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Scope Out Multiband Gravitational-Wave Observations of GW190521-Like Binary Black Holes with Space Gravitational Wave Antenna B-DECIGO

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Abstract: The gravitational wave event, GW190521, is the most massive binary black hole merger observed by ground-based gravitational wave observatories LIGO/Virgo to date. While the observed gravitational wave signal is mainly in the merger and ringdown phases, the inspiral gravitational wave signal of the GW190521-like binary will be more visible to space-based detectors in the low-frequency band. In addition, the ringdown gravitational wave signal will be louder in the next generation (3G) of ground-based detectors in the high-frequency band, displaying the great potential of multiband gravitational wave observations. In this paper, we explore the scientific potential of multiband observations of GW190521-like binaries with a milli-Hz gravitational wave observatory: LISA; a deci-Hz observatory: B-DECIGO; and (next generation of) hecto-Hz observatories: aLIGO and ET. In the case of quasicircular evolution, the triple-band observations of LISA, B-DECIGO, and ET will provide parameter estimation errors of the masses and spin amplitudes of component black holes at the level of order of 1–10%. This would allow consistency tests of general relativity in the strong field at an unparalleled precision, particularly with the “B-DECIGO + ET” observation. In the case of eccentric evolution, the multiband signal-to-noise ratio found in “B-DECIGO + ET” observation would be larger than 100 for a five-year observation prior to coalescence, even with high final eccentricities.

Keywords: gravitational waves; binary black holes; quasinormal modes; general relativity

1. Introduction

Among gravitational wave (GW) events detected by LIGO and Virgo during O1, O2, and O3a runs [1,2], a binary black hole (BBH) merger, GW190521 [3,4], is one of the most striking discoveries. GW190521 is the heaviest BBH merger ever observed, producing a remnant black hole (BH) with a mass of $142^{+28}_{-16} M_{\odot}$ that can be interpreted as an intermediate mass BH; the source parameters of GW190521 (and our notations) are summarized in Table 1. This measurement triggers the intense investigation of GW190521’s unique source property. In this paper, we set $G = 1 = c$ with the useful conversion factor $1M_{\odot} = 1.477 \text{ km} = 4.926 \times 10^{-6} \text{ s}$. We also assume a “Planck” flat cosmology (when it is needed) with the Hubble constant $H_0 = 67.7 \text{ km s}^{-1} \text{ Mpc}^{-1}$, and density parameters $\Omega_M = 0.307$ and $\Omega_{\Lambda} = 0.694$ [5].

Table 1. Summary of the source parameters of GW190521 as a quasicircular BBH merger, reported by the LIGO–Virgo collaboration [3] based on the BBH waveform model [6]. The symmetric 90% credible interval for each parameter is also quoted. Note that the parameters are written in our notation and masses are given in the source’s rest frame; multiply by $(1 + z)$ to convert to the observer frame.

Parameter	Symbol	
Primary mass [M_{\odot}]	m_2^r	85^{+21}_{-14}
Secondary mass [M_{\odot}]	m_1^r	66^{+17}_{-18}
Primary spin magnitude	$ \vec{\chi}_2 $	$0.69^{+0.27}_{-0.62}$
Secondary spin magnitude	$ \vec{\chi}_1 $	$0.73^{+0.24}_{-0.64}$
Total mass [M_{\odot}]	$m_t^r (= m_1^r + m_2^r)$	150^{+29}_{-17}
Mass ratio	$q (= m_1^r / m_2^r \leq 1)$	$0.79^{+0.19}_{-0.29}$
Luminosity Distance [Gpc]	D_L	$5.3^{+2.4}_{-2.6}$
Redshift	z	$0.82^{+0.28}_{-0.34}$

A key element to better understand GW190521 is the precise measurement of the binary parameters (see, for example, Ref. [7] for a possibility of an intermediate mass ratio inspiral). GW190521 is, however, a much heavier BBH system than previously observed GW events, and one of the difficulties here is the short duration and bandwidth of the GW signal that can be observed in the LIGO/Virgo band. In the case of the quasicircular BBH scenario (which is most favored by the LIGO/Virgo analysis [3,4]), the coalescing time and the number of GW cycles at frequency f (in the observer frame) are estimated as:

$$t_c \sim 1.3 (1 + z)^{-5/3} \left(\frac{m_1^r}{66M_{\odot}} \right)^{-1} \left(\frac{m_2^r}{86M_{\odot}} \right)^{-1} \left(\frac{m_t^r}{150M_{\odot}} \right)^{1/3} \left(\frac{f}{10.0 \text{ Hz}} \right)^{-8/3} \text{ s}, \quad (1)$$

$$N_c \sim 1.0 \times 10^1 (1 + z)^{-5/3} \left(\frac{m_1^r}{66M_{\odot}} \right)^{-1} \left(\frac{m_2^r}{85M_{\odot}} \right)^{-1} \left(\frac{m_t^r}{150M_{\odot}} \right)^{1/3} \left(\frac{f}{10.0 \text{ Hz}} \right)^{-5/3}. \quad (2)$$

This indicates the lack of the GW signals from the sufficiently long inspiral phase. Because of the short duration signal dominated by the merger and ringdown phases, for example, only weak constraints are obtained for the component BH spins and their orientations [4]. Furthermore, even alternative interpretations of the observed GW signal other than massive quasicircular BBH merger in general relativity (GR) would become more relevant; several plausible scenarios are assessed in Section 6 of Ref. [4] by the LIGO–Virgo collaboration.

At the same time, the estimation in Equations (1) and (2) suggests a natural way to overcome the hurdle here: observe the inspiral GW signal in the low-frequency band offered by space-based GW detectors. Future GW astronomy in the 2030s will utilize the LISA observatory in the milli-Hz band [8] and deci-Hz GW detectors such as B-DECIGO [9], a prototype GW antenna of the DECIGO mission [10,11] (other proposed GW missions in the low-frequency band, including Taiji [12] and TianQin [13] in the milli-Hz band and MAGIS [14] and TianGO [15] in the deci-Hz band, are concisely summarized in, for example, reviews by Ni [16,17]). Figure 1 plots the track of the strain sensitivity curve of a GW190521-like non-spinning BBH system, assuming the quasicircular evolution and a simple inspiral–merger–ringdown (IMR) amplitude model given in Equation (16). At the GW frequency $f = 0.1$ Hz, for example, Equations (1) and (2) give the coalescence time $t_c \sim 1.1 \times 10^5$ s and the number of GW cycles $N_c \sim 8.1 \times 10^3$ with the dimensionless characteristic strain:

$$h_c \sim 7.6 \times 10^{-22} (1 + z)^{5/6} \left(\frac{\mathcal{M}^r}{65.1M_{\odot}} \right)^{5/6} \left(\frac{f}{0.1 \text{ Hz}} \right)^{-1/6} \left(\frac{D_L}{5.3 \text{ Gpc}} \right)^{-1}, \quad (3)$$

where we introduce the source’s rest-frame chirp mass $\mathcal{M}^r \equiv \nu^{3/5} m_t^r$ with the symmetric mass ratio $\nu \equiv m_1^r m_2^r / m_t^{r2}$. In fact, we find in Section 3 that the sky and polarization

averaged signal-to-noise ratio (SNR) (whose meaning is momentarily clarified in Section 2) accumulated 5 years before the final coalescence would be $\sim 5.9 \times 10^1$ in the B-DECIGO band and ~ 2.7 in the LISA band. These estimations show that the early inspiral signal of GW190521-like BBHs would be sensitive in the LISA band and loud in the deci-Hz band.

