

Oscillations of B_s^0 mesons as a probe of decays with unreconstructed particles

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ABSTRACT: A new experimental technique to study the decays of B_s^0 mesons into final states involving unreconstructable particles is proposed. The method uses a characteristic feature of fast flavour oscillations of B_s^0 mesons. Analysis of the oscillation-frequency spectrum of partially reconstructed flavour-tagged B_s^0 decays provides information about the mass spectrum of unreconstructed decay products and suppresses any background that does not come from the decays of B_s^0 mesons.

KEYWORDS: B physics, Beyond Standard Model, Flavor physics, Hadron-Hadron scattering (experiments), Oscillation

ARXIV EPRINT: [1911.12729](https://arxiv.org/abs/1911.12729)

Contents

1	Introduction	1
2	Formalism	3
3	Simulation studies	6
4	Practical considerations	9
5	Conclusion	14

1 Introduction

Many interesting measurements in heavy flavour physics, both within the Standard Model (SM) framework and beyond the SM, involve decays of heavy-flavoured hadrons into final states with neutral particles that are difficult or impossible to reconstruct with collider detectors. The SM decays may involve decays with neutrinos, neutrons, or K_L^0 mesons. In addition, non-SM phenomena may generate final states involving heavy neutrinos, axion-like particles or other dark matter candidates.

Various experimental techniques have been developed to deal with such decays. In experiments at lepton colliders, one can use kinematic constraints coming from the well-defined initial state. For instance, at B -factories, where pairs of B mesons are produced at threshold without additional particles, full reconstruction of one of the B mesons provides sufficient kinematic constraints to reconstruct undetected particles in the decay of the other B meson [1–3]. At the LHCb experiment [4], which uses hadron collisions, techniques employing the information about the topology of the decays have been adopted for decays with neutrinos in the final state [5, 6]. These techniques take advantage of the large boost of initial-state hadrons in proton interactions at the LHC as well as the excellent vertex resolution of the LHCb detector to constrain the missing kinematic information by using displaced secondary vertices from weak decays.

In this paper, we propose a novel technique that uses the known pattern of B_s^0 oscillations as a tool to study a specific family of B_s^0 decays with invisible (unreconstructed) particles in the final state. The method uses the characteristically high frequency of B_s^0 flavour oscillations, which leads to multiple changes of flavour on the scale of one B_s^0 lifetime. We show that, when combined with the information about the topology of the B_s^0 decay, this feature provides a strong kinematic constraint on decays with invisible particles and efficiently suppresses any non- B_s^0 backgrounds. Contrary to other methods involving topological reconstruction at LHCb, where assumptions are needed on the nature of the missing particle (e.g., that its mass is equal to zero), this technique provides the information about the mass (or even the spectrum of invariant masses) of the invisible state(s).

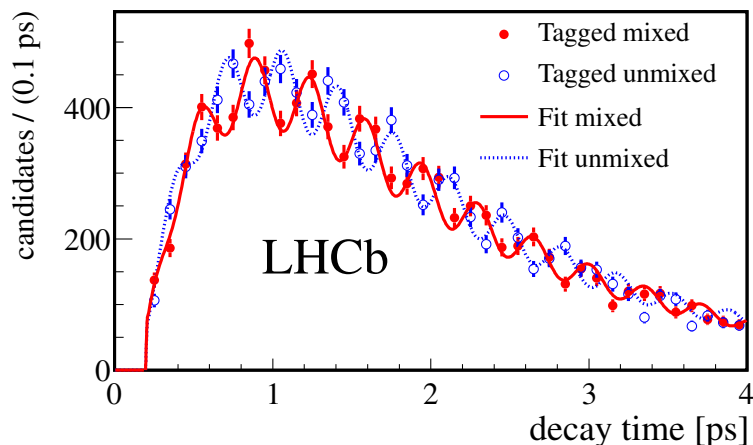


Figure 1. Decay time distribution of flavour-tagged B_s^0 mesons reconstructed with the LHCb detector in the $B_s^0 \rightarrow D_s^- \pi^+$ decays [8].

The angular frequency of B_s^0 meson oscillations Δm_s has been precisely measured to be $17.757 \pm 0.021 \text{ ps}^{-1}$, corresponding to the mixing parameter $x_s = \Delta m_s / \Gamma_s = 26.81 \pm 0.08$ [7] (where Γ_s is the average B_s^0 decay width).¹ The most precise measurement of Δm_s comes from the analysis of $B_s^0 \rightarrow D_s^- \pi^+$ decays² by the LHCb collaboration [8]. The distribution of decay time for mixed (where the flavour of the B_s^0 meson at production and decay are different) and unmixed B_s^0 meson decays in this analysis is shown in figure 1. Observing B_s^0 oscillations requires knowledge of the flavour of the B_s^0 meson at the time of production (flavour tagging), as well as precise measurement of the decay time. LHCb takes advantage of the large boost of B_s^0 mesons in the laboratory frame to transform the measurement of the flight distance of the B_s^0 meson to the measurement of its decay time. This transformation requires knowledge of the B_s^0 momentum.

In our proposed approach, we use the fact that the momentum of the partially reconstructed B_s^0 meson can be determined from the position of its decay vertex, assuming that the mass of the invisible state is known. Similar topological reconstruction has been used in studies of semileptonic decays at LHCb [9]. With the correctly determined B_s^0 momentum, the decay time distribution will exhibit oscillations with a known frequency. One can thus use the amplitude of oscillations at that frequency to constrain the mass (or measure the spectrum of masses) of the invisible particles in partially reconstructed B_s^0 decays.

The method is based on the reconstruction of momenta from the decay topology, thus its applicability is critically dependent on the vertex resolution of the detector. Simulation studies show that the technique is feasible with the value of vertex resolution corresponding to one reached at the LHCb experiment.

Since the method relies on the observation of B_s^0 oscillations, the family of decays that could be studied must satisfy two conditions: the decay should be flavour-specific (i.e., the decays of B_s^0 and \bar{B}_s^0 mesons to the final state under consideration should at least have different probabilities, or one of them should be suppressed), and the decay vertex should be

¹Natural units with $\hbar = c = 1$ are used.

²Charge conjugation is implied throughout this paper.

well-reconstructable. Possible applications of this technique could be studies of decays with heavy invisible (or poorly reconstructed) particles in the final state, both for SM decays, e.g., involving neutrons, and for searches for non-SM particles. One particular example of the processes beyond the SM that satisfy the above conditions is the model suggested in refs. [10–12], where a heavy invisible particle is accompanied by a SM baryon. As a benchmark SM decay mode, one can consider the decay $B_s^0 \rightarrow D_s^- p \bar{n}$, which satisfies both conditions and should have a significant rate since it proceeds via the favoured $b \rightarrow c$ transition. Reconstruction of decays with SM neutrinos, such as $B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu$, can also be considered; in that case, this technique will provide the additional constraint for the reconstruction of kinematic variables related to neutrinos, such as the invariant mass of the $\mu^+ \nu_\mu$ system.

