

SATURATION OF THE DRELL-HEARN-GERASIMOV SUM RULE*

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Abstract

The Drell-Hearn-Gerasimov sum rule for the forward spin-flip amplitude of nucleon Compton scattering is decomposed into separate sum rules originating from different isospin components of the electromagnetic current. The resulting sum rules are reexamined using recently available analyses of single pion photoproduction in the region up to photon laboratory energies of 1.2 GeV. All three sum rules receive important non-resonant as well as resonant contributions. The isovector sum rule whose contributions are known best is found to be nearly saturated, lending support to the assumptions underlying the sum rules. The failure of the isoscalar-isovector sum rule to be saturated is then presumably to be blamed on inadequate data for inelastic contributions.

(Submitted to Phys. Rev.)

*Work supported by the U. S. Atomic Energy Commission.

I. INTRODUCTION

The Drell-Hearn-Gerasimov¹ sum rule for the spin-flip amplitude in forward Compton scattering rests on two assumptions. These are the low energy theorem² for the spin-flip amplitude and the validity of an unsubtracted dispersion relation. Since these are simple and relatively well accepted assumptions, and are often used together with additional stronger assumptions in deriving other sum rules, it is of interest to look into the validity of the Drell-Hearn-Gerasimov sum rule in the light of the present experimental data.

In their original paper, Drell and Hearn¹ did attempt to investigate the validity of the sum rule for a proton target by using an isobar model of single pion photoproduction. Their results were generally encouraging, but some important contribution from high energy (greater than 1 GeV) seemed to be likely. Somewhat later, Chau *et al.*³ extended the examination of the proton sum rule by using an analysis of single pion photoproduction through the second resonance region. They found good agreement, without any high energy contribution. Finally, in the course of an analysis of many sum rules, Fox and Friedman⁴ have considered the Drell-Hearn-Gerasimov sum rule, using Walker's partial wave analysis⁵ of pion photoproduction. They found the somewhat surprising result that while the sum rule involving only the isovector part of the electromagnetic current appeared well satisfied, the sum rule involving one isovector and one isoscalar current (equivalent to the difference of proton and neutron sum rules) was badly violated.

Since that time, there has been a considerable improvement in both the pion photoproduction data and their analysis. In particular, relatively good neutron data are becoming available and have been incorporated in the recent analyses of Pfeil and Schwela⁶ and Moorhouse and Oberlack.⁷

Given this changed situation, we reexamine in this paper the Drell-Hearn-Gerasimov sum rules for both proton and neutron targets, with particular attention to their difference. In the next section, we give the relevant definitions and present the contributions to the sum rules using several recent analyses of pion photoproduction. In the third section, we present some conclusions.

II. ANALYSIS OF CONTRIBUTIONS

The unsubtracted dispersion relation for the forward spin-flip nucleon Compton amplitude f_2 :

$$\text{Re } f_2(\nu) = \frac{2\nu}{\pi} \int_{\nu_0}^{\infty} \frac{d\nu'}{\nu'^2 - \nu^2} \text{Im } f_2(\nu')$$

gives the Drell-Hearn-Gerasimov¹ sum rule

$$\frac{2\pi^2\alpha}{M^2} \kappa^2 = \int_{\nu_0}^{\infty} \frac{\sigma_{3/2}(\nu') - \sigma_{1/2}(\nu')}{\nu'} d\nu' \equiv I_p \quad (1)$$

when the low energy theorem is applied to $f_2(\nu)$. Here κ is the nucleon's anomalous magnetic moment, M the nucleon's mass, ν_0 the threshold energy for a single pion photoproduction, which in the laboratory frame equals:

$$\nu_0 = \frac{m_\pi^2}{2M} + m_\pi \approx 150 \text{ MeV} ,$$

and $\sigma_{3/2}(\sigma_{1/2})$ is the total cross section for the process: photon + nucleon \rightarrow hadrons in the net helicity state 3/2 (1/2). These cross sections enter the sum rule through the unitarity equation which relates the imaginary part of the forward scattering amplitude to the total cross section into the intermediate states.