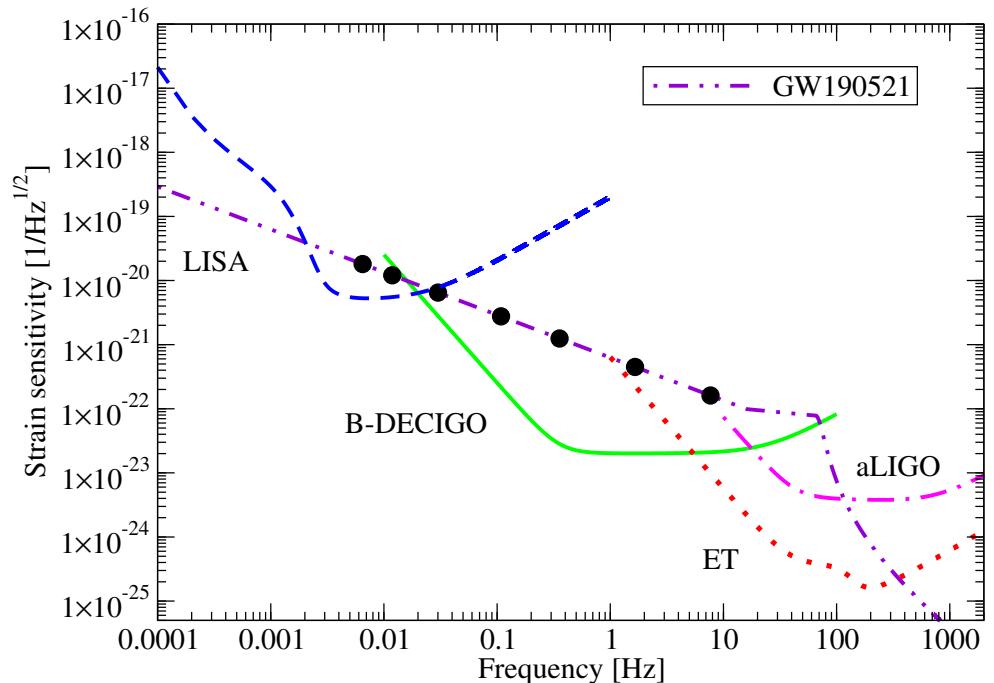


Figure 1. Strain sensitivity curves for the ground-based aLIGO and a next-generation (3G) detector: Einstein Telescope (ET) and space-based B-DECIGO and LISA, together with the GW amplitude of a GW190521-like nonprecessing, quasicircular BBH. We use the median values in Table 1 as source parameters. The noise power spectral density (PSD) for each GW detector is given in Section 2.1, and the spectral density of the BBH amplitude is obtained by the nonspinning, IMR amplitude model in Section 2.3. As a reference, we mark with the black dots with 5 years, 1 year, 1 month, 1 day, 1 h, 1 min, and 1 s before the merger time (1) from the left to the right.

Although the observation with LISA and B-DECIGO alone would provide valuable information on the inspiral GW signal from the GW190521-like BBH system, the true potential of having the low-frequency sensitivity will be revealed only when it is combined with the high-frequency sensitivity in the hecto-Hz band. As illustrated in Figure 1, the late-inspiral and merger-ringdown GW signals of the GW190521-like BBH are best detected with aLIGO, Virgo, and KAGRA [18]. In addition to ground-based GW observatories that are online, the next generation (3G) of ground-based detectors such as the Einstein Telescope (ET) [19,20] (see also Ref. [21] for Cosmic Explorer (CE)) will significantly improve the visibility of GW190521-like BBH systems. We will see in Section 3 that the averaged SNR of the late-inspiral and ringdown signals in the ET band would be $\sim 2.7 \times 10$ and $\sim 1.5 \times 10^2$, respectively. The joint “space + ground” observation across the full GW bands would be therefore the best way to observe the GW190521-like BBH systems; this is the basic idea of *multiband GW astronomy*.

Soon after the first detection of GW150914 [22], the potential of the multiband GW astronomy of BBH systems with LISA and aLIGO was emphasized [23] (see also Ref. [24]). This study was immediately followed up with more detailed analyses. The works include (but not limited to) the improved estimation of source parameters [25,26], tests of GR with high precision [27–32], refined event rate estimations [33], probing environment effects [34,35], and new data analysis ideas [36]; these all prove the scientific values added by the multiband observation. Among these investigations, Refs. [9,37–41] demonstrated that the multiband observation of stellar-mass BBH systems will further benefit from

having (B-)DECIGO in the deci-Hz band, which naturally bridges the gap between the LISA and aLIGO bands.

Goals and Organization of This Paper

Our purpose in this paper is to explore the prospects of the multiband observation of GW190521-like nonprecessing, quasicircular “intermediate-mass” BBHs. The possibility of the multiband observation of intermediate-mass BBHs with deci-Hz GW detectors was pointed out by Amaro-Seoane and Santamaria [42] and Yagi [43]. We consider these observations across the full GW spectrum provided by LISA in the milli-Hz band, B-DECIGO in the deci-Hz band, and aLIGO and ET in the hecto-Hz, looking at two specific aspects of multiband GW astronomy/physics: parameter estimation errors and tests of GR.

We begin in Section 2 by providing our set of basic tools for the signal analysis in the matched filtering technique: (i) the noise PSD in Section 2.1, (ii) the GW signal models in Sections 2.2 and 2.3, and (iii) the multiband Fisher matrix formalism in Section 2.4. In Section 3, we present the parameter estimation errors of GW190521-like nonprecessing, quasicircular BBH systems using multiband GW observations. They are displayed in Table 2 and 3. Based on the estimated errors, we then examine in Section 4 to what extent the future multiband observations of GW190521-like BBH systems will improve tests of GR; this includes an inspiral–ringdown consistency test (see Table 4) and a simple test to discriminate between the remnant BH based on GR and other remnant compact objects (see Figure 2). We conclude in Section 5 with complications and various effects, which are not covered in this paper due to our assumptions and simplifications.

The appendices contain some additional analysis and information. While a quasicircular BBH merger is the most favored scenario in the LIGO/Virgo analysis [3,4], other alternative scenarios such as an eccentric BBH merger may also be consistent with the observed source for GW190521 [44]. We briefly discuss the multiband visibility of a GW190521-like eccentric BBH in Appendix A. Additionally, the additional noise PSDs of ground/space-based current/future GW detectors is provided in Appendix B as a complement to our treatment in Section 2.1.