2 Formalism

Consider the decay $B_s^0 \rightarrow AX$, where the system A (it can be a single state or a combination of several particles) is reconstructed, and X is invisible or partially reconstructed. Consider as well that the system A tags the flavour of B_s^0 meson, i.e., decays of B_s^0 and \bar{B}_s^0 mesons to the final state AX at least do not occur at the same rate or one of them is forbidden. The system A or the combination AX is also required to have a well-reconstructable vertex that coincides with the B_s^0 decay vertex, e.g., it should be reconstructed from at least two charged particles.

To be specific, let us give an example in the context of searches for dark matter (DM) candidates in B meson decays suggested in ref. [10]. In that model, transitions such as $\bar{b} \rightarrow cs\psi$ are expected, where ψ is a heavy neutral DM particle with negative baryonic charge. This transition gives rise to such B_s^0 meson decays as $B_s^0 \rightarrow \Xi_c^0 \psi$, where ψ would act as the invisible state X in the notation used in our formalism. While the ground state Ξ_c^0 is long-lived and does not constrain the B_s^0 decay position, the excited Ξ_c^{*0} state decaying to $\Lambda_c^+ K^-$ forms a well-defined vertex and can serve as the state A in our notation.

Similarly, for the SM decay $B_s^0 \rightarrow D_s^- p \bar{n}$ mentioned above, the antineutron serves as the invisible particle X , while the B_s^0 decay vertex is provided by the combination $D_s^- p$ that acts as the state A . In the reconstruction of semileptonic decays such as $B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu$, where the neutrino is assumed to be massless, one can consider the combination of $\mu^+ \nu_\mu$ as the particle X and treat D_s^- as A . Although D_s^- is long-lived, the decay vertex of the B_s^0 meson is constrained by the D_s^- and μ^+ , thus our technique is still applicable.

Under the conditions formulated above, the momentum of the B_s^0 meson can be reconstructed from the momentum of the system A (p_A) and the angle between the directions of A and B_s^0 in the laboratory frame (θ) up to a two-fold ambiguity:

$$p_B = \frac{(M_B^2 + \Delta) p_A \cos \theta \pm E_A \sqrt{\Delta^2 + M_B^2 (M_B^2 + \Delta - 4M_A^2 - 4p_A^2 \sin^2 \theta)}}{2(M_A^2 + p_A^2 \sin^2 \theta)}, \quad (2.1)$$

where M_B , M_A and M_X are masses of the B_s^0 , A and X states, respectively, $\Delta \equiv M_A^2 - M_X^2$, and $E_A \equiv \sqrt{M_A^2 + p_A^2}$. Figure 2 demonstrates the behaviour of p_B as a function of p_A and the angle θ for fixed $M_A = 2 \text{ GeV}$. The two solutions for the momentum are denoted in the following text as $p_B^{(\min)}$ and $p_B^{(\max)}$.

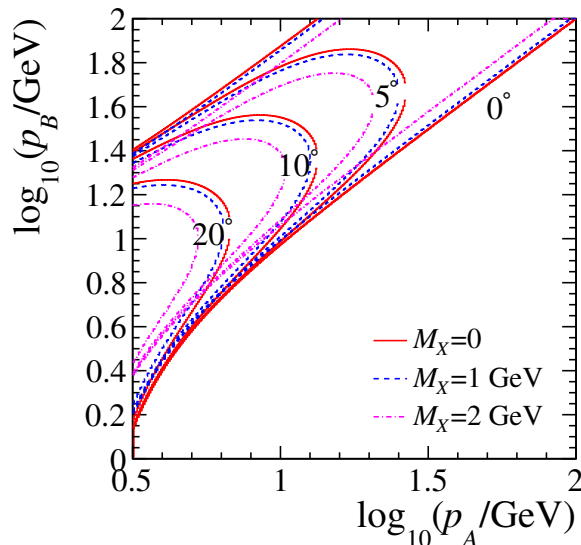


Figure 2. Logarithms of the magnitudes of the reconstructed momentum p_B of the B_s^0 meson as functions of the momentum p_A of the reconstructed state A of mass 2 GeV for the different angles θ and masses m_X of the invisible state X . The p_B solutions for each value of p_A are plotted following eq. (2.1).

The observation of B_s^0 oscillations requires the knowledge of the B_s^0 momentum p_B , as well as the decay length L . This gives the decay time

$$t = L \frac{M_B}{p_B}. \tag{2.2}$$

Oscillations result in a modulation of the rate of flavour-tagged B_s^0 decays. Assuming perfect flavour tagging, the decay time distributions for the mixed and unmixed B_s^0 decays are

$$f_{\pm}(t) \propto e^{-\Gamma_s t} \left[\cosh\left(\frac{\Delta\Gamma_s t}{2}\right) \pm \cos(\Delta m_s t) \right], \tag{2.3}$$

where the sign “+” corresponds to unmixed and “-” to mixed decays. Here, $\Delta\Gamma_s$ is the width difference between the heavy and the light B_s^0 mass eigenstates, and Δm_s is their mass difference, which drives the oscillations.

In the case of partially reconstructed B_s^0 decays, one can use the topologically reconstructed momentum p_B from eq. (2.1) in the calculation of decay time. Since there are two solutions for p_B , a sample of partially reconstructed decays will provide two decay time distributions, $f^{(\min)}(t)$ and $f^{(\max)}(t)$. Only one of the solutions corresponds to the true B_s^0 momentum for each decay, thus wrong solutions will not produce a coherent oscillation pattern and will result in a smooth background contribution in both densities.

Once the correct mass M_X is taken in the calculation of p_B , the distribution of the reconstructed decay time will contain an oscillating term with the frequency corresponding to the mass difference Δm_s . In what follows, lowercase m_X is used to distinguish the X mass assigned in the calculation of p_B from its true mass M_X . A biased value of mass $m_X \neq M_X$ will result in the oscillating term having higher or lower frequency. One can

scan the values of the assigned mass m_X and calculate the amplitude of the harmonic corresponding to the mass difference Δm_s in the spectrum of reconstructed decay times. The maximum of the amplitude will correspond to the true mass $m_X = M_X$.

Instead of the absolute rates of the B_s^0 and \bar{B}_s^0 mesons, it is convenient to study the flavour asymmetry

$$a_{B_s^0}(t) = \frac{N_{\text{unmix}}(t) - N_{\text{mix}}(t)}{N_{\text{unmix}}(t) + N_{\text{mix}}(t)}, \quad (2.4)$$

where N_{mix} and N_{unmix} are the mixed and unmixed yields as a function of decay time. The asymmetry oscillates as a function of decay time. If the correct momentum is used to calculate t from the decay length L , then the oscillations of the flavour asymmetry will occur with known frequency Δm_s .

