

## COMPARISON OF ELECTROWEAK PARAMETERS

## INTRODUCTION

$\sin^2(\theta_{\text{bar}})$ , the square of the Weinberg angle at the Z-pole, has recently been measured both at LEP and by CDF. This note is intended to address, given a measurement of  $\sin^2(\theta_{\text{bar}})$ , how one can compare to other measurements of Electroweak parameters,  $M_W$  and  $1-(M_W/M_Z)^2$  in particular.

There is currently considerable disagreement within CDF about how to make such a comparison (not about the  $\sin^2(\theta_{\text{bar}})$  measurement itself). My hope is that this note will help clarify matters. It is a longer version of a mail message on this subject sent to interested parties on January 25, 1991.

## MINIMAL STANDARD MODEL, BUT IGNORING RADIATIVE CORRECTIONS

When radiative corrections are ignored, the minimal Standard Model has just 3 independent parameters. If these are chosen to be  $\alpha$ ,  $G_F$  and  $\sin^2(\theta)$ , one can express the masses  $M_W$  and  $M_Z$  as

$$M_W^2 = (\pi\alpha)/(\sqrt{2}G_F\sin^2(\theta)) = (37.281 \text{ GeV})^2/\sin^2(\theta)$$

$$M_Z^2 = M_W^2/\cos^2(\theta) = (37.281 \text{ GeV})^2/(\sin^2(\theta)\cos^2(\theta))$$

It is easy to use a measurement of  $\sin^2(\theta)$  to predict  $M_W$  using these relations. If, for instance,  $\sin^2(\theta) = 0.228 \pm 0.016$ , one can deduce that  $M_W = 78.1 \pm 2.7 \text{ GeV}/c^2$ . Note that this is indeed what was done some years ago, before the W and the Z were found, with an error bar on  $\sin^2(\theta)$  not much different from the one used in this example (for the error bar, I am just using  $dM_W = -(M_W/2)(d\sin^2(\theta)/\sin^2(\theta))$  here).

As we will see below, including radiative corrections does not make much difference.

## MINIMUM STANDARD MODEL, INCLUDING RADIATIVE CORRECTIONS

A particularly nice way to look at electroweak radiative corrections was presented by Chris Hill in his January 1991 Academic Training Lectures, see Fermilab-Pub-91/Draft-T, version 2.0, upon which these comments are based. I am grateful to Chris for several useful discussions on this subject. I wish to stress, however, that I (not he) am responsible for any and all misunderstandings and errors that may exist in this CDF note.

Ignoring the weak Higgs mass dependence, there are only 4 parameters needed to calculate all others in the minimal Standard Model, given the fact that the running of the coupling constants is specified by the model, and that the effective masses of all the Fermions, except for the top quark, are known. Let  $\alpha(M)$  denote the value of the running  $\alpha$  at mass  $M$ .

It is natural to pick  $\alpha$ ,  $GF$  and  $M_Z$  as three of the four parameters since all three are now known very accurately. The choice of the fourth is a matter of convenience. In this case, where we have measured  $\sin^2(\theta_{\text{bar}})$  and wish to predict  $M_W$ , it is clearly expedient to use  $\sin^2(\theta_{\text{bar}})$  as the fourth. There are then known functions that give

$$M_W = M_W(\alpha(M_W), GF, M_Z, \sin^2(\theta_{\text{bar}}))$$

$$M_{TOP} = M_{TOP}(\alpha(M_W), GF, M_Z, \sin^2(\theta_{\text{bar}}))$$

$$1 - (M_W/M_Z)^2 = \text{function of } (\alpha(M_W), GF, M_Z, \sin^2(\theta_{\text{bar}})) \text{ , etc.}$$

So we don't need  $M_{TOP}$  to get  $M_W$  from  $\sin^2(\theta_{\text{bar}})$ !!