We decompose the left- and right-hand side of Eq. (1) into components of different isospin character and then relate the corresponding parts. The anomalous magnetic moments of the proton and neutron are defined:

$$\kappa_{p(n)} = \frac{1}{2} \kappa_s \pm \frac{1}{2} \kappa_v \quad (2)$$

where κ_s (κ_v) is the isoscalar (isovector) component.

In this way we obtain the isovector, isoscalar, and the "interference" Drell-Hearn-Gerasimov sum rules:

$$\begin{aligned} I_{VV} &\equiv \int_{\nu_0}^{\infty} \left[\sigma_{3/2}^{VV}(\nu') - \sigma_{1/2}^{VV}(\nu') \right] \frac{d\nu'}{\nu'} = \left(\frac{\kappa_v}{2} \right)^2 \frac{2\pi^2 \alpha}{M^2} = 218.5 \mu b \\ I_{SS} &\equiv \int_{\nu_0}^{\infty} \left[\sigma_{3/2}^{SS}(\nu') - \sigma_{1/2}^{SS}(\nu') \right] \frac{d\nu'}{\nu'} = \left(\frac{\kappa_s}{2} \right)^2 \frac{2\pi^2 \alpha}{M^2} = 0.3 \mu b \\ I_{VS} &\equiv \int_{\nu_0}^{\infty} \left[\sigma_{3/2}^{VS}(\nu') - \sigma_{1/2}^{VS}(\nu') \right] \frac{d\nu'}{\nu'} = \frac{\kappa_v \kappa_s}{4} \frac{2\pi^2 \alpha}{M^2} = -14.7 \mu b \end{aligned} \quad (3)$$

Presently, the data limit the study of the saturation of these sum rules to the contributions (to the total cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$) of a nucleon and one pion final hadronic state. Resonance dominance only allows one to estimate a part of the inelastic contributions. Still, one may well hope that the largest contributions come from not too far from threshold, so that the nucleon plus one pion state at least provides an indication of the sum rules saturation.

Single pion photoproduction amplitudes have a simple isospin decomposition⁵:

$$M_{\gamma p \rightarrow \pi^+ n}^+ = \sqrt{\frac{1}{3}} \left[M^{(3)} - \sqrt{2} (M^{(1)} - M^{(0)}) \right]$$

$$M_{\gamma p \rightarrow \pi^0 p}^0 = \sqrt{\frac{2}{3}} \left[M^{(3)} + \frac{1}{\sqrt{2}} (M^{(1)} - M^{(0)}) \right] \quad (4)$$

$$M_{\gamma n \rightarrow \pi^- p}^- = \sqrt{\frac{1}{3}} \left[M^{(3)} - \sqrt{2} (M^{(1)} + M^{(0)}) \right] .$$

$M^{(3)}$ corresponds to the isospin 3/2 state and $M^{(1)}$ to the isospin 1/2 state created by the isovector part of the photon current; $M^{(0)}$ describes the interaction of the nucleon with the isoscalar part of this current.

The cross sections of a definite isospin character for one-pion photoproduction are then proportional to the following combinations of amplitudes:

$$\begin{aligned} \sigma^{VV} &\propto |M^{(3)}|^2 + |M^{(1)}|^2 \\ \sigma^{SS} &\propto |M^{(0)}|^2 \\ \sigma^{VS} &\propto - \left[(M^{(0)})^* M^{(1)} + M^{(0)} (M^{(1)})^* \right] . \end{aligned} \quad (5)$$

Additionally, the cross sections of a definite helicity in this process are conveniently expressed in terms of $A_{n\pm}$, $B_{n\pm}$ amplitudes^{8,9}:

$$\begin{aligned} \sigma_{1/2} &= \frac{8\pi q}{k} \sum_{n=0}^{\infty} (n+1) \left(|A_{n+}|^2 + |A_{(n+1)-}|^2 \right) \\ \sigma_{3/2} &= \frac{8\pi q}{k} \sum_{n=0}^{\infty} \frac{n(n+1)(n+2)}{4} \left(|B_{n+}|^2 + |B_{(n+1)-}|^2 \right) . \end{aligned}$$

