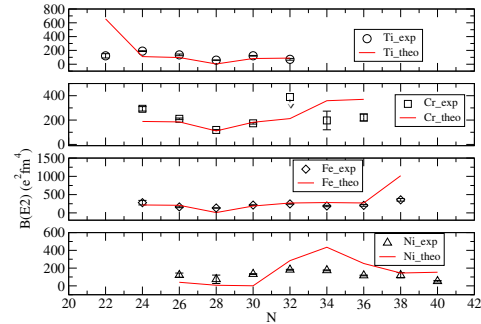


FIG. 1: The experimental values (symbols) of  $B(E2)$  for the Ti, Cr, Fe and Ni isotopes are compared with those obtained from the  $\beta_2$  and  $\gamma$  deformation parameters obtained from the TRS calculations (solid line) in this work.

The theoretical results are summarized in Table I. It contains the deformation parameters  $\beta_2$  and  $\gamma$ , obtained from the present calculation, for various nuclei with even  $Z$  from Ti ( $Z = 22$ ) isotopes to Ni ( $Z = 28$ )

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TABLE I: Deformation parameters ( $\beta_2$  and  $\gamma$ ) of the nuclei near the ground state obtained in this work.

Z	N	Nucleus	$\beta_2$	$\gamma$	$B(E2)(e^2fm^4)$
22	22	$^{44}Ti$	0.4	-21.8	656.83
22	24	$^{46}Ti$	0.19	-2.7	110.55
22	26	$^{48}Ti$	0.17	-5.6	97.06
22	28	$^{50}Ti$	0.03	-30	3.49
22	30	$^{52}Ti$	0.15	-5.4	82.67
22	32	$^{54}Ti$	0.14	-21	88.76
24	24	$^{48}Cr$	0.22	-2.2	188.75
24	26	$^{50}Cr$	0.21	-3.6	184.96
24	28	$^{52}Cr$	0.15	-14.8	110.83
24	30	$^{54}Cr$	0.2	-2.8	182.07
24	32	$^{56}Cr$	0.2	-9.6	212.24
24	34	$^{58}Cr$	0.26	-3.9	359.08
24	36	$^{60}Cr$	0.26	-3.1	370.49
26	24	$^{50}Fe$	0.21	-3.4	216.32
26	26	$^{52}Fe$	0.2	-4	207.50
26	28	$^{54}Fe$	0.04	-18.6	9.30
26	30	$^{56}Fe$	0.18	-5.4	187.32
26	32	$^{58}Fe$	0.2	-12.4	269.95
26	34	$^{60}Fe$	0.21	-4.8	282.47
26	36	$^{62}Fe$	0.2	-6.1	271.44
26	38	$^{64}Fe$	0.38	-0.7	1016.86
28	26	$^{54}Ni$	0.08	-8.7	40.07
28	28	$^{56}Ni$	0.03	-30	6.58
28	30	$^{58}Ni$	0.01	-23.8	0.75
28	32	$^{60}Ni$	0.18	-36.8	284.07
28	34	$^{62}Ni$	0.23	-52.3	435.01
28	36	$^{64}Ni$	0.17	-47.9	251.91
28	38	$^{66}Ni$	0.13	-50.4	144.88
28	40	$^{68}Ni$	0.18	-77.8	153.61

isotopes. These shapes are calculated near the ground state at lower  $\hbar\omega$ .

The calculated values of the deformation parameters are used to calculate the other parameters like the quadrupole moment  $Q_t$  and the transition strengths  $B(E2)$ . The experimental values of these parameters are extracted from the measured lifetimes ( $\tau$ ) reported in the ENSDF [5]. The expressions used for determining these quantities in both the cases are taken from reference [1]. These

measured quantities are for the  $1^{st} 2^+$  states in all these cases. The experimental and theoretical values of  $Q_t$  and  $B(E2)$  are compared in Fig. 1. It can be seen that the experimental values are well reproduced by the calculations.

## Summary and Conclusion

The calculated results predict different shapes for different N/Z. At  $N = 28$ , which is the magic number, in all the isotopes the deformation seems to decrease and shape tends towards triaxiality. For Ti and Cr isotopes, the shapes are mostly prolate except  $N = 28$ . In Fe isotope,  $N=28$  and 32 are triaxial while other N are prolate in shape. Thus the evolution of shapes from prolate to triaxial and again to prolate are predicted for these isotopes from these calculations. The Ni isotope ( $Z = 28$ ) is predicted to have evolution of shape from prolate to oblate via triaxial shape with increase in N. The experimental  $B(E2)$  values are well reproduced by the theoretical calculations which validate the present calculations. Experimental observations of this shape evolution in these nuclei will be interesting.

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## References

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