The complex amplitude $A_{B_s^0}$ of the harmonic corresponding to the measured B_s^0 oscillation frequency,

$$A_{B_s^0} \equiv C_{B_s^0} + iS_{B_s^0} = \int a_{B_s^0}(t) e^{i\Delta m_s t} dt, \quad (2.5)$$

with the real $C_{B_s^0}$ and imaginary $S_{B_s^0}$ parts, will therefore probe the correctness of the mass hypothesis for the X state. Here $A_{B_s^0}$ depends on m_X indirectly, through the dependence of reconstructed decay time t on the momentum p_B which, in turn, depends on m_X . For the correct mass $m_X = M_X$, $C_{B_s^0}(m_X)$ should reach maximum, while $S_{B_s^0}(m_X)$ should cross zero.³

An attractive feature of this approach is that the measurement of the amplitude of high-frequency oscillations does not require detailed knowledge of the low-frequency backgrounds coming from non- B_s^0 decays or random combinations of particles produced in the primary vertex of a proton-proton interaction. As a result, this method is expected to have low systematic uncertainty due to the background description.

If several invisible states X_i contribute to the B_s^0 decay, the mass spectrum $A_{B_s^0}(m_X)$ will consist of several peaks at the masses of each X_i state. The relative magnitudes of the peaks will give relative branching ratios to the final states with X_i states. Finally, if the mass spectrum of X is continuous, the shape of the function $A_{B_s^0}(m_X)$ will be the convolution of the single-mass response $A_{B_s^0}(m_X|M_X)$ for a given M_X with the invariant mass spectrum of the X system $P(M_X)$.

Even assuming perfect tagging and vertex resolution, the mass resolution will be limited since it is determined by a finite number of observed oscillation periods. However, with a large number of B_s^0 decays it should be possible to perform a deconvolution of the $A_{B_s^0}(m_X)$ function to reconstruct the true X mass spectrum. Finally, it is possible that a maximum likelihood fit to the two distributions $f^{(\text{min})}(t)$ and $f^{(\text{max})}(t)$ will provide more information than the study of the Fourier spectrum of B_s^0 flavour asymmetry. This, however, would require a good control of slowly varying components of the decay time distributions.

While the observation of oscillations in flavour asymmetry provides a distinctive experimental signature only for the fast B_s^0 oscillations, the relation between decay time and mass spectrum given by eqs. (2.1) and (2.2) can provide information for partially reconstructed decays of other weakly decaying hadrons. Such generalisation would require a

³Assuming CP conservation in the interference of the B_s^0 mixing and $B_s^0 \rightarrow AX$ decay.

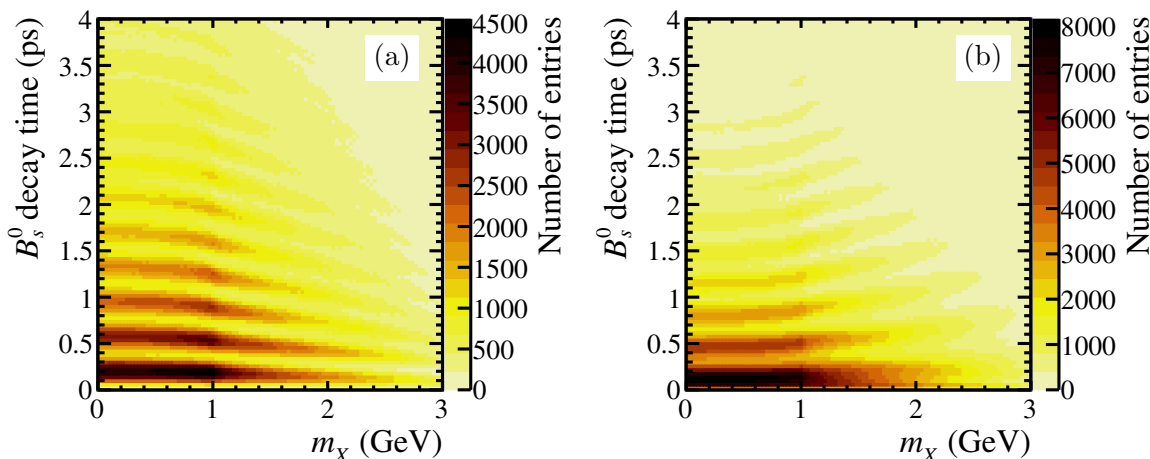


Figure 3. B_s^0 decay time reconstructed from decay length using $p_B^{(\min)}$ (a) and $p_B^{(\max)}$ (b) solutions as a function of the mass m_X used in topological reconstruction. The true mass of the X state is $M_X = 1$ GeV.

completely different study, possibly using real data in the application to a specific final state, since it will be strongly affected by various backgrounds and acceptance effects, and is therefore beyond the scope of this paper.

3 Simulation studies

Simplified Monte Carlo simulations are performed to test the feasibility of the proposed approach. The kinematic properties of two-body $B_s^0 \rightarrow AX$ decays are simulated with $M_A = 2$ GeV and M_X varying between 0 and 2 GeV. The initial B_s^0 mesons are generated with momentum and pseudorapidity distributions roughly matching those obtained in the full and simplified LHCb simulation [13] for the decay modes containing only hadrons in the final state. The transverse momentum p_T is generated following an exponential distribution with a slope of 5 GeV over the threshold of 8 GeV. The distribution of η is a Gaussian function with the mean of 3.2 and the standard deviation of 0.46. The decay time distribution follows eq. (2.3). All the following studies assume perfect flavour tagging; the effect of realistic flavour tagging performance can be taken into account by scaling the size of the data sample.

For each of the generated B_s^0 decays, the momentum (and subsequently the decay time) is calculated from topological information following eqs. (2.1) and (2.2). Figure 3 shows the distributions of B_s^0 decay time as a function of mass m_X for the decays generated with $M_X = 1$ GeV. Each plot corresponds to a solution of eq. (2.1) for the B_s^0 momentum. The dependence of the oscillation frequency on the assigned mass is clearly seen. Once the distributions of the decay time calculated this way for a certain m_X are obtained for the mixed and unmixed decays, the flavour asymmetry as a function of decay time is calculated according to eq. (2.4). Only the solution $p_B^{(\min)}$ is used since it more often gives the right answer for higher B_s^0 momenta (and thus higher decay times). An approach using

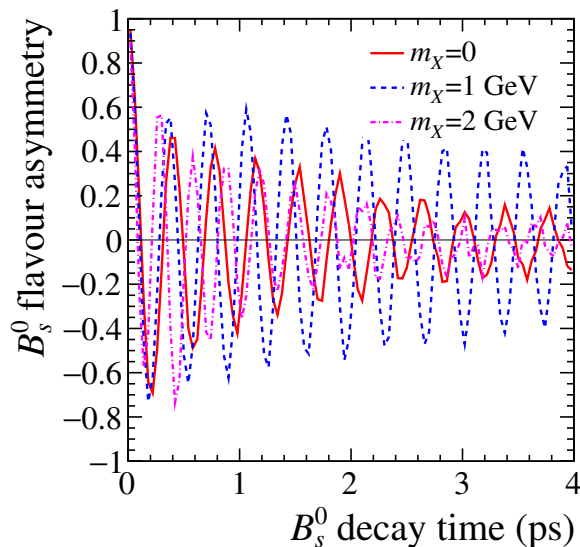


Figure 4. Dependence of B_s^0 flavour asymmetry on decay time t for the true X mass $M_X = 1$ GeV, for different assigned masses m_X .

multivariate regression based on event topology information [14] can be used to make the choice of the correct solution more efficient.