Table 2. Parameter estimation errors of mass parameters (m, ν) and spin parameters (χ_s, χ_a) for the GW190521-like quasicircular BBH inspiral, normalized to the total multiband SNR ρ_{tot} ; for example, the result of “BD” in “BD + aLIGO” is normalized to $\rho_{\text{tot}} = 5.93 \times 10^1$. We assume that the true binary parameters are given by the median values of Table 1—i.e., $m_t = 151.0 M_{\odot}$, $\nu = 0.246$, $\chi_s = 0.71$, and $\chi_a = 0.02$. Since the source location can be determined well by B-DECIGO (BD here) [9,39], we fix $1 + z = 1.82$. Note that only the inspiral SNR of aLIGO is 1.76, which is too small to give meaningful estimation errors. Additionally, any normalized estimation errors $\delta\hat{\theta} > 1.0 \times 10^6$ are discarded from this table.

GW Detector	SNR	$\delta\hat{m}_t/m_t$	$\delta\hat{\nu}/\nu$	$\delta\hat{\chi}_s/\chi_s$	$\delta\hat{\chi}_a/\chi_a$
BD + aLIGO					
BD	5.92×10^1	3.51×10^{-1}	5.84×10^{-1}	4.91	1.12×10^3
aLIGO	1.76
BD + aLIGO	5.93×10^1	3.44×10^{-1}	5.73×10^{-1}	4.73	1.08×10^3
BD + ET					
ET	2.72×10^1	3.70×10^2	6.00×10^2	5.86×10^3	...
BD + ET	6.52×10^1	2.88×10^{-1}	4.80×10^{-1}	3.07	7.01×10^2
LISA + BD + ET					
LISA	2.68	1.73×10^1	2.89×10^1	2.37×10^3	...
LISA + BD	5.93×10^1	1.83×10^{-1}	3.04×10^{-1}	4.01	9.11×10^2
LISA + BD + ET	6.52×10^1	1.45×10^{-1}	2.41×10^{-1}	2.52	5.71×10^2

Table 3. Parameter estimation errors of the central frequency f_c and the quality factor Q of the GW190521-like remnant BH using only the ringdown GW signal. The errors are normalized to the signal-to-noise ratio (SNR) of each corresponding detector.

GW Detector	SNR	$\delta \hat{f}_c/f_c$	$\delta \hat{Q}/Q$
aLIGO	1.08×10^1	3.42×10^{-1}	2.71×10^{-1}
ET	1.47×10^2	3.66×10^{-1}	2.86×10^{-1}

Table 4. Parameter estimation errors of the remnant mass and spin of GW190521-like and GW150914-like BBH systems, using only the ringdown GW signal as well as the inspiral GW signal inferred via the NR remnant fitting formulas; the best fit values are $(M_f/m, \chi_f) \sim (0.904, 0.883)$ for the GW190521-like BBH (recall Equation (27)) and $(M_f/m, \chi_f) \sim (0.891, 0.898)$ for the GW150914-like BBH. Here, BD is an abbreviation for B-DECIGO. The results in the single band are normalized to the SNR for a given GW observatory. In the multiband case, they are normalized to the total SNR of $\rho_{\text{LISA+BD+ET}}$.

GW190521-Like BBH				GW150914-Like BBH		
GW Detector	SNR	$\delta M_f/M_f$	$\delta \chi_f/\chi_f$	SNR	$\delta M_f/M_f$	$\delta \chi_f/\chi_f$
Single band: Ringdown GW signal						
aLIGO	1.08×10^1	3.09×10^{-1}	7.47×10^{-1}	1.78×10^1	3.05×10^{-1}	6.47×10^{-1}
ET	1.47×10^2	3.31×10^{-1}	7.90×10^{-1}	2.66×10^2	3.26×10^{-1}	6.60×10^{-1}
Single band: inspiral GW signal						
BD	5.92×10^1	3.31×10^{-1}	9.63×10^{-1}	2.51×10^2	3.19×10^{-1}	9.43×10^{-1}
Multiband: inspiral GW signal						
BD + ET	6.52×10^1	3.55×10^{-1}	5.91×10^{-1}	5.18×10^2	1.07×10^{-1}	4.48×10^{-1}
LISA + BD	5.93×10^1	2.46×10^{-1}	8.05×10^{-1}	2.51×10^2	6.39×10^{-1}	1.85
LISA + BD + ET	6.52×10^1	1.75×10^{-1}	4.98×10^{-1}	5.18×10^2	1.06×10^{-1}	4.33×10^{-1}

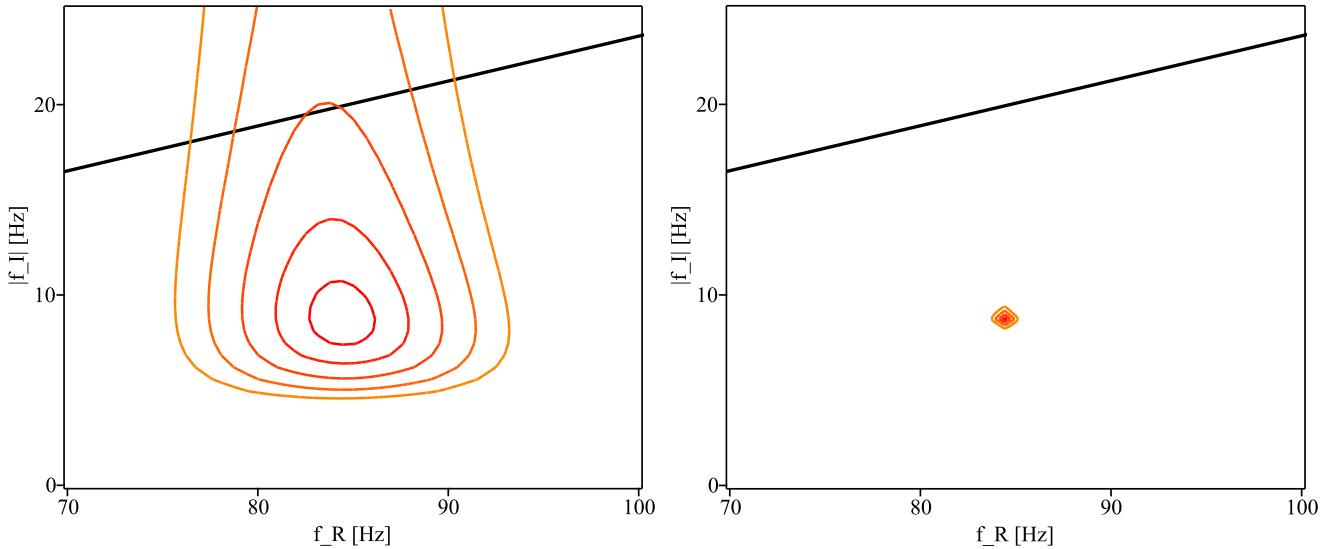


Figure 2. Simple remnant test in the (f_R, f_I) plane. The QNM frequency is given in the observer frame, and we assume the GW190521-like ringdown GW signals with $(M_f/m, \chi_f) \sim (0.904, 0.883)$ (27). The figure shows the parameter estimation error with the aLIGO noise curve (left, $\text{SNR} = 1.08 \times 10^1$) and the ET one (right, $\text{SNR} = 1.47 \times 10^2$). The black lines are the Schwarzschild limit of $|f_I|/f_R \approx 0.236$. The colored ellipses show the contours of $1\sigma, 2\sigma, 3\sigma, 4\sigma$, and 5σ . To obtain these two-dimensional plots, the time and phase parameters (t_0, ϕ_0) have been marginalized out.

