

## Synthesis of the New Element with $Z=117$

J. H. Hamilton<sup>1</sup>, Yu. Ts. Oganessian<sup>2</sup>, F. Sh. Abdullin<sup>2</sup>, P. D. Bailey<sup>3</sup>, D. Benker<sup>3</sup>, M. E. Bennett<sup>4</sup>, S. N. Dmitriev<sup>2</sup>, J. Ezold<sup>3</sup>, R. A. Henderson<sup>5</sup>, M. G. Itkis<sup>2</sup>, Yu. V. Lobanov<sup>2</sup>, A. N. Mezentsev<sup>2</sup>, K. J. Moody<sup>5</sup>, S. L. Nelson<sup>5</sup>, A. N. Polyakov<sup>2</sup>, C. E. Porter<sup>3</sup>, A. V. Ramayya<sup>1</sup>, F. Riley<sup>3</sup>, J. B. Roberto<sup>3</sup>, M. A. Ryabinin<sup>6</sup>, K. P. Rykaczewski<sup>3</sup>, R. N. Sagaidak<sup>2</sup>, D. A. Shaughnessy<sup>5</sup>, I. V. Shirokovsky<sup>2</sup>, M. A. Stoyer<sup>5</sup>, V. G. Subbotin<sup>2</sup>, R. Sudowe<sup>4</sup>, A. M. Sukhov<sup>2</sup>, Tu. S. Tsyganov<sup>2</sup>, V. K. Utyonkov<sup>2</sup>, A. A. Voinov<sup>2</sup>, G. K. Vostokin<sup>2</sup>, and P. A. Wilk<sup>5</sup>

<sup>1</sup>Department of Physics and Astronomy, Vanderbilt University, USA

<sup>2</sup>Joint Institute for Nuclear Research, RU-141980, Dubna, Russia Federation

<sup>3</sup>Oak Ridge National Laboratory, Oak Ridge, TN 37831, USA

<sup>4</sup>University of Nevada Las Vegas, Las Vegas, NV 89154, USA

<sup>5</sup>Lawrence Livermore National Laboratory, Livermore, CA 94551, USA

<sup>6</sup>Research Institute of Atomic Reactors, RU-433510 Dimitrovgrad, Russia Federation  
j.h.hamilton@vanderbilt.edu

**Abstract.** The synthesis of the new chemical element with atomic number  $Z=117$  is presented. The isotopes  $^{293}117$  and  $^{294}117$  were produced in fusion reactions between  $^{48}\text{Ca}$  and  $^{249}\text{Bk}$ . The  $^{249}\text{Bk}$  was produced in the High Flux Isotope Reactor and chemically separated at Oak Ridge. Decay chains involving eleven new nuclei were identified by means of the Dubna Gas Filled Recoil Separator. The measured decay properties show a strong rise of stability for super-heavy nuclei toward  $N=184$ .

### 1. Introduction

The stability of superheavy elements (SHE) is strongly influenced by the shell structure of neutrons and protons, see e.g., [1,2]. In the absence of corrections for shell structure, the liquid drop model predicts nuclei with  $Z \geq 100$  should not exist. However, isotopes of elements having atomic numbers up to  $Z=118$  have been observed [1]. Indeed, beyond the especially stable, spherical double-magic  $^{208}\text{Pb}$  (double closed shells  $Z=82$  and  $N=126$ ), theoretical predictions are that  $N=184$  should be spherical magic number that would give special stability to nuclei. The proton magic number is predicted in different approaches to be 114, 120, and 126 to form an Island of Stability. The synthesis of new elements with neutron number ( $N$ ) approaching 184 provide important tests of the nuclear structure models used to predict closed spherical shells in the heaviest elements.

Reactions between doubly-magic  $^{208}\text{Pb}$  and singly-magic  $^{209}\text{Bi}$  target nuclei and stable neutron-rich projectiles such as  $^{64}\text{Ni}$  or  $^{70}\text{Zn}$  were used for the synthesis of new heavy elements. In these cold fusion reactions isotopes with  $Z \sim 113$  and  $N \sim 165$  [3, 4], stabilized by the  $Z=108$  and  $N=162$  shell gaps for deformed shapes were observed (see Fig. 1). The dramatic drop of the production cross section with increasing  $Z$  practically excludes the continuation of such experiments for heavier elements.