Figure 4 shows the flavour asymmetry as a function of decay time for a signal-only sample of 10^5 $B_s^0 \rightarrow AX$ events of each initial B_s^0 flavour with true $M_X = 1$ GeV and taking different assigned masses m_X of 0, 1, and 2 GeV. Perfect vertex resolution is assumed when calculating the decay time. The amplitude of flavour asymmetry oscillations is maximal for the correct mass hypothesis, while for the wrong masses, the coherence at higher decay times gets worse. In addition, the correct mass hypothesis yields a frequency of oscillations corresponding to the true mass difference Δm_s , while the wrong mass hypotheses yield higher or lower frequencies.

One then obtains the spectrum of X masses by scanning m_X in the range between 0 and 3 GeV and calculating the amplitude of the B_s^0 oscillation harmonic (eq. (2.5)) in the flavour asymmetry distribution. The selection of B_s^0 mesons in a typical analysis involves the requirement of non-zero lifetime and the reconstruction efficiency at low decay times is small (see figure 1). Thus, only the decays with $t > 0.8$ ps, corresponding to around half of the B_s^0 lifetime, are used to calculate the amplitude of the B_s^0 oscillation harmonic. The X mass spectra (in the form of real and imaginary amplitudes as a function of m_X) for true masses M_X of 0, 1, and 2 GeV are shown in figure 5. The mass resolution is better for higher-mass states; this is expected since M_X variations at higher masses lead to greater variations of momentum p_B as follows from eq. (2.1) and is illustrated by figure 2. The mass resolution is found to improve for higher M_A masses for the same reason. While the spectra presented in the plot contain uncertainties, these are not straightforward to take into account: each value of the amplitude is calculated from the same data set, so all the values of the amplitude are correlated. The uncertainties are not plotted in figure 5. The treatment of experimental uncertainties is discussed in detail in section 4.

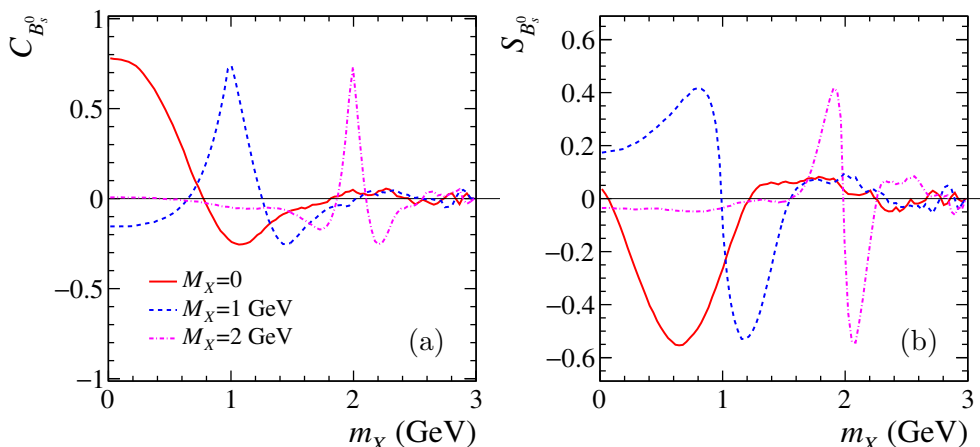


Figure 5. (a) Real $C_{B_s^0}$ and (b) imaginary $S_{B_s^0}$ parts of the amplitude of the harmonic corresponding to B_s^0 oscillation frequency Δm_s as a function of assigned mass m_X for three values of the true mass M_X , for perfect B_s^0 decay vertex resolution and 10^5 events of each B_s^0 flavour.

Good vertex resolution is an important factor for the performance of the proposed technique. Finite precision in the reconstruction of the position of the B_s^0 decay vertex transverse to its flight direction (which is approximately equal to the component σ_x of the vertex resolution perpendicular to the direction of proton beams, since B_s^0 mesons are boosted along the beams) translates to the uncertainty of the angle θ and thus of the momentum p_B . The uncertainty in p_B results in a loss of coherence at higher decay times, which in turn leads to a worse X mass resolution. Finite vertex resolution in the direction along the beams (longitudinal resolution) σ_z affects the measurements to a lesser extent: once the resolution is small compared to the oscillation period, it results in a certain reduction of the flavour asymmetry amplitude. The vertex resolution at LHCb depends on the charge multiplicity and kinematics of the decay, with typical values for a transverse component σ_x of a few tens of μm . The longitudinal component of the vertex coordinate is typically determined with the precision of 200–300 μm [15]. The effect of finite vertex resolution on the flavour asymmetry oscillations is illustrated in figures 6, separately for the longitudinal component of 300 μm and the transverse component of 40 μm .

Figures 7(a–d) demonstrate the effect of a gradual increase of the transverse resolution to 20 μm and 40 μm on the mass spectra $C_{B_s^0}$ and $S_{B_s^0}$. These plots use the same sample size of 10^5 as for figure 5. Figures 7(e,f) show the mass spectra obtained with vertex resolution $\sigma_x = 40 \mu\text{m}$ and $\sigma_z = 300 \mu\text{m}$, and using the smaller sample of 10^4 perfectly tagged B_s^0 decays of each flavour. Taking the typical effective flavour tagging efficiency at LHCb of 6% [16] and assuming the statistical uncertainties to scale as inverse square root of the effective tagging efficiency, this corresponds to around 8×10^4 reconstructed B_s^0 decays of both flavours. This is a realistic number of decays with, e.g., single-charm final states for the currently available LHCb data set recorded during the 2011–2018 data taking period. The possible benchmark mode $B_s^0 \rightarrow D_s^- p \bar{n}$ mentioned above should have sufficient yield in the currently accumulated LHCb data sample.