The above relations allow us to consider the isovector, isoscalar, and interference sum rules separately, to insert the single pion plus nucleon intermediate states' cross sections into the integrals, and to separate the contributions due to each partial wave.¹⁰ Since we can extract the contributions of definite isospin and of definite total and orbital angular momentum, we are able to

separate the contributions due to resonances of known inelasticity and subsequently to evaluate the corresponding part of the inelastic contributions to the integrals.

A. The Isovector Sum Rule

Table I presents the contributions to the isovector, I_{VV} , sum rule. For the photon laboratory energy, ν_{lab} , below .45 GeV we find good agreement¹¹ between the results obtained with the analyses of ref. 5, 6, and 7. Since the Pfeil-Schwela fit extends to the lowest energy, we have chosen to list the contributions obtained with their analysis, where possible.¹² For ν_{lab} above .45 GeV the values resulting from ref. 5 and 7 also agree well, and listed are contributions obtained with the very recent Moorhouse-Oberlack fit.

The inelastic part is evaluated as a sum of inelastic contributions of N(1520), N(1670), and N(1688), (D_{13} , D_{15} , and F_{15} , respectively).

The behavior of $\sigma_{3/2}^{VV} - \sigma_{1/2}^{VV}$ as a function of energy is shown in Fig. 1. One can easily recognize on this graph the dominant features of a big negative non-resonant s-wave contribution and a large positive contribution from the P_{33} (1236) and the two other resonances.

B. The Isoscalar Sum Rule

Table II presents the results for the isoscalar Drell-Hearn-Gerasimov sum rule. There are serious discrepancies between the results obtained with different fits; the signs of various contributions even disagree between analyses. This difficulty has been previously noted by Fox and Friedman.⁴ The isoscalar amplitudes cannot be extracted directly from the data. Instead, they are obtained indirectly from the sum and difference of the proton and neutron data, which results in relatively large errors, as the isoscalar amplitudes themselves are small compared to the isovector ones.

The behavior of $\sigma_{3/2}^{ss} - \sigma_{1/2}^{ss}$ in Fig. 2 combines the Pfeil-Schwela and the Moorhouse-Oberlack results. There is no systematic trend and further data are necessary, both more accurate and including higher energies, in order to evaluate this sum rule.

C. The Interference Sum Rule

Table III presents the results for the isovector-isoscalar interference sum rule. This part corresponds to $I = 1$ exchange in the t -channel, or in other words, to the difference between the proton and neutron Drell-Hearn-Gerasimov sum rules. It should have a negative value, being proportional to $\kappa_p^2 - \kappa_n^2$.

We find agreement between the results obtained with different analyses, where they overlap, and present the values resulting from the Pfeil-Schwela fit for $\nu_{lab} \leq .45$ GeV,¹² and from the Moorhouse-Oberlack fit for $\nu_{lab} \geq .45$ GeV.

Each of these fits results in a big, non-resonant, s -wave contribution of the wrong, positive sign. This contribution, however, is cancelled almost entirely by the non-resonant part of 1^+ partial wave in the Moorhouse-Oberlack analysis. The second and third resonance cause the total result for this sum rule to be of wrong, positive sign, which reflects the stronger coupling of these resonances to the proton than to the neutron. Figure 3 shows this behavior in terms of $\sigma_{3/2}^{vs} - \sigma_{1/2}^{vs}$.

