

Molecular Design for Optically Induced Magnetization: Targeting Excited State Orbital Degeneracy in Tungsten(V) Complexes

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ABSTRACT: The rise of quantum information science has spurred chemists to prepare new molecules that serve as useful building blocks in quantum technologies of the future. Implementation of molecular spin-based qubits requires new methods to induce high spin polarization of samples. Herein, we report design criteria to develop axially symmetric spin-1/2 molecules amenable to optically induced magnetization (OIM), a technique using circularly polarized (CP) excitation to deliver spin polarization. We apply these criteria to develop a series of tungsten(V) chalcogenide complexes that are demonstrated to have large spin-sensitive responses to CP light using magnetic circular dichroism (MCD) that could allow up to ~20% spin polarization through OIM. Pulsed electron paramagnetic resonance (EPR) spectra reveal these systems have improved relaxation times over molecules like K_2IrCl_6 , a species recently investigated by OIM, and field-swept electron spin-echo (FS-ESE) experiments show they have a remarkable lack of anisotropy in their phase-memory T_m times. The design criteria are general and point toward future ways to improve OIM-initializable qubits.

The development of molecular spin-based qubits¹ requires that spins fulfill several criteria outlined by DiVincenzo for quantum computing² or Degen for quantum sensing.³ Of these criteria, use of molecular spin most prominently demands advancement in methods of “initialization,” the process of selective spin polarization into a pure $|0\rangle$ or $|1\rangle$ state. Spin polarization is usually done thermally: application of an external magnetic field splits the M_S levels, and a Boltzmann population of these levels is then achieved on the timescale of the spin–lattice (T_1) relaxation time.⁴ Conceptually, greater spin polarization can be easily reached by applying a larger magnetic field or cooling a sample further.⁵ Practically, this is difficult. The combination of low temperatures and high fields does provide high polarization, but such instrumentation⁶ can be prohibitively costly, and cooling below helium temperatures can cause increasingly slow T_1 relaxation.^{7,8}

Nonthermal ways to polarize spin sidestep these issues, and optical methods have proven especially promising.^{9–11} These techniques leverage electronic excitation and molecular design to influence spin, allowing chemists to untether spin polarization from Boltzmann population and spin T_1 relaxation. We are interested in a spin-sensitive optical phenomenon known as “optically induced magnetization” (OIM). OIM is deeply intertwined with a technique more familiar to chemists, magnetic circular dichroism (MCD): MCD measures the differential response of M_S levels to circularly polarized (CP) light ($\Delta\epsilon = \epsilon_{\text{LCP}} - \epsilon_{\text{RCP}}$),^{12,13} whereas OIM exploits this differential response to preferentially bleach one M_S level. Despite demonstration of OIM in a wide variety of systems,^{14–18} the method has not been well-explored to manipulate molecular spin. This will surely change after a recent report demonstrating the usefulness of OIM to explore ultrafast spin dynamics of K_2IrCl_6 .¹⁹

Selection of K_2IrCl_6 relied on its octahedral geometry and its $^2T_{2g}$ ground state (GS) orbital angular momentum (OAM), promoting a large MCD/OIM response but also causing exceedingly short coherence (phase-memory) T_m times, even at 20 K.¹⁹ Such rapid decoherence impairs many applications in quantum technologies by limiting the computations that can be performed or the information that may be sensed by a qubit.^{2,3} Herein, we report a more flexible molecular design strategy to optimize OIM responses that allows us to explore the preparation of $S = 1/2$ molecules of axial symmetry that lack any GS OAM. Our design strategy led us to a series of tungsten(V) chalcogenide complexes $[\text{E}]^-$ (Figure 1a) that could give up to ~20% spin polarization by OIM. Pulsed electron paramagnetic resonance (EPR) experiments reveal improved spin relaxation characteristics for these complexes due to their lack of GS OAM or a low-lying excited state (ES), and field-swept electron spin-echo (FS-ESE)²¹ EPR measurements reveal a remarkable insensitivity of phase-memory (T_m) times to magnetic field.