A new method of synthesizing superheavy elements, with  $Z \geq 112$  and neutron numbers closer to the predicted spherical shell closure at  $N=184$ , was pioneered at the Flerov Laboratory of Nuclear Reactions (FLNR) of Joint Institute for Nuclear Research (JINR) [1, 5]. Four new isotopes of

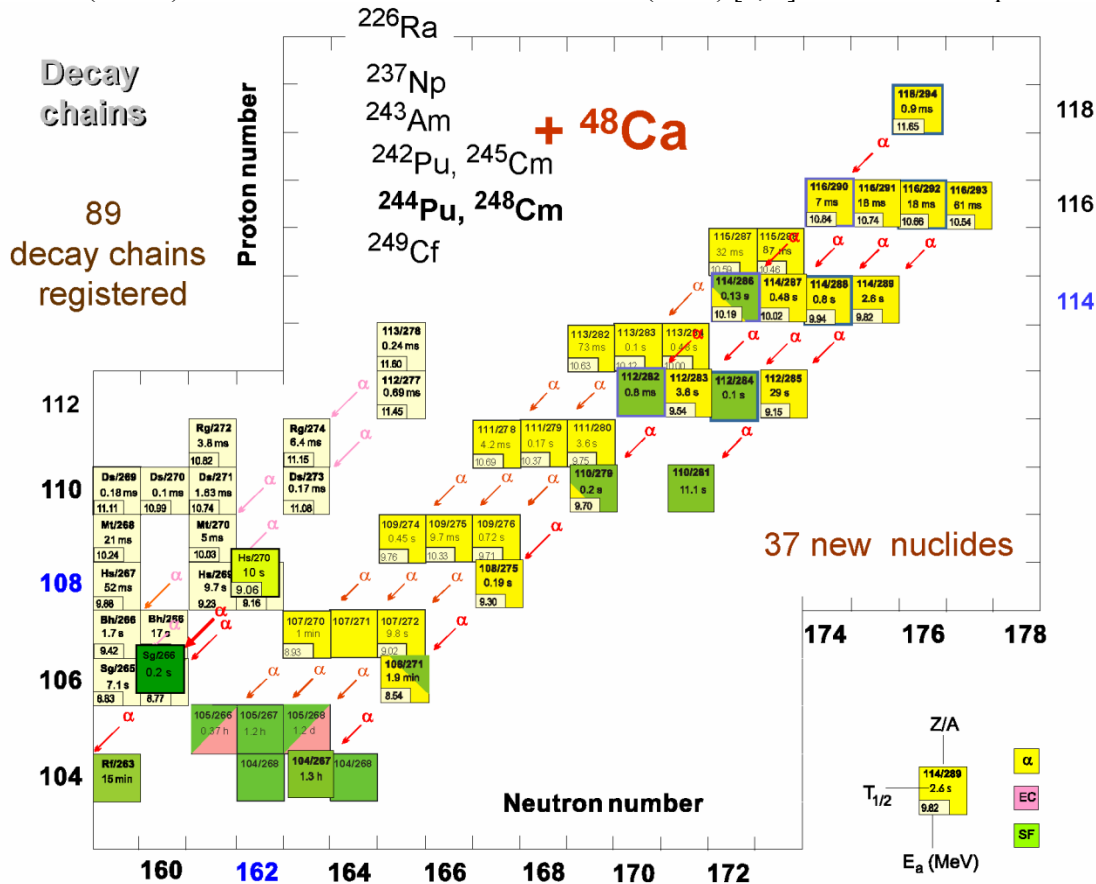


Fig. 1 Superheavy elements produced in cold fusion (pale yellow) [3, 4] and hot fusion (bright yellow on right) [1] reactions. The number of decay chains and nuclides were prior to this work.