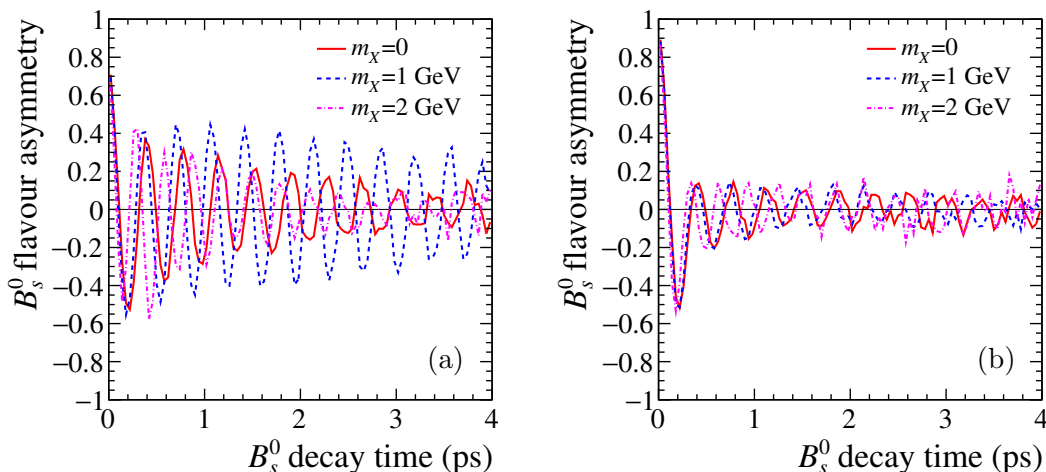


Figure 6. Dependence of B_s^0 flavour asymmetry on decay time t for the true X mass $M_X = 1$ GeV, for different assigned masses m_X , with (a) non-zero longitudinal vertex resolution $\sigma_z = 300 \mu\text{m}$ and perfect transverse resolution, and (b) non-zero transverse vertex resolution $\sigma_x = 40 \mu\text{m}$ and perfect longitudinal resolution.

4 Practical considerations

Analyses using the technique proposed above with experimental data will need to deal with such effects as background contamination, imperfect flavour tagging, finite vertex resolution and non-uniform acceptance, as well as be able to provide statistical and systematic uncertainties of the measured parameters. The exact approaches to be used will depend strongly on the decay mode under study. In this section, possible ways to address these issues are demonstrated using the same fictional two-body decay $B_s^0 \rightarrow AX$ with the fully reconstructed state A of mass 2 GeV and the unreconstructed state X of a well-defined mass.

There are several possible background contributions which may affect the correctly tagged $B_s^0 \rightarrow AX$ signal considered in section 3. One source of backgrounds is the events where the flavour tag is uncorrelated with the flavour of the fully reconstructed A state (incoherent backgrounds). These could be due either to true signal events with poor tags, or to random combinations of tracks that do not come from a single beauty decay. Since these events will not introduce a flavour asymmetry, their contribution results in a reduction of the statistical precision of the measurement: it can be accounted for as a factor that effectively reduces the signal sample size, similar to the effective tagging efficiency factor introduced to account for imperfect flavour tagging [16].

Another source of backgrounds could come from the decays of beauty hadrons other than the B_s^0 . In that case, a non-zero flavour asymmetry may emerge, however, it will either be slowly oscillating as a function of decay time (in the case of B^0 decays) or be constant (for other types of beauty hadrons). This may introduce a systematic bias to the $C_{B_s^0}$ and $S_{B_s^0}$ terms, which slowly changes as a function of m_X , and is clearly distinguishable from the peaking signal of B_s^0 oscillations.

The backgrounds from partially reconstructed B_s^0 decays could constitute a more serious problem for this method. The situation here is similar to the traditional analyses