Throughout this paper, the binary parameters and GW frequencies measured in the source's rest frame are denoted with the index 'r', explicitly distinguished from those in the observer frame.

2. Method for Signal-to-Noise Ratio and Fisher Analysis

In this section, we summarize our methodology of multiband GW data analysis in the specific context of a nonprecessing, spinning, quasicircular GW190521-like BBH system, where the spins are aligned to the orbital angular momentum; an alternative scenarios of non-zero orbital eccentricity in GW190521-like BBH system will be discussed in Appendix A. Our simple framework here is largely based on the previous B-DECIGO works [9,38].

2.1. aLIGO, ET, B-DECIGO and LISA

For our multiband GW observation, we follow Ref. [38] by considering the four GW observatories: aLIGO and ET in the hecto-Hz band, B-DECIGO in the deci-Hz, and LISA in the milli-Hz band (see also Appendix B for some of other GW detectors in these bands). It should be noted that we shall use the "non sky-averaged" PSD $S_n(f)$; we will account for the average over the GW detector's antenna pattern function at the level of the waveform.

For aLIGO in the hecto-Hz band, we use Equation (4.7) in Ref. [45],

$$S_n^{\text{aLIGO}} = 10^{-48} \left(0.0152 x^{-4} + 0.2935 x^{9/4} + 2.7951 x^{3/2} - 6.5080 x^{3/4} + 17.7622 \right) \text{ Hz}^{-1};$$

$$x = \frac{f}{245.4 \text{ Hz}}. \quad (4)$$

For a third generation (3G) GW interferometer—ET in the hecto-Hz band—we find Table 1 in Ref. [46] as:

$$S_n^{\text{ET}} = 1.5 \times 10^{-52} \left(y^{-4.1} + 186.0 y^{-0.69} \right. \\ \left. + 233.0 \times \frac{1.0 + 31.0 y - 65.0 y^2 + 52.0 y^3 - 42.0 y^4 + 10.0 y^5 + 12.0 y^6}{1.0 + 14.0 y - 37.0 y^2 + 19.0 y^3 + 27.0 y^4} \right) \text{ Hz}^{-1};$$

$$y = \frac{f}{200.0 \text{ Hz}}. \quad (5)$$

For B-DECIGO in the deci-Hz band, we use Equation (20) in Ref. [38], originally proposed by Nakamura et al. [9]:

$$S_n^{\text{BD}} = \left[(2.01 \times 10^{-23})^2 + \left(\frac{2.53 \times 10^{-18}}{F^2} \right)^2 + \left(8.00 \times 10^{-22} F \right)^2 \right] \text{ Hz}^{-1};$$

$$F = \frac{f}{1.0 \times 10^{-3} \text{ Hz}}. \quad (6)$$

For LISA in the milli-Hz band, we use Equation (1) of Ref. [47], which is based on the 2018 LISA Phase-0 reference design parameters. It reads:

$$S_n^{\text{LISA}} = \left[\left(2.4602 \times 10^{-41} + 4.0504 \times 10^{-38} z^2 + \frac{4.7850 \times 10^{-48}}{z^2} \right. \right. \\ \left. \left. + \frac{2.8485 \times 10^{-51}}{z^4} + \frac{3.9412 \times 10^{-58}}{z^6} \right) + S_c^{\text{LISA}} \right] \text{ Hz}^{-1};$$

$$z = \frac{f}{1.0 \text{ Hz}}. \quad (7)$$

where $S_c^{\text{LISA}}(f)$ represents the effective PSD due to the unresolved galactic binaries; the explicit expression is given in Equation (14) of Ref. [47], and we have assumed a 4-year mission for Figure 1. Here, our expression S_n^{LISA} (7) is smaller than Equation (1) in Ref. [47]

by a overall factor of 5; the difference arises simply because Equation (7) does not account for the sky-averaging.

2.2. Waveform Models

We employ as our BBH waveform model the frequency domain, “restricted” waveform in the stationary phase approximation for the inspiral phase, and the frequency domain, single-mode waveform for the ringdown phase. We shall restrict our waveform model to these two phases to simplify our analysis as far as possible; the complete IMR treatment at the level of waveform (using, for example, the “effective-one-body” (EOB) approach [48,49] and the “phenomenological” (IMRPhenom) model [50,51]) will be left for future work.

The GW waveform from a BBH inspiral in the frequency domain has the well-known form (see, for example, Ref. [52]):

$$\tilde{h}_{\text{Insp}}(f) = \mathcal{A} f^{-7/6} e^{i\Psi_{\text{Insp}}(f)}, \quad (8)$$

where \mathcal{A} is the “Newtonian” amplitude averaged over all sky positions and binary orientations (see, for example, Ref. [53]), so that:

$$\mathcal{A} \equiv \frac{2}{5} \sqrt{\frac{5}{24}} \pi^{-2/3} \frac{\mathcal{M}^{5/6}}{D_L}. \quad (9)$$

The waveform’s frequency domain phase, $\Psi_{\text{Insp}}(f)$, in the post-Newtonian (PN) approximation is given by (see, for example, Ref. [54]):

$$\Psi_{\text{Insp}}(f) = 2\pi f t_c - \Psi_c - \frac{\pi}{4} + \frac{3}{128\pi v^5} \left(\Delta\Psi_{3.5\text{PN}}^{\text{pp}} + \Delta\Psi_{3.5\text{PN}}^{\text{pp-spin}} + \Delta\Psi_{3.5\text{PN}}^{\text{BH-tidal}} \right), \quad (10)$$

where $v \equiv (\pi m t)^{1/3}$ is the PN parameter (in terms of the observer-frame total mass), and t_c and Ψ_c are the time and phase at coalescence. The phase terms $\Delta\Psi_{3.5\text{PN}}^{\text{pp}}$ and $\Delta\Psi_{3.5\text{PN}}^{\text{pp-spin}}$ are the 3.5PN spin-independent, point-particle contributions derived in Ref. [55] and the 3.5PN spin-dependent, point-particle contributions that include linear spin-orbit [56,57], quadratic-in-spin [58], and cubic-in-spin [59] effects, respectively. The remaining phase term $\Delta\Psi_{3.5\text{PN}}^{\text{BH-tidal}}$ is related to the tidal response of a spinning BH as a finite-size body—i.e., BH-absorption corrections such as the GW energy and angular momentum fluxes down to the horizons and the associated evolution of the BH itself [60–62].