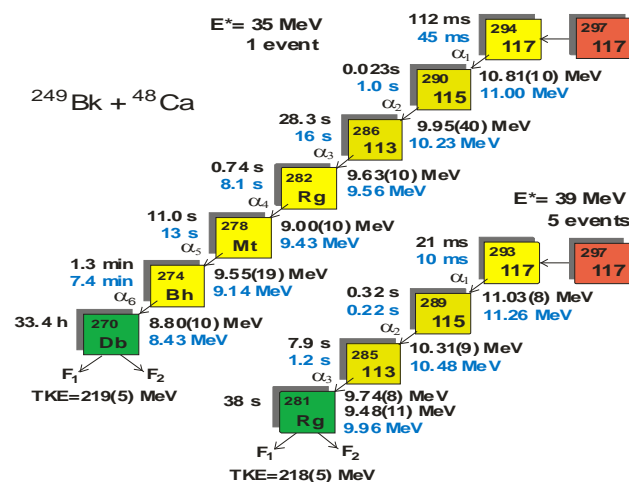


Fig. 2 Observed decay chains interpreted as originating from the isotopes  $A=294$  and  $A=293$  (average of five events) of the new element  $Z=117$  [10]. The deduced and predicted [11] lifetimes ( $\tau = T_{1/2}/\ln 2$ ) and  $\alpha$ -particle energies are shown in black and blue, respectively. Taken from [10].

element  $Z=112$  and fourteen isotopes of new elements with  $Z=113-116$  and  $118$  were identified [1] by using heavy-ion fusion reactions of doubly-magic  $^{48}\text{Ca}$  projectiles and actinide targets of U-Cm and Cf, respectively (see Fig. 1). The element with  $Z=117$  was missing because of difficulty in obtaining the short lived  $^{249}\text{Bk}$  for target material. The probabilities of formation and the decay properties of these 18 new nuclei provide evidence of a considerable increase in nuclear stability with increasing neutron number in the nucleus. The identification and decay properties of the  $Z=112$ ,  $114$  isotopes obtained at Dubna (1, 5) have been recently confirmed in several independent experiments [6-9]. Here we describe the synthesis of  $^{293,294}117$  ( $N=176,177$ ) isotopes in the  $^{48}\text{Ca} + ^{249}\text{Bk}$   $4n$  and  $3n$  reactions (see Fig. 2). If  $(p,xn)$  reactions had occurred, then known isotopes of  $116$  would be seen and they were not, so  $117$  is correct. The observed  $\alpha$ -decay chains show increasing stability (longer half-lives) for eleven new isotopes that end in the spontaneous fission (SF) of  $^{281}\text{Rg}$  ( $T_{\text{SF}} = 26\text{s}$ ) and  $^{270}\text{Db}$  ( $T_{\text{SF}} = 1\text{d}$ ) [10]. The decay properties of the observed eleven neutron-rich isotopes demonstrate the decisive role of the shell effects in the stability of the heaviest nuclei approaching  $N = 184$ .

## 2. Experimental Procedures and Results

The  $^{249}\text{Bk}$  material ( $T_{1/2}$  of 320 d) was produced at the High Flux Isotope Reactor (HFIR) at Oak Ridge National Laboratory (ORNL). Targets of Cm and Am were irradiated with a flux of  $2.5 \times 10^{15} \text{ n/cm}^2\text{s}$  as part of an ongoing  $^{252}\text{Cf}$  campaign. Following a three-month cooling period, extensive chemical separations were carried out over the next three months. The Bk chemical fraction, separated and purified at the Radiochemical Engineering Development Center (REDC) at ORNL, shown in Fig. 3, contained 22.2 mg of  $^{249}\text{Bk}$ , only 1.7 ng of  $^{252}\text{Cf}$ , and no other detectable impurities.



Fig. 3 The 22 mg of  $^{249}\text{Bk}$  at the end of the separation at REDC (Oak Ridge).

Six arc-shaped targets, each with an area of  $6.0 \text{ cm}^2$ , were made at the Research Institute of Atomic Reactors (Dimitrovgrad, RF) by depositing Bk nitrate onto  $0.74\text{-mg/cm}^2$  Ti foils to a thickness of  $0.31 \text{ mg/cm}^2$  of  $^{249}\text{Bk}$ . The targets were mounted on the perimeter of a disk that was rotated at 1700 rpm perpendicular to the beam direction while the beam was wiggled vertically up and down over the target. The experiments were performed employing the Dubna Gas-Filled Recoil Separator [1, 12] and the heavy-ion cyclotron U-400 at JINR.