Using the language of MCD theory, maximization of OIM is equivalent to maximization of the “C term” MCD response, something driven by spin–orbit coupling (SOC) in transition-metal systems.²² The K_2IrCl_6 system experienced strong SOC due to the residual (“fictitious”) OAM in its GS, but there is no reason SOC splitting must come from the GS. Instead, we were inspired by ruby, where OIM can be performed from its

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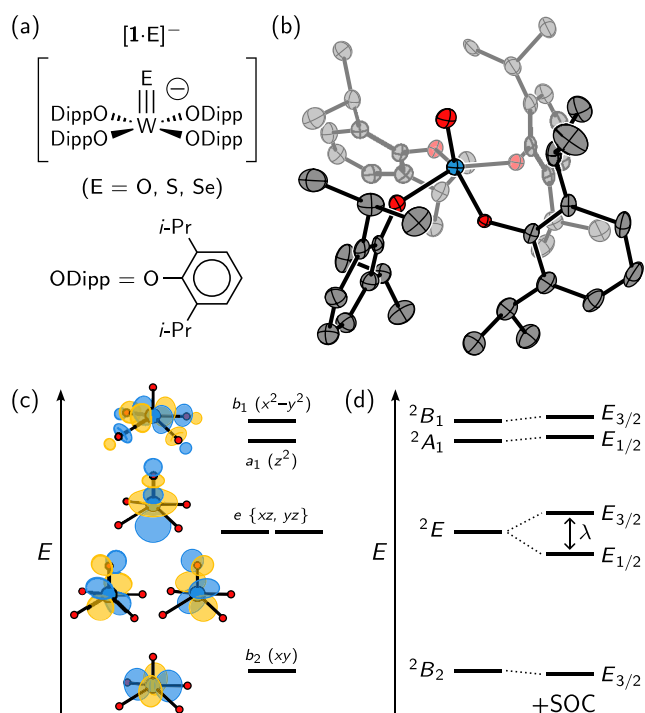


Figure 1. (a) The $[1\cdot E]^-$ compounds have bulky aryloxide ligands that enforce pseudo- C_{4v} geometry, confirmed by (b) the crystal structure of $[\text{Na}(\text{THF})_6][1\cdot\text{O}]^-$. (c) The d orbital splittings for $[1\cdot\text{O}]^-$ are shown with C_{4v} irrep labels. (d) SOC splits the 2E state into two components labeled by C_{4v}^* double group irreps.²⁰

orbitally nondegenerate 4A_2 GS to a degenerate 2E ES.¹⁴ This 2E ES of ruby lacks residual OAM; yet, SOC still causes splitting into two Kramers doublets. Because SOC is the dominant driver of this splitting, spin and orbital angular momenta strongly mix to deliver a highly M_S -selective MCD/OIM response.²³ Clearly, threefold orbital degeneracy is not required for OIM.

Simply put, any point group possessing degenerate irreducible representations (irreps) may allow for orbital-based optimization of OIM. This certainly includes cubic point groups (O_h , T_d , etc.), but also all axial point groups of threefold-or-higher symmetry. At its simplest, OIM maximization should therefore incorporate

- axial ($C_{(n \geq 3)(v/h)}$, $D_{(n \geq 3)(h)}$, $D_{(n \geq 2)d}$, $S_{2n \geq 4}$) or isometric ($T_{(d/h)}$, $O_{(h)}$, $I_{(h)}$) point groups
- an optical transition involving an orbitally degenerate state
- SOC splitting of similar or greater magnitude to the line width of the electronic transition

¹⁸ An OIM-initializable qubit should additionally target increased spin coherence through

- avoidance of GS OAM and low-lying ESs
- a large absorptivity (ϵ) to allow more dilute samples

Of the few examples of molecular OIM we are aware of,^{18,19} every system fails to satisfy at least one of these five criteria.

We selected tungsten(V)-oxo complex $[1\cdot\text{O}]^-$ as a promising system: tetragonal d^1 oxos generally have a low-lying 2E ES (Figure 1c),²⁴ and $5d$ metals have strong SOC ($\zeta_{W(V)} = 3483 \text{ cm}^{-1}$).²⁵ Heating a THF solution of $\text{WOCl}_3(\text{THF})_2$ ²⁶ and sodium 2,6-diisopropylphenoxide (NaODipp , 4 equiv) overnight (69 °C, 18 h) provided deep

blue $[\text{Na}(\text{THF})_6][1\cdot\text{O}]^-$ in 50% recrystallized yield. A preliminary X-ray diffraction study (Figure 1b) revealed a nearly C_{4v} -symmetric anion with a geometry close to an ideal square pyramid ($\tau_5 = 0.052(6)$).²⁷ An optical $^2B_2 \rightarrow ^2E$ transition was observed in the UV-vis-NIR absorption spectrum ($16\,100 \text{ cm}^{-1}$, Figure 2a) with a prominent shoulder due to tungsten(V) SOC, and a relatively large absorptivity for a $d-d$ transition ($270 \text{ M}^{-1} \text{ cm}^{-1}$). Thus, $[1\cdot\text{O}]^-$ satisfied all our design requirements.