III. DISCUSSION AND CONCLUSIONS

The Drell-Hearn-Gerasimov sum rule rests on two assumptions: The low energy theorem and the validity of an unsubtracted dispersion relation for the f_2 (spin-flip) Compton amplitude. The first is quite general and has a very solid theoretical basis. Therefore, the validity of the sum rule is presumably dependent on the second assumption. This one is equivalent to the absence of a fixed pole (with $J^P = 1^+$) in the f_2 amplitude.¹³

Strong evidence for the correctness of this hypothesis comes from the evaluation of the contributions to the isovector sum rule, shown in Table I. This sum rule seems to be rather well saturated by the results obtained with the single pion plus nucleon data, complemented by some estimate of the inelastic contributions, up to $\nu_{\text{lab}} = 1.2$ GeV. Even more pleasing is that this saturation occurs in a non-trivial way. There are strong cancellations, mainly between the large negative s-wave and the positive P (1236) contributions (see Fig. 1). The large non-resonant s-wave contribution is of interest in itself, however, as it violates two-component duality. The imaginary part of a non-diffractive amplitude like $f_2(\nu)$ should contain only s-channel resonances.

Our total numerical result for the isovector sum rule is in general agreement with previous analyses.^{3,4} While contributions from still higher values of ν_{lab} need not be small, we expect the contributions listed in Table I to be the largest individual ones, particularly since the sum rule integrands involve the factor $1/\nu_{\text{lab}}$ times a difference of total cross sections which is expected to vanish at high energies.

The isoscalar sum rule presents some problems. Because of the sensitivity of the isoscalar amplitudes to the relatively small differences between the neutron and proton photoproduction data, it is very difficult to achieve reliable values for the isoscalar contributions. Furthermore, Table II shows that individual, non-resonant, partial waves make (cancelling) contributions, each of which is of the order of the expected total value. In such a situation, small shifts in the data or the contributions of inelastic states may easily remove the present disagreement between the total value shown in Table II and that predicted (I_{ss}) by Eq. (3).

In this light, the behavior of the isovector-isoscalar sum rule is puzzling. Since the isovector sum rule is almost saturated, we have every reason to expect

the validity of the underlying assumptions for the interference sum rule as well. The total contribution in Table III is of the wrong sign, however. This general difficulty had been previously noted by Fox and Friedman⁴ using an earlier analysis⁵ of pion photoproduction. We note in particular that there are large non-resonant contributions in the 0^+ and 1^+ partial waves, which tend to cancel. The second and third resonances contribute to the sum rule with the wrong sign. This again is of interest in itself as it violates local two-component duality. Global duality would still seem to be satisfied for this part of $f_2(\nu)$, because the various non-resonant partial waves are cancelling in the full amplitude. This leads one to believe that if the sum rule is to work, it may well be the contributions of quite inelastic resonances, many in low partial waves, that saturate the sum rule. Unfortunately, the determination of these contributions would be very difficult experimentally.

In summary, we find no reason to doubt the validity of the Drell-Hearn-Gerasimov sum rule. The isovector sum rule's near-saturation even furnishes some direct evidence of support. There seems little reason to be alarmed at the non-saturation of the isoscalar and isovector-isoscalar sum rules at the present stage of photoproduction analysis. What is needed is the more direct experimental determination of $\sigma_{1/2}(\nu)$ and $\sigma_{3/2}(\nu)$, using a polarized beam and target, something which is now becoming a real possibility.

Acknowledgments

I would like to thank Professor F. J. Gilman for the constant help and encouragement given me at all stages of work on this paper, and R. G. Moorhouse and H. Oberlack for making their results available before publication.