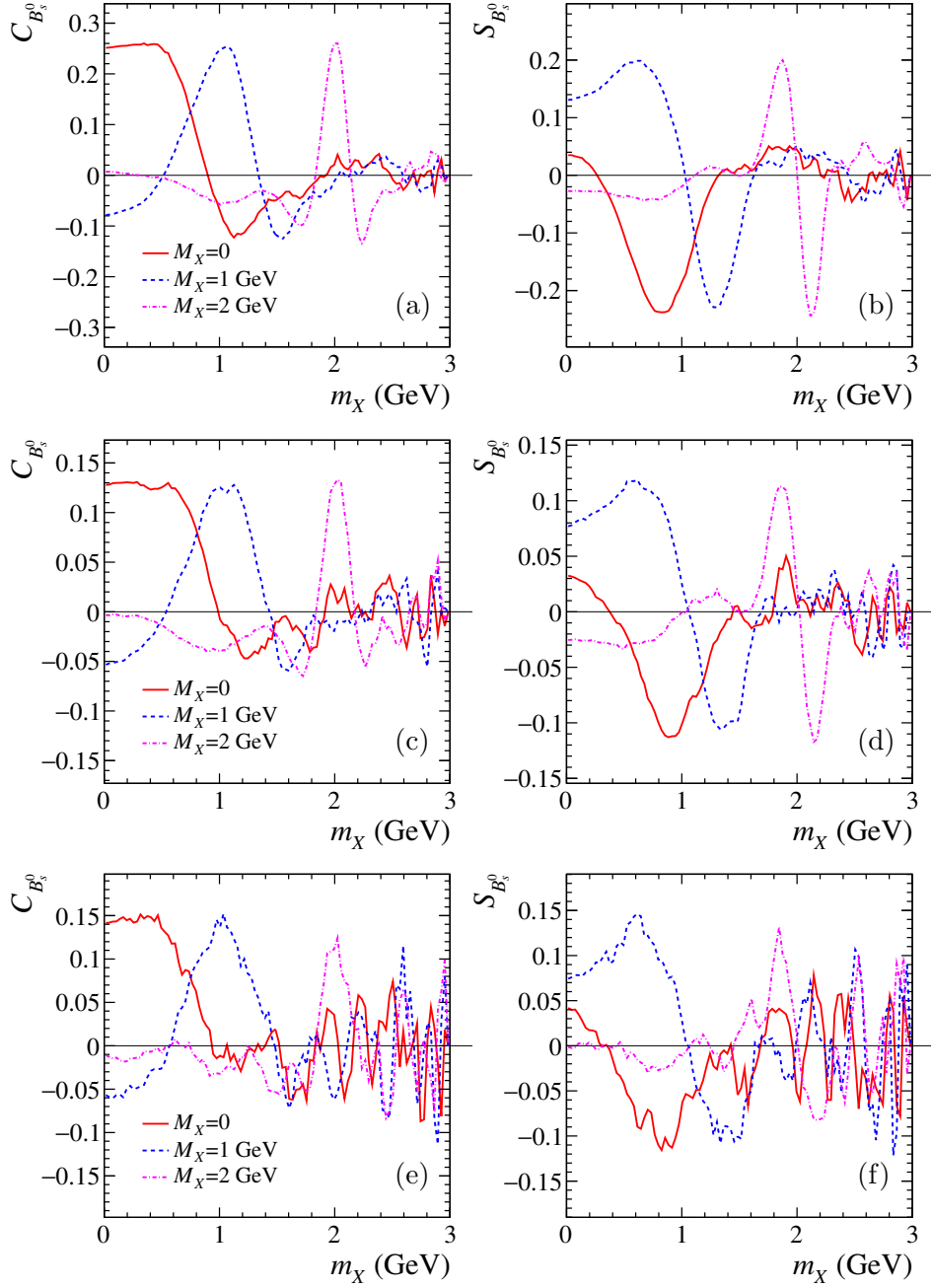


Figure 7. (a,c,e) Real $C_{B_s^0}$ and (b,d,f) imaginary $S_{B_s^0}$ parts of the amplitude of the harmonic corresponding to B_s^0 oscillation frequency Δm_s as a function of assigned mass m_X for three values of the true mass M_X , for (a,b) vertex resolution components $\sigma_x = 20 \mu\text{m}$, $\sigma_z = 0$ and 10^5 events of each B_s^0 flavour, (c,d) $\sigma_x = 40 \mu\text{m}$, $\sigma_z = 0$ and 10^5 events of each B_s^0 flavour, and (e,f) $\sigma_x = 40 \mu\text{m}$, $\sigma_z = 300 \mu\text{m}$ and 10^4 events of each B_s^0 flavour.

that use invariant masses to statistically separate signal and background components. A dedicated background study would need to be done for every specific final state. The spectra for the well-known decay modes can be obtained from simulation (or even from data if they can be fully reconstructed) and added to the function fitted to the $C_{B_s^0}(m_X)$ and $S_{B_s^0}(m_X)$ spectra. Depending on the quantum numbers (e.g., baryon or lepton number) and the masses of the unreconstructed states, certain backgrounds might be prohibitively high, but this needs to be studied on a case-by-case basis.

To demonstrate how one can evaluate the significance of the observation of the $B_s^0 \rightarrow AX$ decay and X mass resolution, and to show the effect of non- B_s^0 backgrounds, 1000 pseudoexperiments are generated with a sample size of 10^4 , 3×10^4 and 10^5 decays in each pseudoexperiment. The decays generated are $B_s^0 \rightarrow AX$ with the correct flavour tag for the signal, $B_s^0 \rightarrow AX$ with a random flavour tag for the incoherent background, and $B^0 \rightarrow AX$ to illustrate the effect of non- B_s^0 backgrounds. The mass of the X state is fixed to either $M_X = 1 \text{ GeV}$ or $M_X = 2 \text{ GeV}$. The vertex resolution is taken to be $\sigma_z = 300 \mu\text{m}$ in the longitudinal and $\sigma_x = 40 \mu\text{m}$ in the transverse direction. The production properties of the initial B_s^0 are the same as in section 3.

Figure 8 shows the distributions of the $C_{B_s^0}(m_X)$ and $S_{B_s^0}(m_X)$ values collected after adding up all pseudoexperiments with the sample size of 10^4 events. The spectra are obtained for four cases: $B_s^0 \rightarrow AX$ decay with $M_X = 1 \text{ GeV}$ (a,b) and with $M_X = 2 \text{ GeV}$ (c,d), for the non- B_s^0 background decay $B^0 \rightarrow AX$ (e,f) and for the incoherent background with the B_s^0 lifetime, where the flavour of the B_s^0 is uncorrelated with the final state flavour (g,h).

As a result, one obtains the distributions of the terms $C_{B_s^0}(m_X)$ and $S_{B_s^0}(m_X)$ at each point m_X . These distributions are well parametrised by a Gaussian function, and the corresponding standard deviation $\sigma_C(m_X)$ as a function of m_X is shown in figure 9(a). The function $\sigma_S(m_X)$ is the same within statistical precision. Moreover, the values of $\sigma_{C,S}(m_X)$ are the same within precision for all four modes considered. The evaluation of the significance of the observed oscillation signal will also need to account for correlations between the measurements of the $C_{B_s^0}$ and $S_{B_s^0}$ terms at different values of m_X . The correlation function $\rho_C(m_X, m'_X)$ obtained from the same pseudoexperiments is shown in figure 9(b); the function $\rho_S(m_X, m'_X)$ is identical within the precision of our simulation.

The evaluation of mass resolution is performed with a simple approach, with the value m_X at which the function $S_{B_s^0}(m_X)$ crosses zero taken as the estimate of the measured mass $M_X^{(\text{meas})}$. Figures 8(b,d) show that this estimate gives a value that is somewhat biased towards higher masses (the true mass is indicated with a dashed line). The distributions of the measured masses $M_X^{(\text{meas})}$ for $M_X = 1 \text{ GeV}$ and $M_X = 2 \text{ GeV}$ and 10^4 events of each B_s^0 flavour are shown in figure 10. The mass resolutions are $39.0 \pm 1.1 \text{ MeV}$ and $16.7 \pm 0.4 \text{ MeV}$, and the means of the mass distributions are $1.0659 \pm 0.0012 \text{ GeV}$ and $2.0195 \pm 0.0005 \text{ GeV}$, respectively. The mass resolution demonstrates a $1/\sqrt{N}$ scaling as a function of the sample size N (see figure 11(a)), while the mass bias is found to be independent of N .

The bias in the measurement of mass using zero-crossing $S_{B_s^0}(m_X)$ point originates from the not precisely sinusoidal behaviour of the flavour asymmetry due to reduction of coherence at higher decay times as a result of finite vertex resolution. An additional mass bias might be introduced in a real experiment by, e.g., non-uniform decay time acceptance.