More details of the experiment than in our initial publication [10] will be given in a forthcoming paper [13]. The basic features are presented here. Evaporation residues (ER) passing through the separator with an overall transmission about 35% were registered by a time-of-flight system with a detection efficiency of 99.9%, and were implanted in a  $4 \text{ cm} \times 12 \text{ cm}$  Si-detector array with 12 vertical position-sensitive strips surrounded by eight  $4 \text{ cm} \times 4 \text{ cm}$  side detectors. The position-averaged detection efficiency was 87% of  $4\pi$  for  $\alpha$ -particles emitted from implanted nuclei. If an  $\alpha$ -particle was detected only by a side detector (its position was lost), the total energy was estimated as a sum of energy measured by the side detector and half of the threshold energy ( $\bullet 0.5 \text{ MeV}$ ), with its total

energy uncertainty increased to  $\pm 0.4$  MeV. The position resolution (FWHM) of the strip detector was  $\sim 1.2$  mm when registering correlated decay chains of the ER- $\alpha_1$ - $\alpha_2$ - $\alpha_3$ -SF type. The background rate in the detector was reduced by switching off the beam for at least 3.5 minutes after a recoil signal was detected with an implantation energy expected for  $Z=117$  ERs, followed by an  $\alpha$ -like signal with an energy between 10.7 MeV and 11.4 MeV, in the same strip, within a 2.2-mm-wide position window. For 252-MeV  $^{48}\text{Ca}$  projectiles, the excitation energy of the compound nucleus  $^{297}117$  is estimated to be  $E^* = 39$  MeV, near the expected maximum for the total ER cross section (sum of 3n and 4n evaporation channels [1]). A 70 day irradiation at this beam energy was performed from July 27 to October 29 to give a total beam dose of  $2.4 \times 10^{19}$ .

Five position-correlated decay chains were observed in the 252-MeV  $^{48}\text{Ca}$  irradiation; in each case, two or three  $\alpha$ -decays were observed between the time of arrival of the ER and the detection of SF. The averaged decay properties of the five events assigned to the  $^{293}117$  isotope are shown in Fig. 2. The average  $\alpha$  energy of the first five  $\alpha$ -particles emitted following the recoil implantations is  $E_{\alpha 1} = 11.03 \pm 0.08$  MeV and  $T_{\alpha 1} = 14(+11, -4)$  ms. The average  $\alpha$  energy emitted by the daughter nuclei and detected in three out of five chains is  $E_{\alpha 2} = 10.31 \pm 0.09$  MeV with  $T_{\alpha 2} = 0.22(+0.26, -0.08)$  s. The third  $\alpha$ -transition was observed to have  $E_{\alpha 3} = 9.74 \pm 0.08$  MeV and  $E_{\alpha 3} = 9.48 \pm 0.11$  MeV, and  $T_{\alpha 3} = 5.5(+5.0, -1.8)$  s. All five decay chains ended in spontaneous fission with  $T_{\text{SF}} = 26(+25, -8)$  s.

The maximum cross section is expected for the 4n evaporation channel for the  $E^* = 39$  MeV excitation energy. The observed decay chains are assigned to originate from the isotope  $^{293}117$ . This assignment is thus supported by the systematics of the cross sections  $\sigma_{\text{xn}}(E^*)$  measured previously for the production of superheavy isotopes with  $Z=108, 112-116$ , and 118 in  $^{48}\text{Ca}$ -induced reactions [1], by calculations made for the evaporation residues of the reaction  $^{249}\text{Bk} + ^{48}\text{Ca}$  [14-16], and by the result of our  $^{249}\text{Bk} + ^{48}\text{Ca}$  experiment performed at lower beam energy (discussed next). In the  $E_{\alpha}$  energy range between 8.8 MeV and 11.3 MeV, where we expect  $\alpha$ -particles of the first five transitions  $117 \rightarrow 115 \rightarrow 113 \rightarrow 111 \rightarrow 109 \rightarrow 107$ , the counting rate was 0.17/s (with beam on) and  $10^{-3}$ /s (beam off) for the whole area of the front detector. The position window defined by the resolution of the detector is about 0.005 of the entire detector area, so background rates in a given detector "pixel" are small. The total numbers for random sequences [17] imitating each of the observed five decay chains were calculated to be  $6 \times 10^{-6}$ ,  $10^{-3}$ ,  $10^{-5}$ ,  $3 \times 10^{-11}$  and  $3 \times 10^{-11}$ .