Meanwhile, limited to only fundamental ($n = 0$), $\ell = 2 = m$ mode, the time-domain, single-mode ringdown waveform measured at a GW observatory is written as [63–65]:

$$h_{\text{Ring}}(f_c, Q, t_0, \phi_0; t) = \begin{cases} e^{-\frac{\pi f_c (t-t_0)}{Q}} \cos[2\pi f_c (t-t_0) - \phi_0] & \text{for } t \geq t_0, \\ 0 & \text{for } t < t_0, \end{cases} \quad (11)$$

where t_0 and ϕ_0 are the initial time and phase of the ringdown, respectively, and we have ignored the overall amplitude so that Equation (11) is not normalized; the initial ringdown amplitude may be determined by matching the ringdown GW waveform with the merger one (refer to, for example, Equation (16) below). When the final remnant object is a Kerr BH, the central frequency f_c and the quality factor Q are given in terms of quasinormal mode (QNM) frequencies ($f_{\text{QNM}} = f_{\text{R}} + i f_{\text{I}}$) of the remnant BH as (see Equation (7) of Ref. [66]):

$$f_c := f_{\text{R}}, \quad Q := -\frac{f_{\text{R}}}{2f_{\text{I}}}. \quad (12)$$

It should be noted that $f_{\text{QNM}} = f_{\text{QNM}}^{\text{r}}(1+z)^{-1}$ here is given in the observer frame. The Fourier transforms of Equation (11) provide the corresponding frequency-domain waveform, which takes the form (here we follow the convention of Ref. [63]):

$$\tilde{h}_{\text{Ring}}(f_c, Q, t_0, \phi_0; f) = \frac{Q(f_c \cos \phi_0 - 2iQf \cos \phi_0 + 2Qf_c \sin \phi_0)}{\pi(f_c + 2iQf_c - 2iQf)(f_c - 2iQf_c - 2iQf)} e^{2i\pi f t_0}. \quad (13)$$

We note that this frequency-domain GW waveform is not normalized either, following the approach presented in Appendix B of Ref. [64]. The maximized SNR over the initial ringdown phase ϕ_0 is discussed in Ref. [67].

The ringdown waveforms (11) and (13) require input data for the QNM frequency of the final remnant BH. Because we do not consider the complete IMR phase at the level of waveforms, we employ the numerical relativity (NR) remnant-fitting formulas (see, for example, Refs. [68–73] and references therein), from which the final mass M_f and spin \mathbf{S}_f of the remnant BH are consistently inferred for a given initial BH masses $m_{1,2}$ and dimensionless spin parameters $\vec{\chi}_{1,2}$ (in the inspiral phase) as:

$$M_f = M_f(m_1, m_2, \vec{\chi}_1, \vec{\chi}_2), \quad \chi_f \equiv \frac{|\mathbf{S}_f|}{M_f^2} = \chi_f(m_1, m_2, \vec{\chi}_1, \vec{\chi}_2). \quad (14)$$

The mass and spin of the remnant BH in the remnant formulas are derived in the isolated horizon framework (see, for example, Ref. [74]), not obtained from the ringdown GWs. The latter is used only for checking the internal consistency of the formulas.

We then generate the accurate numerical data of f_{QNM} of the inferred remnant BH with the Black Hole Perturbation Club (B.H.P.C.) code [75] to obtain f_c and Q (for accurate numerical data of f_{QNM} , see also Ref. [76], Emanuele Berti’s “Ringdown” website [77], Ref. [78], and the Black Hole Perturbation Toolkit [79]). In practice, it is also convenient to present f_c and Q by means of a compact analytical formula. Such a formula for the ($\ell = 2, m = 2, n = 0$) mode is obtained in Ref. [65] (by performing fits to the numerical QNM frequency data), and it reads:

$$f_c = \frac{1}{2\pi M_f} \left[f_1 + f_2 (1 - \chi_f)^{f_3} \right], \quad Q = q_1 + q_2 (1 - \chi_f)^{q_3}, \quad (15)$$

with $f_1 = 1.5251$, $f_2 = -1.1568$, $f_3 = 0.1292$, $q_1 = 0.7000$, $q_2 = 1.4187$ and $q_3 = -0.4990$.

2.3. Signal-to-Noise Ratio

We estimate the SNR of a complete IMR GW signal, making use of the simple frequency-domain, IMR “amplitude” model [9]. This model is motivated by the waveform amplitude in the IMRPhenomB model [80] and it is given by:

$$\text{IMR}(f) = \mathcal{A} \times \begin{cases} f^{-7/6} & \text{for } f < f_{\text{max}}, \\ f_{\text{max}}^{-1/2} f^{-2/3} & \text{for } f_{\text{max}} \leq f < f_{\text{R}}^{\text{r}}/(1+z), \\ \frac{f_{\text{max}}^{-1/2} [f_{\text{R}}^{\text{r}}/(1+z)]^{-2/3} [f_{\text{I}}^{\text{r}}/(1+z)]^2}{\{f - [f_{\text{R}}^{\text{r}}/(1+z)]\}^2 + [f_{\text{I}}^{\text{r}}/(1+z)]^2} & \text{for } f_{\text{R}}^{\text{r}}/(1+z) \leq f, \end{cases} \quad (16)$$

where the overall constant \mathcal{A} is chosen to be the (averaged) GW signal’s amplitudes in the inspiral phase (9). We set $f_{\text{max}} = 1/[6^{3/2}\pi(1+z)m_t^{\text{r}}]$ as the GW frequency at the innermost stable circular orbit (ISCO) of a test particle in the Schwarzschild spacetime with the total mass of BBH m_t^{r} , and f_{R}^{r} and f_{I}^{r} are the real and imaginary parts of the QNM frequency (i.e., $f_{\text{QNM}}^{\text{r}} = f_{\text{R}}^{\text{r}} + i f_{\text{I}}^{\text{r}}$) of the fundamental ($n = 0$), $\ell = 2 = m$ mode, which are determined by the final mass and spin of the remnant BH; the redshift dependence appears because the model parameters (m_t^{r} , f_{R}^{r} , f_{I}^{r}) are all given in the source’s rest frame (while the GW

frequencies f and f_{\max} are in the observer frame). The averaged SNR (in the above sense) can be then obtained by

$$\rho_{\text{ave}} = 2 \left\{ \int_{f_{\text{in}}}^{f_{\text{end}}} \frac{[\text{IMR}(f)]^2}{S_n(f)} df \right\}^{1/2}, \quad (17)$$