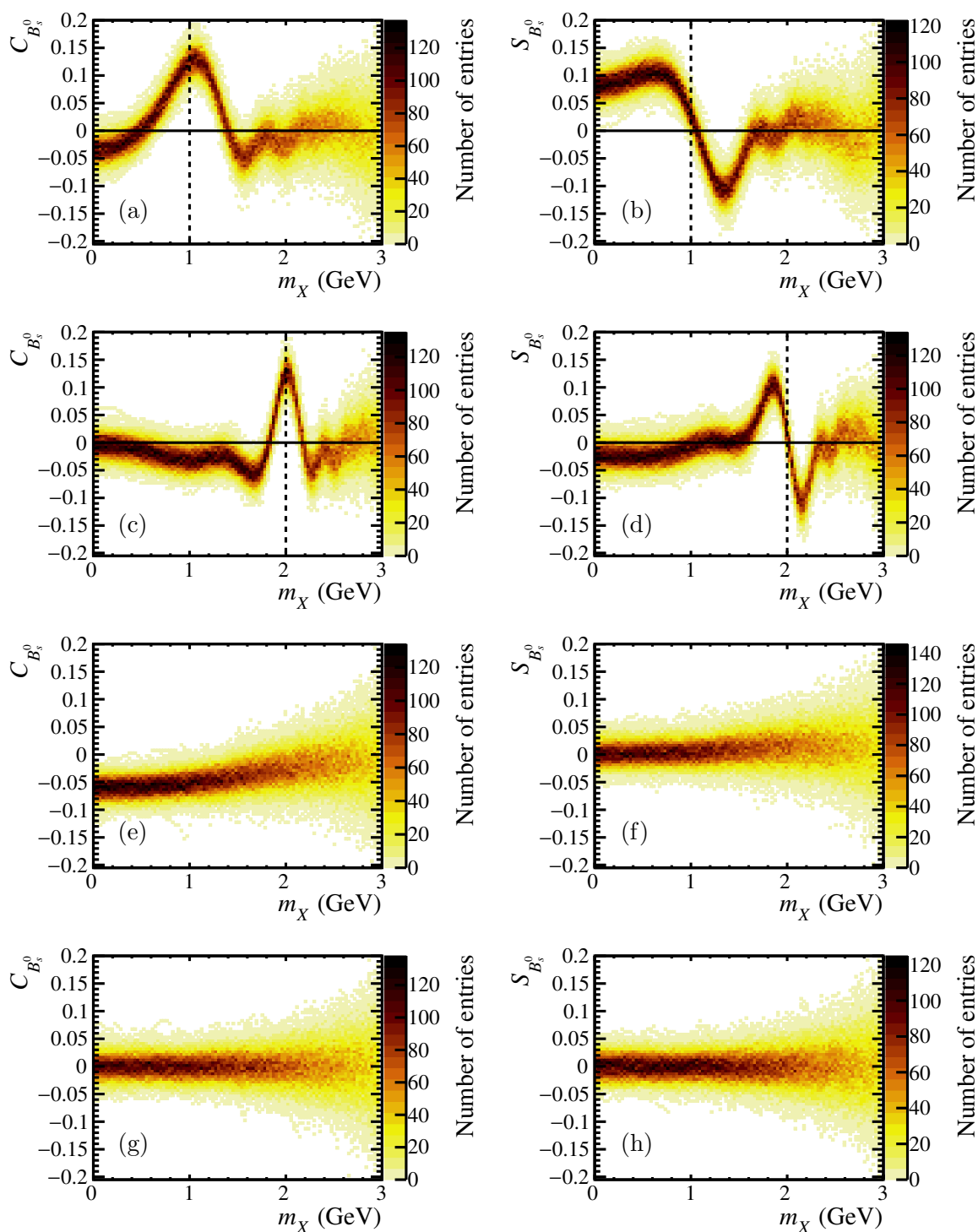


Figure 8. Distributions of the $C_{B_s^0}$ and $S_{B_s^0}$ terms based on 1000 pseudoexperiments of 10^4 events of each flavour for (a,b) $B_s^0 \rightarrow AX$ decay with $M_A = 2$ GeV and $M_X = 1$ GeV, (c,d) the same with $M_X = 2$ GeV, (e,f) $B^0 \rightarrow AX$ decay with $M_X = 1$ GeV, and (g,h) $B_s^0 \rightarrow AX$ decay with $M_X = 1$ GeV and random flavour tag.

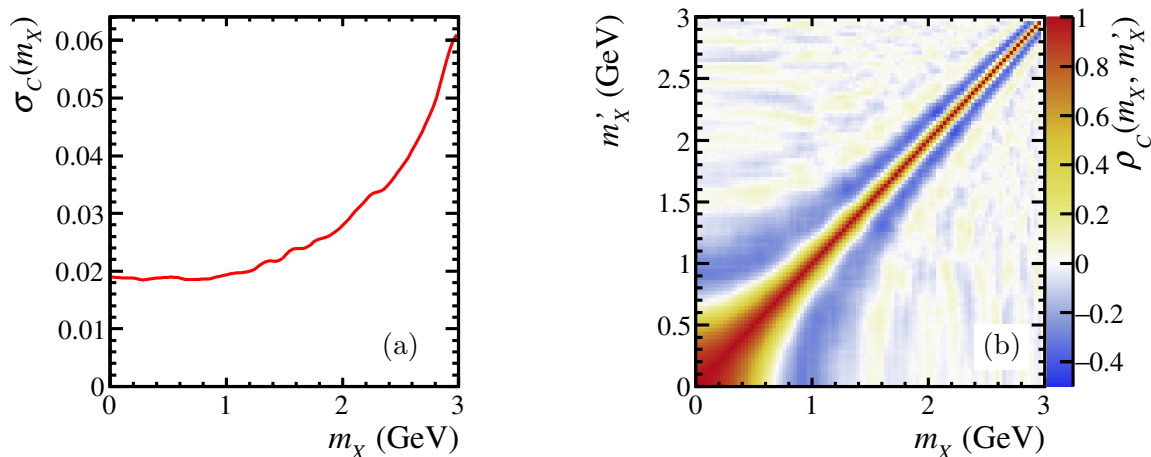


Figure 9. (a) Statistical uncertainty $\sigma_C(m_X)$ and (b) correlation function $\rho_C(m_X, m'_X)$ of the $C_{B_s^0}$ coefficient for the sample of $10^4 B_s^0 \rightarrow AX$ decays of each flavour.

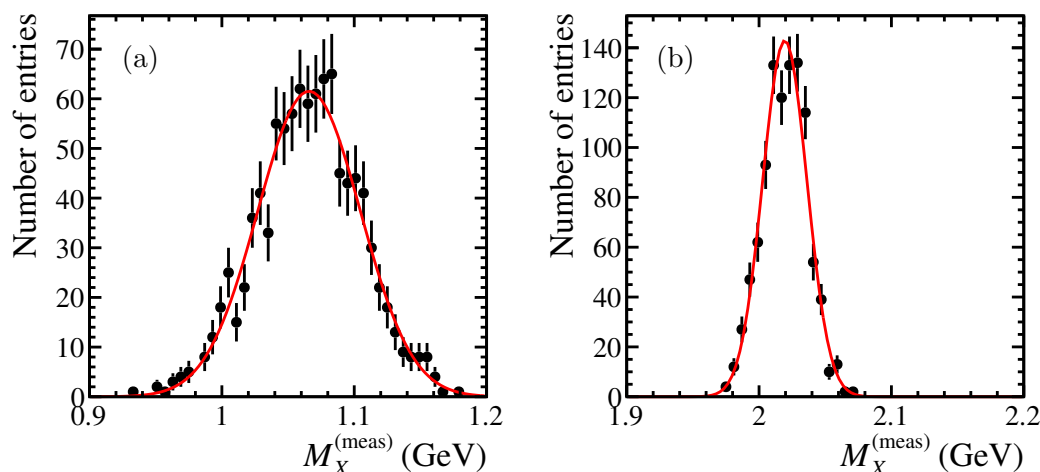


Figure 10. Distributions of the X mass estimate $M_X^{(\text{meas})}$, evaluated as the zero crossing point of the $S_{B_s^0}(m_X)$ function for true X mass (a) $M_X = 1$ GeV and (b) $M_X = 2$ GeV. Smooth line is the result of the Gaussian fit.