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7. R. G. Moorhouse and H. Oberlack, Contribution to the 16th International Conference on High Energy Physics, Chicago, Illinois, September 6-13, 1972, and private communication.
8. $A_{n\pm}$ and $B_{n\pm}$ correspond to a state with final pion orbital momentum n , definite parity $P = (-1)^{n+1}$ and total angular momentum $j = n \pm 1/2$.
9. See also, for example, F. J. Gilman, "Photoproduction and Electroproduction," Physics Reports 4C, 95 (1972).
10. For example, the single pion part of $\sigma_{3/2}^{vs}$ equals:
$$\sigma_{3/2}^{vs} = \frac{8\pi q}{k} \sum_{n=0}^{\infty} -\frac{n(n+1)(n+2)}{4} \left[\left(B_{n+}^{(0)} \right)^* B_{n+}^{(1)} + B_{n+}^{(0)} \left(B_{n+}^{(1)} \right)^* \right. \\ \left. + \left(B_{(n+1)-}^{(0)} \right)^* B_{(n+1)-}^{(1)} + B_{(n+1)-}^{(0)} \left(B_{(n+1)-}^{(1)} \right)^* \right]$$
11. I_{vv} for $\nu_{lab} \leq .45$ due to P_{33} (1236) has the following values resulting from different analyses: 260 μb with Walker's fit, 210 μb with the Moorhouse-Oberlack fit, and 237 μb with the Pfeil-Schwela fit.
12. The Pfeil-Schwela analysis is limited to the first resonance region and to 0^+ , 1^- , and 1^+ partial waves.
13. H. D. I. Abarbanel and M. L. Goldberger, Phys. Rev. 165, 1594 (1968).

Figure Captions

- Figure 1. Single pion photoproduction contribution to the difference between the photon-nucleon cross sections in the helicity $3/2$ and $1/2$ states $(\sigma_{3/2} - \sigma_{1/2})$ for the isovector photons.
- Figure 2. As in Fig. 1 but for the isoscalar photons.
- Figure 3. As in Fig. 1 but for the interference term between the isoscalar and isovector photons.