Next, the experiment was run at a  $^{48}\text{Ca}$  energy of 247 MeV for 70 days with a total beam dose of  $2 \times 10^{19}$ . The excitation energy of the compound nucleus  $^{297}117$  was approximately 35 MeV, which favors the 3n reaction channel. A new decay chain with six consecutive  $\alpha$ -decays and ending in SF was detected, see Fig 2. In this chain, the  $Z=111$  great-granddaughter nucleus emitted an  $\alpha$ -particle with  $E_{\alpha 4} = 9.00$  MeV instead of undergoing SF. Then at least two more  $\alpha$ -transitions followed, and after about 33 hours, a fission event was recorded. The latter observation is significantly different from the known decay properties of  $^{294}(118)$  nucleus [1] and it was made when about 75% of the  $^{249}\text{Bk}$  still remained in the target. Therefore, we assign the chain to the decay of the neighboring odd-odd nucleus  $^{294}117$ . The chance probability for this chain was  $6 \times 10^{-4}$ .

The decay properties of the neighboring isotopes  $^{293}117$  and  $^{294}117$ , their daughters  $^{289}115$  and  $^{290}115$ , and granddaughters  $^{285}113$  and  $^{286}113$  are essentially the same but change significantly for the great-granddaughter nuclei. In spite of a strong hindrance resulting in a relatively long half-life, SF is a principal decay mode of the odd-even nucleus  $^{281}111$  (see fig. 2). However, the heavier isotope  $^{282}111$  undergoes  $\alpha$ -decay. The SF decay of  $^{281}111$  can be understood by comparing the present results with the properties of the neighboring even- $Z$  nuclei. In the  $T_{\text{SF}}(N)$  systematics, the decrease in the half-life with increasing neutron number for nuclei with  $N > 162$  changes to a strong increase in stability as  $N$  approaches the spherical shell at  $N=184$  [18]. Minimum values of  $T_{\text{SF}}$  are characteristic in the transition region  $N=168-170$ . These  $T_{\text{SF}}$  have minimum values because the effect of nuclear shells is at a minimum. For example, the  $Z=110$  darmstadtium isotopes with  $N=169$  and 171 and the

$Z=112$ ,  $N=170$ ,  $172$  copernicium isotopes, undergo SF rather than  $\alpha$ -decay [1]. The odd- $Z$  isotopes of elements  $113$  and  $115$  with  $N=169$ - $173$  have a preference for  $\alpha$ -decay [19, 20] because of their high hindrance of SF for nuclei with odd number of protons and the relatively low  $T_\alpha$ . Only in isotopes of elements  $105$  is SF observed where the  $\alpha$ -decay half-life exceeds  $10^5$  s for  $^{268}\text{Db}$ . The reaction  $^{249}\text{Bk}+^{48}\text{Ca}$  yields daughter nuclei that originate from the evaporation residues  $^{293}117$  and  $^{294}117$  with one or two extra neutrons compared with those produced in the lower- $Z$  reactions. Approaching closer the  $N=184$  shell should yield a decrease in their decay energy  $Q_\alpha$  and an increase in  $T_\alpha$  with respect to the neighboring lighter isotopes at the same  $Z$ . This behavior is clearly observed experimentally for all the isotopes with  $Z\geq 111$ , for the  $^{293}117$  chain and in  $Q_\alpha$  for  $^{294}117$ . The decay times for  $Z\geq 112$  in the  $^{294}117$  chain are far longer than those in the  $^{293}117$  chain, see Fig.4. By analogy with the neighboring even- $Z$  isotopes, all the nuclei in the  $^{293}117$  and  $^{294}117$  decay chains with  $Z>111$  and  $N\geq 172$  should undergo  $\alpha$ -decay. With its odd proton, the nucleus  $^{281}111$  ( $N=170$ ) is in a "critical" region, and may avoid SF only because of the hindrance produced by its odd proton. However, in spite of a hindrance of  $3 \times 10^4$  with respect to its even-even neighbor  $^{282}112$ , the isotope  $^{281}111$  has a probability  $b_{SF} \geq 83\%$  for SF. So, even the high hindrance caused by its odd proton does not "save" this nucleus from SF because of the weakening of the stabilizing effect of the  $N=162$  and  $N=184$  neutron shells. The presence of an extra, unpaired neutron in the neighboring isotope  $^{282}111$  further hinders SF relative to the  $\alpha$ -decay of this nucleus. The experimental  $Q_\alpha$  and half-lives  $T_\alpha$  are presented in Figs. 4a,b for isotopes with  $Z=111, 113, 115$  and  $117$ . As  $N$  increases,  $Q_\alpha$  decreases with a considerable increase in  $T_\alpha$ . The isotopes of elements  $111$  and  $113$  exhibit an especially strong growth of  $T_\alpha(N)$ . Except for  $^{281}111$ , all the nuclides shown in Fig. 4, are  $\alpha$ -emitters; with  $T_\alpha$  smaller than  $T_{SF}$  to indicate the high stability of