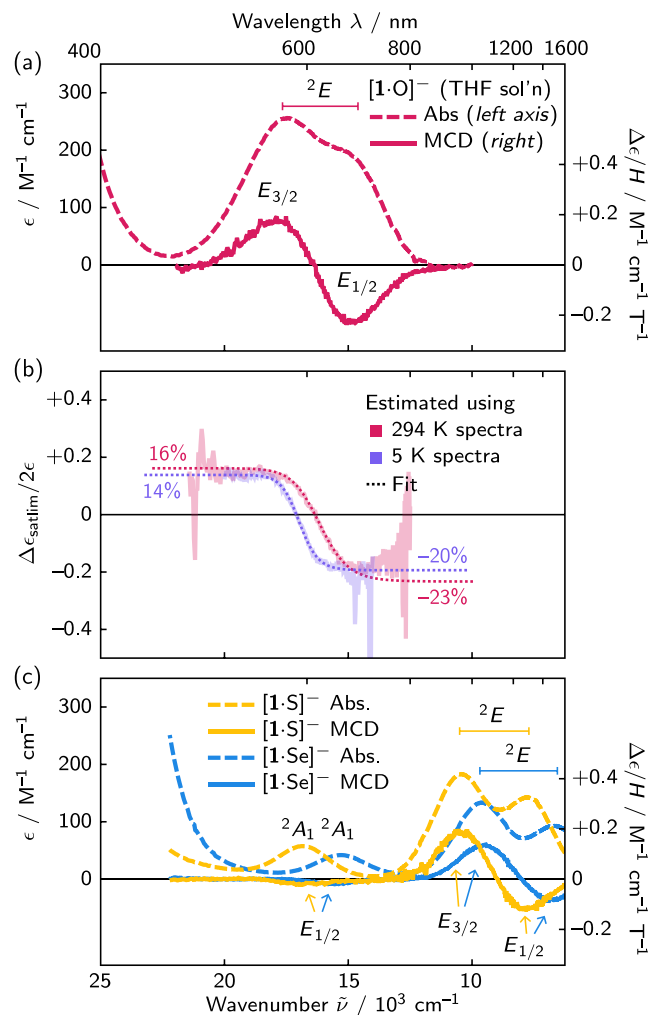
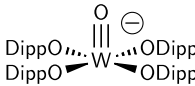
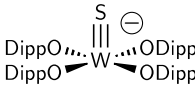
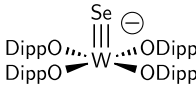
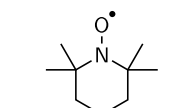
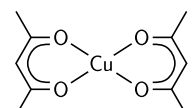
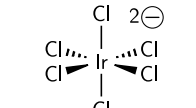


Figure 2. (a) Room-temperature UV-vis-NIR absorption and MCD spectra of $[\text{Na}(\text{THF})_6][1\cdot\text{O}]^-$ (THF solution). (b) The maximum spin polarization through OIM is $|\Delta\epsilon_{\text{satlim}}/2\epsilon|$, estimated here for $[1\cdot\text{O}]^-$ in two ways (see text and the SI). (c) The absorption and MCD spectra of the heavier chalcogenide complexes reveal a 2A_1 transition.

We were pleased to find the MCD spectrum of $[1\cdot\text{O}]^-$ shows large intensities associated with its 2E ES. The SOC splitting of this state was clearly indicated by the characteristic bisignate (“pseudo-A”) line shape, the negative and positive features corresponding to $E_{1/2}$ and $E_{3/2}$ levels, respectively (see the SI). The MCD intensities of these levels were quantified through their “ C_0/D_0 ” ratios, a ratio of MCD C term intensity (“ C_0 ”) to absorption intensity (“ D_0 ”). For a simple spin-1/2 GS lacking OAM and low-lying ESs, the maximum possible magnitude of $|C_0/D_0|$ is $g/2$, half the GS g value (see the derivation given in the SI); however, molecules typically feature ratios far smaller (Chart 1). The values for the $[1\cdot\text{O}]^-$

Chart 1. $^2B_2 \rightarrow ^2E$ Transitions of $[1\cdot E]^-$ Species Have Far Stronger MCD Intensities than Other $S = 1/2$ Organic or Transition Metal Species with Smaller SOC; Their C_0/D_0 Ratios Approach Those of $[\text{IrCl}_6]^{2-}$, a System with GS OAM

								
Solution in THF			Solution in THF			Solution in THF		
λ / nm	C_0/D_0		λ / nm	C_0/D_0		λ / nm	C_0/D_0	
676	−0.46(3)		1300	−0.39(3)		1520	−0.41(2)	
562	+0.32(4)		952	+0.44(1)		1040	+0.43(1)	
			595	−0.21(4)		654	−0.26(1)	