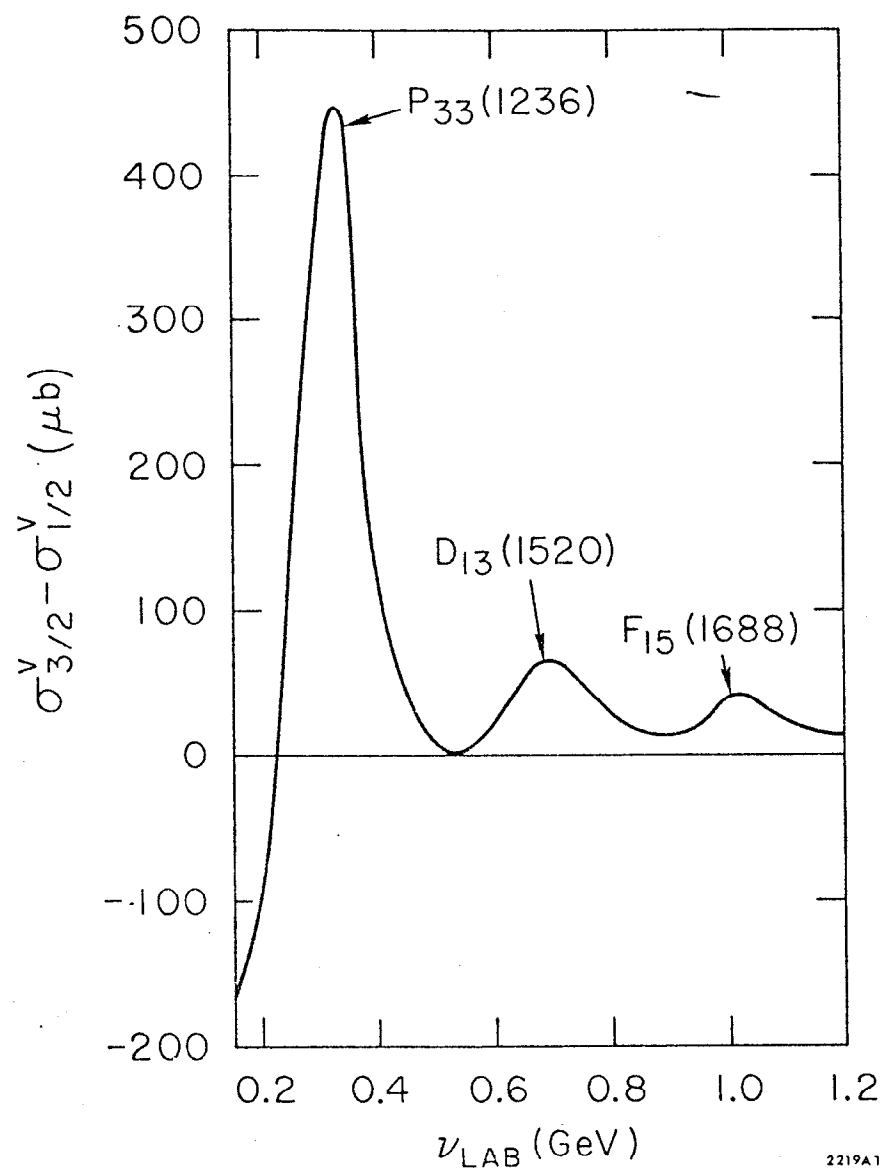
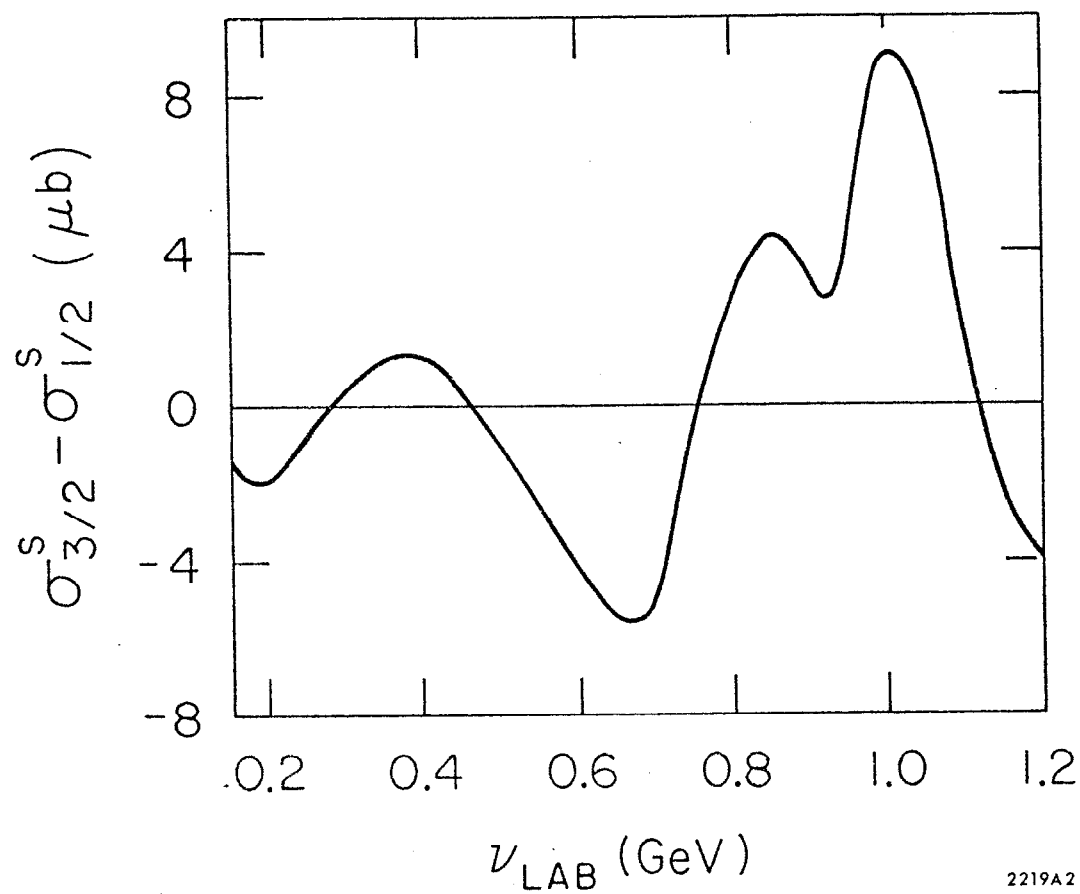