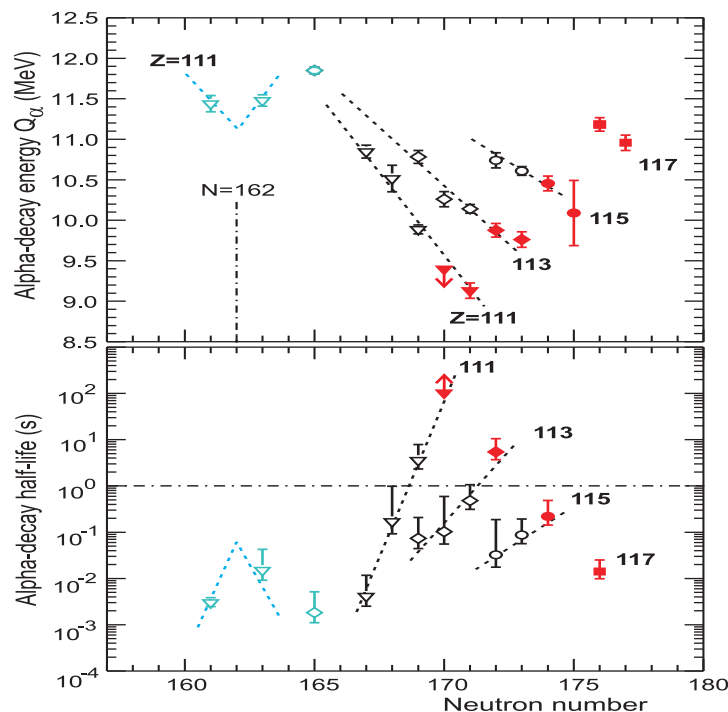


Fig. 4 Alpha-decay energies and b) half lives vs. neutron number for the isotopes of elements with  $Z=111$ - $117$ . All the nuclides with  $N>165$  were produced in  $^{48}\text{Ca}$  induced reactions. In red are shown the results obtained in the present work. In the plot (b) the values of  $T_\alpha(\text{exp})$  are given only for the nuclei produced in the decay of the isotope of  $^{293}117$  (5 events). Figures are taken from [10].

the superheavy nuclei with respect to SF. By comparing the experimental and theoretical  $\alpha$ -particle energies in Fig. 2, one sees that the macroscopic-microscopic calculations of the masses of the superheavy nuclei [11] are in a good agreement with our experiment for all the nuclei in the decay chains of the isotopes of element 117. The production cross sections for the nuclei of element 117 in the reaction  $^{249}\text{Bk}+^{48}\text{Ca}$  are  $\sigma = 0.5(+1.1, -0.4)$  pb and  $\sigma = 1.3(+1.5, -0.6)$  pb at  $E^* = 35$  MeV and  $E^* = 39$  MeV, respectively. These results are similar to previously measured cross sections for the reactions of  $^{233,238}\text{U}$ ,  $^{237}\text{Np}$ ,  $^{242,244}\text{Pu}$ ,  $^{243}\text{Am}$ ,  $^{245,248}\text{Cm}$  and  $^{249}\text{Cf}$  targets with  $^{48}\text{Ca}$  projectiles [1].