Using relation 3.21 from Chris Hill, it is easy to work out the details:

$$M_W^2 = (PI \cdot \alpha(M_W)) / (SQRT(2) \cdot GF \cdot \sin^2(\theta_{\text{bar}}))$$

which is very similar to the relation without radiative corrections. This leads to the following table:

Source	$\sin^2(\theta_{\text{bar}})$	$M_W$	$1 - (M_W/M_Z)^2$
LEP	0.2302+- 0.0021	80.13+-0.37	0.2276+-0.0070
CDF	0.228 +- 0.016	80.52+-2.83	0.220 +-0.055
CDF	-----	79.91+-0.39	0.2319+-0.0075

I have used  $\alpha(M_W) = 1/128.85$  here, ignoring any error in  $\alpha$ . This is certainly OK as far as the new CDF measurement is concerned, since the error on  $\alpha$  is probably less than 0.4%. The LEP numbers, also taken from Chris Hill's note, are shown for comparison to our own. The last column is obtained using  $M_Z = 91.177 \text{ GeV}/c^2$

So if one wanted to use our new  $\sin^2(\theta_{\text{bar}}) = 0.228+-0.016$  measurement to derive  $M_W$  for comparison with our  $M_W$  measurement, the correct statement would be:

The measurement of  $\sin^2(\theta_{\text{bar}}) = 0.228+-0.016$  implies a  $W$  mass of  $80.5+-2.8 \text{ GeV}/c^2$ , to be compared with our previous measurement  $M_W = 79.91+-0.39 \text{ GeV}/c^2$ . The measurements are clearly consistent.

I have two comments (opinions) about this:

- 1) Given the size of the error bar on  $M_W$ , this comparison seems rather uninteresting: It belongs to the time before accurate  $W$  and  $Z$  masses were known.
- 2) The comparison, if done, should be in terms of  $M_W$ , not in terms of  $1 - (M_W/M_Z)^2$ , since this quantity, especially when labeled  $\sin^2(\theta_{\text{Sirlin}})$  or similar, is known to generate confusion with other  $\sin^2(\theta)$ .

There is another important reason why it is better to make the comparison in terms of  $M_W$ . To understand this point, look at Fig. 3.3 from the CDF upgrade proposal, which is included in this note. You will see that  $\sin^2(\theta_{\text{bar}})$  is weakly dependent on both  $M_{TOP}$  and  $M(\text{Higgs})$ : Once you know  $M_Z$  accurately (from LEP), and you calculate within the Minimal Standard Model,  $\sin^2(\theta_{\text{bar}})$  can be predicted within a narrow range,  $0.228 < \sin^2(\theta_{\text{bar}}) < 0.236$  according to the figure,



if you restrict  $M_{TOP}$  to be below  $230 \text{ GeV}/c^{**2}$ . This range is substantially smaller than the range allowed by this measurement:

$\sin^2(\theta_{\text{bar}}) = 0.228 \pm 0.016$ . When comparing this measurement with MW measurements, it is therefore important to realize that:

a) To display the uncertainty from this measurement alone, one has to avoid using the implications of the LEP  $M_Z$  measurement, directly or indirectly.

b) If  $M_Z(\text{LEP})$  is used, the uncertainty will be substantially reduced, but it no longer reflects the result of this measurement alone.

As shown above, one can cleanly compare to the CDF MW measurement. If one floats  $M_Z$ , however, how exactly does one get from the CDF MW to  $1 - (MW/M_Z)^{**2}$  ?

## SO WHAT IS WRONG WITH THE COMPARISON IN THE DRAFT PRL?

---

The draft contains a figure showing "the value of  $\sin^2(\theta_{\text{Sirlin}})$  ( $= 1 - (MW/M_Z)^{**2}$ ) derived from the asymmetry as a function of the top quark mass" plus corresponding text. Our MW measurement, in the guise of  $1 - (MW/M_Z)^{**2}$  with MW from CDF and  $M_Z$  from LEP, is shown on the right of the figure to allow the comparison.