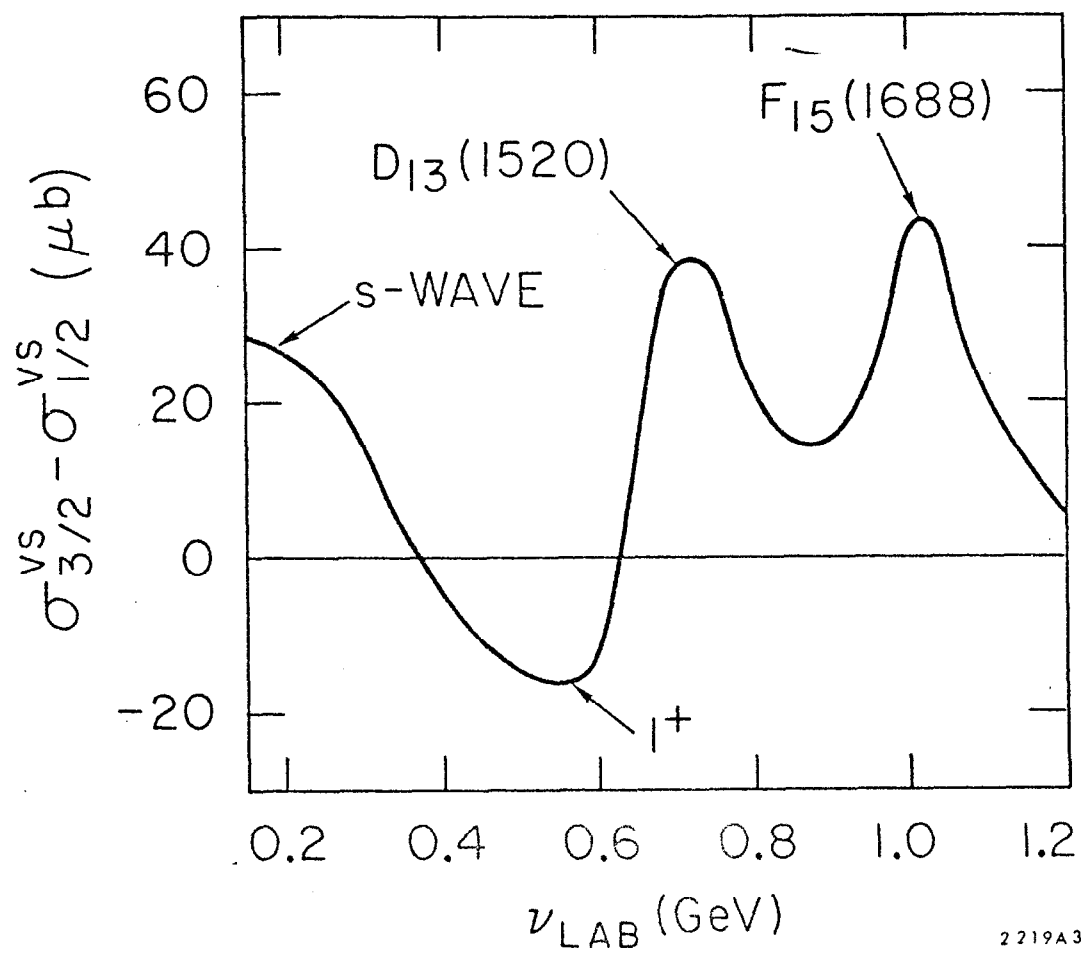