								
Solution in MeCN			Solution in DCM			Solution in Water		
λ / nm	C_0/D_0		λ / nm	C_0/D_0		λ / nm	C_0/D_0	
460	+0.040(4)		730	+0.08(1)		490	+0.65(2)	
			510	−0.13(1)		430	−0.68(1)	

$^2B_2 \rightarrow ^2E$ transitions (−0.46(3) and +0.32(4)) were quite large, multiple times larger than those of TEMPO $^\bullet$ and Cu(acac) $_2$ exemplars, and approach the values for $[\text{IrCl}_6]^{2-}$,^{13,28} an ion specifically chosen for OIM studies due to its residual GS OAM.¹⁹

The MCD C_0/D_0 ratio is a useful heuristic, but more direct insight for OIM would come from measuring $\Delta\epsilon$ for each individual M_S level. This can be accomplished in the saturation limit. Magnetic saturation occurs under sufficiently low temperatures or high fields that most molecules populate a single M_S level. In the limit of complete saturation, the MCD spectrum reveals the intrinsic $\Delta\epsilon$ for the single lowest-energy M_S level. MCD intensity in the presence of saturation is nonlinear,²⁹

$$\frac{\Delta\epsilon}{E} = \gamma \frac{2C_0}{g} \tanh\left(\frac{g\mu_B B}{2k_B T}\right) f(E) \quad (1)$$

where E is the excitation energy, γ a proportionality constant, μ_B the Bohr magneton, B the magnetic field, k_B the Boltzmann constant, T the temperature, and $f(E)$ a line shape function (see the SI). We used eq 1 to estimate MCD in the saturation limit $\Delta\epsilon_{\text{satlim}}$ in two ways: from room-temperature data, and from cryogenic MCD data (5 K, 7 T). The maximum spin polarization that can be achieved by OIM is $|\Delta\epsilon_{\text{satlim}}/2e|$ (see the SI), and both methods of estimating $\Delta\epsilon_{\text{satlim}}$ reveal that CP excitation of the $^2B_2 \rightarrow ^2E$ transition can deliver up to 15% ($E_{3/2}$) or 20% ($E_{1/2}$) spin polarization, depending on wavelength of excitation (Figure 2b). For context, such a large polarization can only be achieved thermally at 1–1.5 K in X-band EPR experiments.

Analogous chalcogen congeners $[1\cdot\text{S}]^-$ and $[1\cdot\text{Se}]^-$ were also synthesized to explore their MCD intensities. The former was prepared by treating $\text{W}(\text{ODipp})_4$ ³⁰ with NaSCPh_3 (1 equiv, THF), a convenient source of $\text{S}^{\bullet-}$ upon release of $\text{CPh}_3^\bullet$. Stirring overnight then washing with ether provided teal $[\text{Na}(\text{THF})_6][1\cdot\text{S}]$ in 36% yield after recrystallization. The latter complex was prepared in a two-step one-pot procedure: $\text{W}(\text{ODipp})_4$ was first treated with PPh_3Se to form the terminal selenide complex, and then reduction with NaCPh_3 formed the tungsten(V) anion. Recrystallization provided vibrant green $[\text{Na}(\text{THF})_x(\text{Et}_2\text{O})_y][1\cdot\text{Se}]$ in 26% isolated yield.

The MCD spectra of $[1\cdot\text{S}/\text{Se}]^-$ show the appearance of the $^2B_2 \rightarrow ^2A_1$ ligand field transition, likely caused by weaker $\sigma_{\text{W}\equiv\text{E}}$ interactions in the heavier chalcogenides. This assignment follows from its negative MCD response, indicating $E_{1/2}$ symmetry, and is corroborated by model CASSCF/RI-NEVPT2 calculations (see the SI).³¹ We ascribe the weaker C_0/D_0 ratios (Chart 1) of the $^2B_2 \rightarrow ^2A_1$ transitions to a lack of ES orbital degeneracy. For such nondegenerate states, the sum-over-states perturbation theory formalism²² says C term intensity is proportional to the reciprocal energy difference to other states coupling through SOC. This is an important distinction from 2E states, underscoring that design of molecules with large MCD/OIM responses should target transitions involving orbital degeneracy.