This way of comparing is WRONG because it does not obtain the correct result,  $M_W = 80.5 \pm 2.8 \text{ GeV}/c^{**2}$ , which was obtained above (in terms of  $1 - (MW/M_Z)^{**2}$ , the correct result is  $0.220 \pm 0.055$ , where the central value uses  $M_Z = 91.177$ , while the size of the error bar is rather insensitive to the  $M_Z$  value).

It is also highly MISLEADING: The construct "the value of MW derived from the asymmetry as a function of the top quark mass" says that you need  $M_{TOP}$  in addition to  $\sin^2(\theta_{\text{bar}})$  to get MW, whereas we know that  $\sin^2(\theta_{\text{bar}})$  alone (together with  $\alpha(M_W)$ , GF and  $M_Z$ ) is enough! There are only 4 independent parameters! (ignoring the weak Higgs mass dependence). Identical statements hold when MW is replaced by  $1 - (MW/M_Z)^{**2}$ .

Remark to the alert reader of the draft:

To obtain the functional  $M_{TOP}$  dependence in the figure, one has to replace either  $M_Z$  or GF by  $M_{TOP}$  as one of the four independent input parameters. This means that a measured, fixed constant ( $M_Z$  or GF as the case might be) is treated as a dependent, multivalued variable (i.e. the important constraint of fixed  $M_Z$  is ignored). Using the approximation  $M_Z = 38.45 \text{ GeV}/\sqrt{\sin^2(\theta_{\text{bar}}) \cos^2(\theta_{\text{bar}})}$ , for instance, I calculate approximate values of  $M_Z = 89.5 \text{ GeV}$ ,  $91.6 \text{ GeV}$  and  $94.1 \text{ GeV}$  for the three curves in the figure representing the CDF result (i.e. for  $\sin^2(\theta_{\text{bar}})$  of 0.244, 0.228 and 0.212, respectively).

Finally, I think, as stated above, that it would be much easier to express and understand if the comparison were done with MW rather than with  $1 - (MW/M_Z)^{**2}$ .

## CONCLUSION

---

I conclude that this figure plus the associated text must be removed from the paper. It is fortunate that this can be done easily: This comparison to MW has nothing to do with the  $\sin^2(\theta_{\text{bar}})$  measurement itself!

The corresponding plot must of course also be removed from the file of "blessed" CDF plots.

#### FINAL COMMENT

---

Let me conclude by advocating what I think is the nicest way for CDF to show comparisons of EWK variables: That is to use the plot of MW versus M<sub>TOP</sub>, both of which are quantities that we aim to measure. This plot can be used to show:

- 1) The calculated dependence  $MW = MW(\alpha(MW), GF, MZ, M_{TOP})$
- 2) Our (and other) MW measurements as horizontal bands, see for instance Fig. 3-1 in the CDF upgrade proposal.
- 3) The LEP measurement of  $\sin^2(\theta_{\text{bar}})$ , converted to a predicted MW as shown above, as a horizontal band
- 4) Any and all other relevant measurements, such as  $\nu$ -e scattering, again converted to a horizontal band via the relation

$$MW = MW(\alpha(MW), GF, MZ, \text{new measurement})$$

Note: To deal with precision experiments, one will have to include the Higgs mass dependence.