where $S_n(f)$ is the noise PSD and the frequency range $[f_{\text{in}}, f_{\text{end}}]$ is determined by the GW detector with which we observe the GW signals. For the given Fourier transform of a GW signal $\tilde{h}(f)$, the SNR can be written in terms of either $\tilde{h}(f)$ itself with the dimension of 1/Hz, the spectral density of the source amplitude $\sqrt{S_h(f)} \equiv 2\sqrt{f}|\tilde{h}(f)|$ with the dimension of $1/\sqrt{\text{Hz}}$, or the dimensionless characteristic strain $h_c(f) \equiv 2f|\tilde{h}(f)|$ [81]:

$$\rho_{\text{ave}} = \left(\int_{f_{\text{in}}}^{f_{\text{end}}} \frac{|2\tilde{h}(f)|^2}{S_n(f)} df \right)^{1/2} = \left(\int_{f_{\text{in}}}^{f_{\text{end}}} \frac{|\sqrt{S_h(f)}|^2}{S_n(f)} \frac{df}{f} \right)^{1/2} = \left(\int_{f_{\text{in}}}^{f_{\text{end}}} \frac{[h_c(f)]^2}{f S_n(f)} \frac{df}{f} \right)^{1/2}.$$

Note that the amplitude spectral density (i.e., the square root of the PSD of the source amplitudes) is $\sqrt{S_h(f)} \equiv 2\sqrt{f}|\text{IMR}(f)|$ in this model, and its track for the GW190521-like BBH is plotted in Figure 1 with the (square root of) the noise PSD $\sqrt{S_n(f)}$, assuming the median values of Table 1 and the remnant formulas of Ref. [68].

2.4. Multiband Fisher Analysis

We approximate the variance (i.e., uncertainty squared) associated with the measurement of a set of signal parameters, making use of the standard Fisher matrix formalism. The Fisher information matrix for a single-band GW detector is defined by:

$$\Gamma_{ab} \equiv \left(\frac{\partial \tilde{h}}{\partial \theta_a} \middle| \frac{\partial \tilde{h}}{\partial \theta_b} \right) \bigg|_{\theta=\theta_0}, \quad (18)$$

where $\tilde{h}(f, \theta)$ is the frequency-domain GW signal described by the set of parameters θ , and θ_0 are the best-fit values of the binary parameters. The bracket defines the noise-weighted inner product over the frequency range of $[f_{\text{in}}, f_{\text{end}}]$ [82]:

$$(a \mid b) \equiv 2 \int_{f_{\text{in}}}^{f_{\text{end}}} \frac{\tilde{a}^*(f)\tilde{b}(f) + \tilde{b}^*(f)\tilde{a}(f)}{S_n(f)} df, \quad (19)$$

where an asterisk '*' denotes the complex conjugation. The inverse Fisher matrix defines the corresponding variance-covariance matrix $\Sigma^{ab} \equiv (\Gamma_{ab})^{-1}$. In the limit of suitably high SNR [83], the variance of the parameter θ^a is given by:

$$\sigma_a^2 = \Sigma^{aa}. \quad (20)$$

In the case of the multiband analysis (to combine the information from, for example, aLIGO + B-DECIGO), we simply construct a multiband SNR and Fisher matrix by adding individual (averaged) SNRs and Fisher information matrices for each GW detector:

$$\rho_{\text{tot}}^2 \equiv \sum_I (\rho_{\text{ave}}^I)^2, \quad \Gamma_{ab}^{\text{tot}} \equiv \sum_I \Gamma_{ab}^I, \quad (21)$$

where ρ_{ave}^I and Γ_{ab}^I are the averaged SNR and the Fisher matrix for the I -th detector. The multiband variance-covariance matrix is defined by:

$$\Sigma_{\text{tot}}^{ab} \equiv (\Gamma_{ab}^{\text{tot}})^{-1}, \quad (22)$$

The variance of θ^a is then obtained by $\sigma_a^2 = \Sigma_{\text{tot}}^{aa}$.

3. Parameter Estimation Errors via Multiband Observation

In this section, we summarize the parameter estimation errors of the inspiral and ringdown phases for a nonprecessing, spinning, GW190521-like BBH system using the multiband GW network (LISA, B-DECIGO, aLIGO, and ET) detailed in Section 2.1. We follow Refs. [38,66] in our treatment of the Fisher matrix calculation for BBH GW signals (the setup here is slightly different from Ref. [38]; (i) we will assume a 5 yr observation rather than a 4 yr observation and (ii) we will use the new LISA sensitivity curve proposed by Ref. [47] and displayed in Equation (7), not the earlier eLISA sensitivity curve presented in Ref. [84]). We will also continue to neglect the contribution from the merger GW signal (using our IMR amplitude model in Equation (16), we have the (averaged) merger SNR of 17.0, 1.97×10^2 , and 3.42 for aLIGO, ET, and B-DECIGO, respectively, and the addition of these contributions to the signal analysis will improve the parameter estimations), for the reason that one cannot separate the inspiral, merger, and ringdown phases cleanly in the strict sense, although we have presented a simple IMR amplitude model in Equation (16). When the merger contribution is introduced into the inspiral or ringdown signal analysis, it causes some bias in the inspiral or ringdown parameter estimation. For example, the merger–ringdown waveform is parametrized by binary parameters, not solely by the remnant BH parameters after the merger. Similarly, in the inspiral–merger waveform, the “late” merger phase can be described by the overtones of QNMs [85]—i.e., the remnant BH parameters. Thus, we shall perform a inspiral–ringdown consistency test of GR without the merger contribution in Section 4.1.

3.1. Setup of Fisher Analysis

We set the default frequency interval of each GW detector $[f_{\text{low}}, f_{\text{up}}]$ as $[10.0, 3.0 \times 10^3]$ Hz (aLIGO), $[2.0, 3.0 \times 10^3]$ Hz (ET), $[0.01, 1.0 \times 10^2]$ Hz (B-DECIGO), and $[1.0 \times 10^{-4}, 1.0]$ Hz (LISA), respectively. We shall adopt the $T_{\text{obs}} = 5$ -year observation time and assume that the binary merges at the GW frequency of the Schwarzschild ISCO, $f_{\text{ISCO}} = 1/[6^{3/2}\pi(1+z)m_t^r]$. In this setup, the minimum frequency of the GW signal is:

$$f_{\text{min}} = 9.24 \times 10^{-3} (1+z)^{-5/8} \left(\frac{\mathcal{M}^r}{65.1 M_{\odot}} \right)^{5/8} \left(\frac{5 \text{ yr}}{T_{\text{obs}}} \right)^{3/8} \text{ Hz}, \quad (23)$$

where \mathcal{M}^r is the chirp mass in the source’s rest frame, normalized to that of GW190521. Therefore, the GW signal observed by each GW detector is truncated at the corresponding initial frequency $f_{\text{in}} \equiv \max(f_{\text{min}}, f_{\text{low}})$ as well as the end frequency $f_{\text{end}} \equiv \min(f_{\text{ISCO}}, f_{\text{up}})$; recall Figure 1.