Such effects can be corrected using simulated data, either by introducing event-by-event weights to recover the sinusoidal behaviour of the flavour asymmetry, or by parametrising the bias and applying the correction to $M_X^{(\text{meas})}$.

To illustrate the evaluation of significance of the observed oscillation signal, a χ^2 -like test statistic accounting for correlation in the measured $C_{B_s^0}$ and $S_{B_s^0}$ values is constructed. Assuming that the contribution of peaking backgrounds is small, incoherent background with the same statistical properties (the uncertainties and the correlations of $C_{B_s^0}$, $S_{B_s^0}$ coefficients) as for the signal is taken as the null hypothesis. In that case, the test statistic that characterises the global significance of any $B_s^0 \rightarrow AX$ signal, with respect to the null hypothesis of no signal, is defined as

$$\chi_{\text{corr}}^2 = \sum_{ij} \left(C_{B_s^0}(m_i) \Sigma_{C,ij}^{-1} C_{B_s^0}(m_j) + S_{B_s^0}(m_i) \Sigma_{S,ij}^{-1} S_{B_s^0}(m_j) \right), \quad (4.1)$$

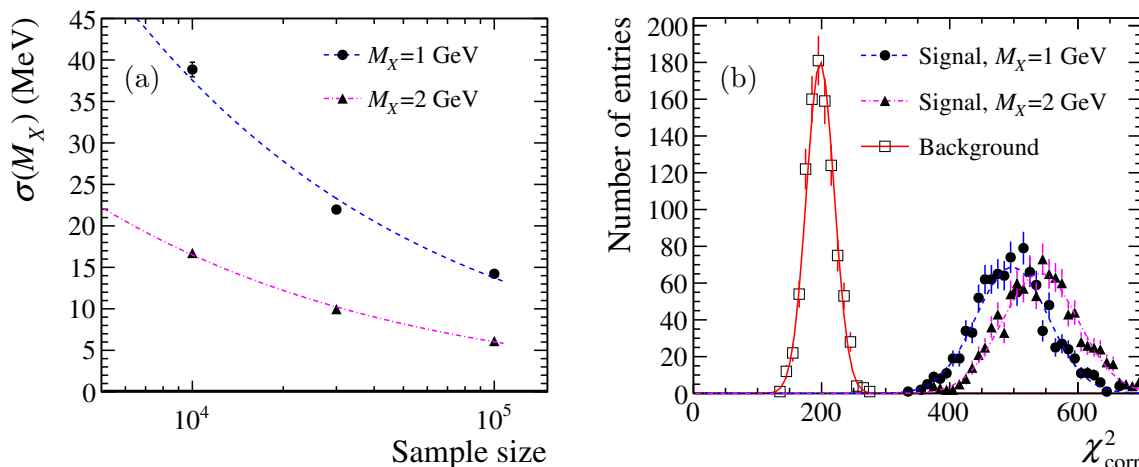


Figure 11. (a) Mass resolutions for $M_X = 1$ GeV and $M_X = 2$ GeV cases as a function of sample size. Smooth lines are the result of the fits with the function proportional to inverse square root of the sample size. (b) Distributions of the test statistic χ_{corr}^2 for background samples and signal with $M_X = 1$ GeV and $M_X = 2$ GeV. Smooth lines are results of the Gaussian fits.

where the terms $C_{B_s^0}$ and $S_{B_s^0}$ are measured at the values of mass m_i (in our case, in 100 points uniformly distributed in the range from 0 to 3 GeV), and $\Sigma_{C,S}^{-1}$ are the inverse covariance matrices for the $C_{B_s^0}$ and $S_{B_s^0}$ values at these masses:

$$\begin{aligned}\Sigma_{C,ij} &= \sqrt{\sigma_C(m_i), \sigma_C(m_j)} \rho_C(m_i, m_j), \\ \Sigma_{S,ij} &= \sqrt{\sigma_S(m_i), \sigma_S(m_j)} \rho_S(m_i, m_j).\end{aligned}\tag{4.2}$$

The distributions of the test statistic χ_{corr}^2 for the signals and for the background are shown in figure 11(b). The distributions for signal events and the background-only hypothesis are clearly separated, their difference provides a measure of the statistical significance of the effect. In this case, the significance of any deviation from the null hypothesis is around 10 standard deviations on average. In the real experiment, a similar test statistic can be used to assess the goodness of fit for a particular signal model, in which case one would need to calculate instead

$$\chi_{\text{corr}}^2 = \sum_{ij} \left(\Delta C_{B_s^0}(m_i) \Sigma_{C,ij}^{-1} \Delta C_{B_s^0}(m_j) + \Delta S_{B_s^0}(m_i) \Sigma_{S,ij}^{-1} \Delta S_{B_s^0}(m_j) \right),\tag{4.3}$$

where $\Delta C_{B_s^0}$, $\Delta S_{B_s^0}$ are the differences between the observed $C_{B_s^0}$ and $S_{B_s^0}$ functions and those expected from the model.

5 Conclusion

A novel technique to study B_s^0 meson decays with invisible particles at a hadron collider has been presented. It uses the characteristic pattern of fast B_s^0 flavour oscillations to obtain information about the B_s^0 meson momentum and suppress any non- B_s^0 backgrounds.

The technique has been studied with pseudoexperiments for a particular case of a decay $B_s^0 \rightarrow AX$, where the system A is reconstructed and constrains the B_s^0 vertex, while X is invisible. In this case, the proposed procedure provides sensitivity to the spectrum of X masses which would be impossible with other approaches.

The proposed technique can find its applications in the studies of both SM decays, e.g., involving neutrons, and for searches beyond the SM, such as decays with heavy neutrinos, axion-like particles and other dark matter candidates (see, e.g., refs. [10–12]). The benchmark measurements that should be feasible already with the currently available LHCb data set could employ Cabibbo-favoured open-charm decays such as $B_s^0 \rightarrow D_s^- p \bar{n}$. A similar approach is probably feasible for the reconstruction of decays with neutrinos (e.g., for the semileptonic decays involving muons or τ -leptons) where B_s^0 oscillations can provide additional kinematic constraints compared to conventional techniques used at LHCb.

Acknowledgments

The authors would like to thank Matthew Charles, Tim Gershon, Alexey Petrov, the CPPM LHCb group (Olivier Leroy, Julien Cogan) and members of the LHCb physics working group “semileptonic decays” (in particular, Greg Ciezarek, Lucia Grillo and Michel De Cian) for stimulating discussions and useful feedback on the paper. This work was supported by the A*MIDEX international project TAUFU, Aix-Marseille University.

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