Figure 3-3 from the Upgrade Proposal

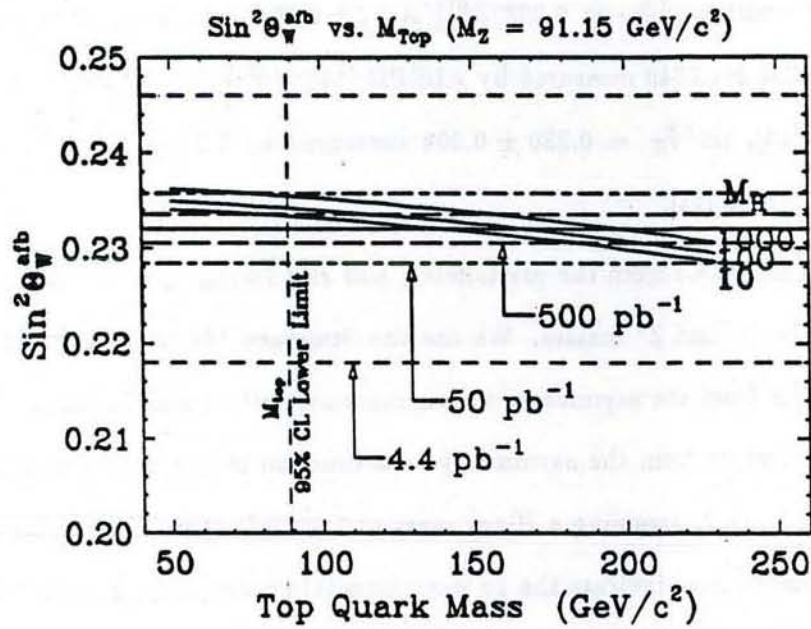


Figure 3-3: Predicted  $M_{top}$  dependence of  $\sin^2 \theta_W^{afb}$  from the combined results of forward-backward asymmetry measurements in  $Z \rightarrow e^+e^-$  and  $Z \rightarrow \mu^+\mu^-$ . Uncertainties are given for several values of integrated luminosity. Example given is for current CDF measurement.



Our final result,  $\sin^2 \bar{\theta}_W = 0.228^{+0.017}_{-0.015}(\text{stat}) \pm 0.002(\text{sys})$ , is in good agreement with  $\sin^2 \bar{\theta}_W = 0.2291 \pm 0.0040$  measured by ALEPH [14],  $\sin^2 \bar{\theta}_W = 0.2309 \pm 0.0048$  measured by DELPHI [15],  $\sin^2 \bar{\theta}_W = 0.230 \pm 0.004$  measured by L3 [16],  $\sin^2 \bar{\theta}_W = 0.233^{+0.007}_{-0.008}$  measured by OPAL [17].

$\sin^2 \bar{\theta}_W$  is measured from the asymmetry, and  $\sin^2 \theta_W|_{\text{Sirlin}} = 1 - M_W^2/M_Z^2$  [18] from the ratio of the  $W$  and  $Z^0$  masses. We use the Standard Model to convert our measurement of  $\sin^2 \bar{\theta}_W$  from the asymmetry into a measurement of  $\sin^2 \theta_W|_{\text{Sirlin}}$  [19]. A plot of  $\sin^2 \theta_W|_{\text{Sirlin}}$  derived from the asymmetry as a function of the top mass is shown by the solid line in Figure 2, assuming a Higgs mass of 250 GeV (the Higgs mass dependence is small); the dashed lines indicate the  $1\sigma$  experimental uncertainty. Figure 2 also shows the  $1\sigma$  confidence region derived from recent  $Z^0$  mass measurements [20] as well as a direct measurement of  $1 - M_W^2/M_Z^2 = 0.232 \pm 0.008$  [21] determined from the CDF  $W$  and LEP  $Z^0$  masses.  $A_{FB}$  is not strongly dependent on the top mass; the top mass dependence in Figure 2 comes from higher order corrections in the Standard Model incurred in the conversion from  $\sin^2 \bar{\theta}_W$  to  $\sin^2 \theta_W|_{\text{Sirlin}}$ .

In summary, we have measured  $A_{FB} = (5.3 \pm 5.9(\text{stat}) \pm 0.4(\text{sys}))\%$  after background and QCD corrections, and  $\sin^2 \bar{\theta}_W = 0.228^{+0.017}_{-0.015}(\text{stat}) \pm 0.002(\text{sys})$  after background and radiative corrections. The systematic uncertainties are summarized in Table 1. Our measurement of  $\sin^2 \bar{\theta}_W$  is consistent, both without and with order  $\alpha^3$  radiative corrections,

---

with previous measurements of  $\sin^2 \bar{\theta}_W$  at the  $Z^0$  mass. Our measurement of  $\sin^2 \theta_W|_{\text{Sirlin}}$  from the asymmetry is consistent with measurements of  $\sin^2 \theta_W|_{\text{Sirlin}}$  from the  $W$  and  $Z^0$  mass ratio over a broad range of top quark masses.