The parameters of the inspiral waveform (8) are:

$$\theta_{\text{Insp}} = (f_0, t_c, \Psi_c, \ln m_t, \nu, \chi_s, \chi_a), \quad (24)$$

where we define the symmetric and anti-symmetric combinations of BH spins by $\chi_s \equiv (\chi_1 + \chi_2)/2$ and $\chi_a \equiv (\chi_1 - \chi_2)/2$ with the component (aligned) BH spins $\chi_{1,2} \equiv |\vec{\chi}_{1,2}| \equiv |\mathbf{S}_{1,2}|/m_{1,2}^2$. At the same time, the parameters of the ringdown waveform (13) are:

$$\theta_{\text{Ring}} = (t_0, \phi_0, f_c, Q). \quad (25)$$

It should be noted that the amplitude parameters are left out from the set of our independent parameters in both Equations (24) and (25). They are entirely uncorrelated with other parameters θ_a because the variance-covariance matrix Σ^{ab} gives the variance $\sigma_{\ln \mathcal{A}}^2 = \rho_{\text{ave}}^{-2}$ and the correlation $c^{\ln \mathcal{A}, a} \equiv \Sigma^{\ln \mathcal{A}, a} / (\sigma_{\ln \mathcal{A}} \sigma_a) = 0$ for the inspiral GW signal (8) (see, for example, Ref. [86]), and similar for the ringdown GW signal (if we explicitly introduce the amplitude to the normalized waveform of Equation (13)). For simplicity, we consider that all (other) parameters are unconstrained.

The best fit values of inspiral parameters are given by the median values in Table 1 with $t_c = 0.0 = \Psi_c$, while those of the ringdown parameters are assumed to be $\phi_0 = 0.0 = t_0$ and:

$$(f_c, Q) = (85.061 \text{ Hz}, 4.8354), \quad (26)$$

for the (observer-frame) central frequency and quality factor. Note here that the parameter estimation errors are independent of the value of the initial time t_0 . In our frequency-domain, single-mode ringdown waveform in Equation (13), the t_0 dependence is factorized as $e^{2i\pi f t_0}$ and it does not contribute to the noise-weighted inner product in Equation (19). On the other hand, the estimation errors depend on the initial phase ϕ_0 weakly [64].

The values in Equation (26) are obtained via Equation (12) for the given QNM frequency data of the remnant BH, assuming that the remnant mass M_f and spin χ_f of the final BH are inferred via the remnant formula in Equation (14) with the parameters of each component BH in Table 1. Specifically, we use the remnant formula provided by Ref. [71], and we quote:

$$\left(\frac{M_f}{m_t}, \chi_f \right) = (0.90356, 0.88269). \quad (27)$$

Note that these values are different from the mass and spin $(M_f, \chi_f) = (142_{-16}^{+28} M_\odot, 0.72_{-0.12}^{+0.09})$ of the remnant BH reported in the LIGO/Virgo GW190521 detection paper [3]; we have assumed that the individual spins of GW190521-like BBH are nonprecessing and completely aligned to the orbital angular momentum for simplicity, while the observed GW190521 is actually considered to be a precessing BBH.

The associated QNM frequency data of this remnant BH are then generated by the B.H.P.C. code [75], yielding the results in Equation (26).

The root-mean-square of parameter estimation error scales like the inverse of SNR, $\sim 1/\rho$, and it depends on both the ρ and bandwidth over which the SNR is accumulated. To see the benefit of having wider bandwidth due to the multiband observation, we finally introduce the normalized root-mean-square error as:

$$\delta\hat{\theta} \equiv \rho \sigma, \quad (28)$$

We display our error estimations in terms of $\delta\hat{\theta}$ with the total averaged SNR $\rho_{\text{ave}}^{\text{tot}}$ accumulated over the multi-frequency bands.

3.2. Result: Inspiral Phase

In Table 2, we present the parameter estimation errors of mass parameters (m, ν) and spin parameters (χ_s, χ_a) for the GW190521-like BBH inspiral in various combinations of ground/space-based GW observatories (while suppressing those of t_c and Ψ_c). Here, it should be noted that only the inspiral phase is analyzed, and the merger–ringdown phase which contributes to the SNR for ground-based observatories is ignored; recall the merger SNR shown in Section 3. Therefore, the SNR for aLIGO quoted here is much smaller than the observed LIGO/Virgo network SNR of 14.5 [3].

In the single-band case, as seen in Figure 1, the inspiral GW signal of the GW190521-like BBH is best observed by B-DECIGO because it can cover both the early (1 yr before merger) and late (around the ISCO) phases. However, B-DECIGO observation alone is not enough benefit to discern the spin parameters.

In the multiband cases, thanks to the wider bandwidth, the observation with B-DECIGO and ET gives a factor of 2 improvement in all the parameter estimation, even if we observe only the inspiral GW signal. This is further refined if we combine the data from LISA, forming a triple-band network (although the estimated LISA SNR of ~ 2.68 is likely too small compared to the ‘realistic’ detection SNR threshold of ~ 15 [87]). Importantly, the normalized errors of BH spins $(\delta\hat{\chi}_s/\chi_s, \delta\hat{\chi}_a/\chi_a) = (2.52, 5.71 \times 10^2)$ in this best case imply that the magnitudes of individual BH spins σ_χ can be recovered with the fractional

statistical errors (i.e., now accounting for the total multiband SNR of $\rho_{\text{ave}}^{\text{tot}} = 6.52 \times 10^1$ to the uncertainties; recall Equation (28)) ($\sigma_{\chi_1}/\chi_1, \sigma_{\chi_2}/\chi_2$) = $(2.43 \times 10^{-1}, 2.51 \times 10^{-1})$.

The joint LISA, B-DECIGO, and ET observatories would be only viable network to measure all the binary parameters of GW190521-like BBH system in the inspiral phase, including BH's component spins.

3.3. Result: Ringdown Phase

In Table 3, we show the parameter estimation errors of the central frequency f_c and quality factor Q (while suppressing those of t_0 and ϕ_0) for the ringdown phase of the GW190521-like BBH in ground-based observatories. For simplicity, we estimate the ringdown amplitude and associated SNR via the IMR amplitude model in Equation (16) (we should note that this approximation is likely too raw because the ringdown amplitude strongly depends on the starting time t_0 of the ringdown phase, which is difficult to determine in practice). We see that ET will be able to measure the QNM frequency of the remnant BH with the statistical error $\sim 10^{-3}$. At the same time, however, the difference in the normalized errors $\delta \hat{f}_c$ and $\delta \hat{Q}$ between the aLIGO and ET observations is not so evident. We speculate that this arises from the difference in the spectrum shapes (not the overall amplitudes) because the ringdown waveform (13) is narrow-banded in the frequency domain.

4. The Implications for Tests of GR via Multiband Observation

In this section, we explore to what extent the multiband observation of the GW190521-like BBH system discussed in Section 3 could improve tests of GR. A handful of tests have been already formulated and performed with merging BBH systems (see, for example, Refs. [88–91]), and we follow the (very) simple tests proposed by Nakano et al. [66,92].