FIG. 1



2219A2

FIG. 2



2219A3

FIG. 3

Table I
The Isovector-Isovector Contributions
to the Drell-Hearn-Gerasimov Sum Rule¹²

Partial Wave	$.18 \text{ GeV} \leq \nu_{\text{lab}} \leq .45 \text{ GeV}$ (μb)	$.45 \text{ GeV} \leq \nu_{\text{lab}} \leq 1.2 \text{ GeV}$ (μb)	Total (μb)
0^+	-115 ^a	-51 ^b	-166
1^-	-2 ^a	-9 ^b	-11
1^+ $\left\{ \begin{array}{l} I=3/2; \text{ includes} \\ \quad P_{33}(1236) \\ I=1/2 \end{array} \right.$	+237 ^a +12 ^a	+2 ^b +17 ^b	+239 +29
2^- $\left\{ \begin{array}{l} I=3/2 \\ I=1/2; \text{ includes} \\ \quad D_{13}(1520) \end{array} \right.$	+7 ^b +13 ^b	+5 ^b +41 ^b	+12 +54
2^+ $\left\{ \begin{array}{l} I=3/2 \\ I=1/2; \text{ includes} \\ \quad D_{15}(1670) \end{array} \right.$	+1 ^b +2 ^b	-2 ^b +1 ^b	-1 +3
3^- $\left\{ \begin{array}{l} I=3/2 \\ I=1/2; \text{ includes} \\ \quad F_{15}(1688) \end{array} \right.$	+1 ^b +2 ^b	+1 ^b +7 ^b	+2 +9
Inelastic	+12	+37 ^b	+49
Total	+170	+49 ^b	+219

a - Pfeil and Schwela, ref. 6

b - Moorhouse and Oberlack, ref. 7

Table II
The Isoscalar-Isoscalar Contributions
to the Drell-Hearn-Gerasimov Sum Rule¹²

Partial Wave	$.18 \text{ GeV} \leq \nu_{\text{lab}} \leq .45 \text{ GeV}$ (μb)	$.45 \text{ GeV} \leq \nu_{\text{lab}} \leq 1.2 \text{ GeV}$ (μb)	Total (μb)
0^+	-1.07^a	—	
	$-.80^c$		-1.30^c
	-1.93^b		-2.99^b
1^-	$-.33^a$		
	$-.29^c$		$-.63^c$
	$-.26^b$		-1.87^b
1^+	$+.30^a$		
	$-.28^c$		$-.88^c$
	$+1.77^b$		$+4.39^b$
2^- (includes $D_{13}(1520)$)	$+.30^c$	$-.21^c$	$+.09^c$
	$+1.30^b$		$-.45^b$
2^+ (includes $D_{15}(1670)$)	$+.01^c$	$+.04^c$	$+.05^c$
	$-.10^b$		$+.15^b$
3^- (includes $F_{15}(1688)$)	$+.25^c$	$+1.38^c$	$+1.63^c$
	$+.01^b$		$+2.50^b$
Inelastic	$+.41^c$	$+.62^c$	$+1.03^c$
	$+.95^b$		$+1.19^b$
Total	$-.40^c$	$+.39^c$	$-.01^c$
	$+1.74^b$		$+2.92^b$

a - Pfeil-Schwela analysis, ref. 6

b - Moorhouse-Oberlack analysis, ref. 7

c - Walker analysis, ref. 5

Table III
The Isovector-Isoscalar Contributions
to the Drell-Hearn-Gerasimov Sum Rule¹²

Partial Wave	$.18 \text{ GeV} \leq \nu_{\text{lab}} \leq .45 \text{ GeV}$ (μb)	$.45 \text{ GeV} \leq \nu_{\text{lab}} \leq 1.2 \text{ GeV}$ (μb)	Total (μb)
0^+	+17 ^a	+5 ^b	+22
1^-	+1 ^a	+2 ^b	+3
1^+	-6 ^a	-17 ^b	-23
2^- (includes D_{13} (1520))	-3 ^b	+16 ^b	+13
2^+ (includes D_{15} (1670))	0	+1 ^b	+1
3^- (includes F_{15} (1688))	0	+8 ^b	+8
Inelastic	-2	+17	+15
Total	+7	+32	+39

a - Pfeil-Schwela analysis, ref. 6

b - Moorhouse-Oberlack analysis, ref. 7