All three chalcogenide complexes can be observed by continuous-wave (CW) EPR experiments at temperatures at least up to 80 K due to their lack of GS OAM and thus slower

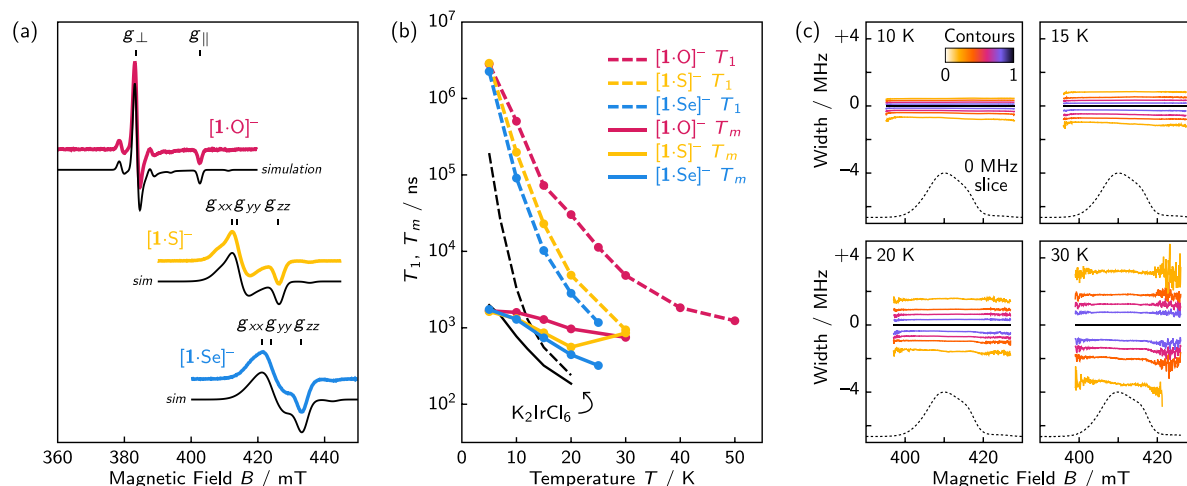


Figure 3. (a) The CW EPR spectra show approximate C_4 symmetry for all anions with slight rhombic distortions for $[1\cdot\text{S}/\text{Se}]^-$. (b) Spin T_1 and T_m measurements show slower relaxation for $[1\cdot\text{E}]^-$ anions than K_2IrCl_6 (values from ref 19) due to the latter's GS OAM. (c) Normalized contours in the FS-ESE spectrum of $[1\cdot\text{Se}]^-$ show nearly flat T_m times independent of field.

relaxation properties. Inversion recovery and Hahn-echo X-band EPR experiments were used to directly measure T_1 and T_m times (Figure 3), giving values higher than those of K_2IrCl_6 (10–100 times higher at 20 K);¹⁹ in fact, even higher values may be attainable through techniques like deuteration,³² picket-fence saturation recovery experiments,³³ and Carr–Purcell–Meiboom–Gill (CPMG) pulse sequences.³⁴ Despite this increase in coherence for $[1\cdot E]^-$ complexes, these T_m relaxation times are still modest likely due to the strong SOC within the systems.³⁵ A balance clearly needs to be struck between maximizing MCD/OIM using heavy atom SOC while avoiding the collapse of spin coherence, and this will be a fascinating area for future exploration.

The T_1 and T_m measurements also showed unexpectedly consistent relaxation times when measured in the $g_{||}$ and $g_{\perp}$ regions of the spectra. Such behavior differs from related species,³⁶ which often show large variations across the spectra.³⁷ This insensitivity to magnetic field was demonstrated using a 2D FS-ESE EPR experiment²¹ conducted on $[1\cdot Se]^-$, showing nearly constant T_m times across the entire spectrum. Studies are ongoing to uncover correlations between relaxation times, molecular symmetry, and MCD/OIM intensity.

With this work, we hope to highlight that a large synthetic space remains to be explored in the development of optimal molecules for OIM. Our results ease the restrictive O_h design strategy for OIM, showing that chemists can exploit axial symmetry to deliver large MCD responses. This increased flexibility should allow targeting of other properties important for qubits, especially spin relaxation times. Ligand design may be particularly useful, as the ODipp ligands here were selected for ease of use³⁰ rather than promotion of spin coherence.^{38–40} We are conducting experiments toward these ends, and toward balancing spin and optical lifetimes for future implementation of OIM in molecular qubit systems on the nanosecond time scale.⁴¹

■ ASSOCIATED CONTENT

SI Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/jacs.5c03783>.

Experimental details, spectroscopic characterization, crystallographic information, computational details, and theoretical derivations (PDF)

Accession Codes

Deposition Numbers 2428564–2428566 and 2444691 contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via the joint Cambridge Crystallographic Data Centre (CCDC) and Fachinformationszentrum Karlsruhe [Access Structures service](#).

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Notes

The authors declare no competing financial interest.

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