We thank the Fermilab Accelerator Division and the technical staffs of the participating institutions. We thank Y. Srivastava, C. Hill, and W. Bardeen for assistance. This work was supported by the Department of Energy, the National Science Foundation, the Istituto Nazionale di Fisica Nucleare, the Ministry of Science, Culture, and Education of Japan, and the A. P. Sloan Foundation.

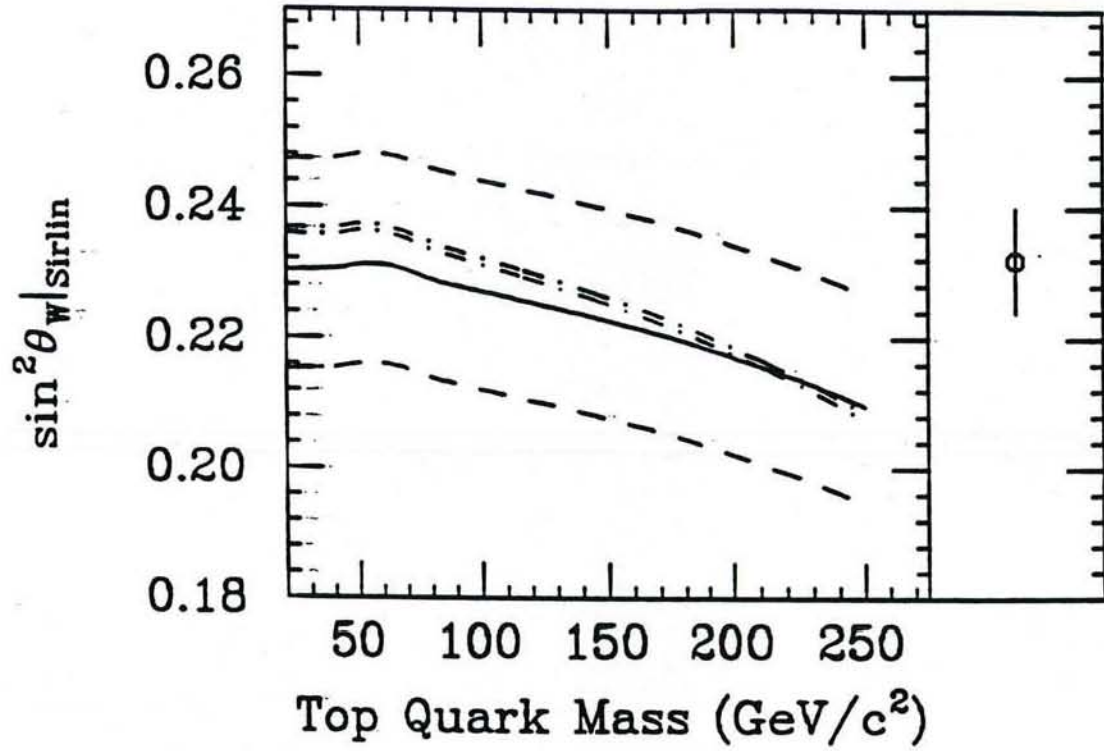


Figure 2: The solid line shows the central value of  $\sin^2 \theta_W |_{sirlin}$  derived from the asymmetry as a function of the top quark mass; the dashed lines indicate the  $1\sigma$  experimental uncertainty. The dot-dashed lines show the  $1\sigma$  uncertainty on  $\sin^2 \theta_W |_{sirlin}$  determined from  $Z^0$  mass measurements. At right is the CDF value from  $1 - M_W^2/M_Z^2$  [21].