In summary, we have synthesized a new chemical element with atomic number 117 in the fusion of  $^{249}\text{Bk}$  and  $^{48}\text{Ca}$ . Two isotopes of element 117, with atomic masses 293 and 294 were observed to undergo  $\alpha$ -decay with  $E_\alpha = 11.03(8)$  MeV and  $10.81(10)$  MeV and half lives  $14(+11, -4)$  ms and  $78(+370, -36)$  ms, respectively. Their sequential  $\alpha$ -decay chains ended in spontaneous fission of  $^{281}\text{Rg}$  ( $T_{\text{SF}} \sim 26$  s) and  $^{270}\text{Db}$  ( $T_{\text{SF}} \sim 1$  ds), respectively. Our knowledge of the properties of odd- $Z$  nuclei in the region of the most neutron-rich isotopes of elements 105 to 117 is significantly expanded by our eleven newly identified isotopes which have increased stability with larger neutron number  $N$ . Investigations of the chemistry of superheavy elements and their place in the Periodic Table are opened up by their longer half-lives. The new isotopes, with superheavy nuclides previously synthesized in  $^{48}\text{Ca}$  reactions, demonstrate the critical role of nuclear shells and provide experimental verification for the existence of the predicted "Island of Stability" for superheavy elements.

We are grateful to the JINR Directorate and the U-400 cyclotron and ion source crews for their continuous support of the experiment. We acknowledged the support of the Russian Federal Agency of Atomic Energy, grants RFBR Nos. 07-02-00029, 09-02-12060, 09-03-12214, the U.S. Department of Energy through Contracts DE-AC05-00OR2272 (ORNL) and DE-AC52-07NA27344 (LLNL), grants DE-FG-05-88ER40407 (Vanderbilt University) and DE-FG07-01AL67358 (UNLV). These studies were performed in the framework of the Russian Federation/U.S. Joint Coordinating Committee for Research on Fundamental Properties of Matter.

## References

- [1] Yu. Ts. Oganessian, Jour. Phys. G **34**, R165 (2007) and earlier references therein.
- [2] S. Ćwiok, P.-H. Heenen and W. Nazarewicz, Nature **433**, 705 (2005).
- [3] S. Hofmann and G. Munzenberg, Rev. Mod. Phys. **72**, 733 (2000).
- [4] K. Morita *et al.*, J. Phys. Soc. Jpn. **76**, 045001 (2007).
- [5] Yu. Ts. Oganessian, this conference proceedings.
- [6] R. Eichler *et al.*, Nature, **447**, 72 (2007).
- [7] S. Hofmann *et al.*, Eur. Phys. J. A **32**, 251 (2007).
- [8] L. Stavsetra *et al.*, Phys. Rev. Lett., **103**, 132502 (2009).
- [9] Ch. E. Düllmann *et al.*, Phys. Rev. Lett., **104**, 252701 (2010).
- [10] Yu. Ts. Oganessian *et al.*, Phys. Rev. Lett., **104**, 142502 (2010).
- [11] A. Sobczewski, Acta. Phys. Pol. B41, **157** (2010).
- [12] Yu. Ts. Oganessian *et al.*, in Proc. of the Fourth International Conference on Dynamical Aspects of Nuclear Fission, World Scientific, Singapore, 2000, p.334.-
- [13] Yu. Oganessian *et al.*, Phys. Rev. in press (2011).
- [14] C. Shen *et al.*, Int. J. Modern Phys. E **17**, 66 (2008).
- [15] V. Zagrebaev and W. Greiner, Phys. Rev C **78**, 034610 (2008).
- [16] -Z.H. Liu and Jing-Dong Bao, Phys. Rev. C **80**, 034601 (2009).
- [17] K. H. Schmidt *et al.*, Z Phys. A **316**, 19 (1984).
- [18] R. Smolanczuk, J. Skalski, and A. Sobczewski, Phys. Rev. C **52**, 1871 (1995).
- [19] Yu. Ts. Oganessian, *et al.*, Phys. Rev. C **72**, 034611 (2005).
- [20] Yu. Ts. Oganessian *et al.*, Phys. Rev. C **76** 011601(R) (2007).