4.1. A Consistency test of GR with the Inspiral and Ringdown GW Signals

One possible test of GR with a BBH system is to establish the consistency of the mass and spin of the final remnant BH determined by two different parts of the GW signals. Thanks to the recent advancement in NR simulations of BBH systems [93–95] (see also Refs. [96–100]), one can infer these values from the initial component masses and spins measured from the inspiral GW signal (in the low-frequency band), making use of the NR fitting formulas for the remnant properties of the final BH; recall Section 2.2. At the same time, they are directly estimated from the succeeding merger–ringdown GW signal (in the high-frequency band). This type of test is now known as the “IMR consistency test” [101,102]. By formulation, multiband observations of heavy BBH mergers such as GW190521-like BBH systems will be “golden binaries” [103] of such an IMR consistency test.

Given that the early inspiral and late ringdown GW signals will be best observed in a different frequency band (such as “B-DECIGO + aLIGO” network etc.), we here perform the multiband version of the “inspiral–ringdown consistency test” formulated by Nakano et al. [66] (see also Refs. [103,104]), solely using the inspiral and ringdown parts, and test the consistency of GR across the merger part, which is a highly dynamical phase in a strong-field regime. We estimate the statistical errors on M_f and χ_f from the inspiral GW signal by using the (normalized) statistical errors $\delta \hat{\theta}$ in Table 2 and applying a standard variance propagation of non-linear functions to the specific NR remnant formulas (“UIB formulas”) [71] publicly available in LALInference [105,106]. Here, we ignore the systematic bias due to our specific choice of the NR remnant formulas, for simplicity (see, for example, Refs. [99,102] for details). The errors from the ringdown GW signal are estimated from Table 3, through the dependence of f_c and Q on (M_f, χ_f) . Technically, this procedure requires the evaluation of the partial derivatives $(\partial f_c / \partial \chi_f)_{M_f}$ etc. to compute the variance propagation. We construct the numerical function of QNM frequencies $M_f f_{\text{QNM}}(\chi_f)$ with the B.H.P.C. code [75] around the best fit values of Equation (27), from which these derivatives are extracted.

In Table 4, we present the parameter estimation errors of the remnant mass M_f and spin χ_f from both the inspiral and ringdown phases. For the reference, we also present the same errors of the GW150914-like BBH with the redshifted component masses $(m_1, m_2) = (30 M_\odot, 40 M_\odot)$, spin magnitudes $(\chi_1, \chi_2) = (0.9, 0.7)$, and the luminosity distance $D_L = 0.4$ Gpc (i.e., the redshift $z \sim 0.085$). This system was analyzed as “System B” in Ref. [38]. We see that the “LISA + B-DECIGO + ET” observation allows us the IMR consistency test at the sub-percent precision both for the GW150914-like, and GW190521-like BBH systems. The result of the GW190521-like BBH system is slightly worse than the case of the GW150914-like BBH system because the luminosity distance of GW190521 ($D_L = 5.3$ Gpc) is much larger than that of GW150914 ($D_L = 0.4$ Gpc) [22,107,108]. Additionally, the GW190521-like BBH system has the central frequency $f_c \sim 85$ Hz lower than that of the GW150914-like BBH system ($f_c \sim 350$ Hz), missing the most sensitive frequency of aLIGO and ET ~ 250 Hz. If a GW190521-like BBH was observed at the same distance as GW150914, the accuracy of its inspiral–ringdown consistency test would be at the level of $\sim O(0.01\%)$.

4.2. A Simple Test of the Remnant Compact Object with Quasinormal Modes

Another simple test of GR is to bracket whether the remnant object should be a BH predicted by GR or not, making use of the parameter estimation errors of QNM frequencies ($f_{\text{QNM}} = f_R + i f_I$) obtained from the ringdown GW signals [66].

Figure 2 plots the 1σ , 2σ , 3σ , 4σ , and 5σ error contours on the fundamental ($n = 0$), $\ell = 2 = m$ mode of the QNM frequency in the (f_R, f_I) plane, in the case of the GW190521-like BBH system observed by aLIGO (left) and ET (right). The errors are estimated through the results in Table 3 concerning the parameter estimation error on the ringdown GW signal (after t_0 and ϕ_0 being marginalized out), and the outermost contour in each panel shows the 5σ error. The black lines in each panel depict the Schwarzschild limit of $|f_I|/f_R$, which may be obtained by setting $\chi_f = 0$ in Equation (15) with Equation (12):

$$\frac{|f_I|}{f_R} \approx 0.236. \quad (29)$$

This is marginally consistent with the exact QNM frequency of the fundamental ($n = 0$), $\ell = 2 = m$ mode in the Schwarzschild limit. For example, the B.H.P.C. code gives $|f_I|/f_R = 0.23808 \dots$, and the difference is negligibly small here. The key point of this test is that the top-left side of the black line becomes the prohibited region in GR—i.e., for the ($n = 0$), $\ell = 2 = m$ mode the QNM frequencies (f_R, f_I) of any rotating Kerr BHs must sit in the bottom-right side of the black line.

In the aLIGO case, due to the low SNR ($= 10.8$) the parameter estimation errors already go beyond the Schwarzschild limit at the 3σ level, and we cannot confirm whether the remnant object is a BH predicted by GR with the 5σ level. For such low SNR events, the “coherent mode stacking method” [109] will be useful. On the other hand, thanks to the high SNR ($= 147$) in the ET case, this simple test can necessarily confirm that the remnant object is a GR-predicted BH. Note that the (M_f, χ_f) plane has been discussed as Figure 5 in Ref. [3]. The remnant BH spin is restricted to $0 \leq \chi_f < 1$ in the analysis. Therefore, the simple test presented in this paper is not applicable.

5. Summary and Discussion

This work underlines the multiband observation of the GW190521-like nonprecessing, quasicircular “intermediate-mass” BBH with LISA and B-DECIGO in the low-frequency band, combining it with aLIGO and ET in the high-frequency band. Our first result for the parameter estimation errors is displayed in Tables 2 and 3; the statistical errors of the binary’s mass parameters by B-DECIGO observation will be $\sim 10^{-2}$, and the multiband observation with LISA, B-DECIGO, and ET will further improve them to a factor of 2, even allowing the statistical errors of component BH spins at the accuracy level of $\sim 10^{-1}$. Based on the ringdown analysis, ET will measure the QNM frequency with about $O(0.1)\%$

precision. Our second result for the GR tests is presented in Table 4 as well as Figure 2. We showed that the multiband observation of the GW190521-like BBH system by LISA, B-DECIGO, and ET would perform the inspiral–ringdown consistency test at a percent level of precision.

The main point of our analysis is that there is the principal advantage of measuring GW190521-like BBH systems using *the full GW spectrum* from milli-Hz to deci-Hz bands, and to hecto-Hz band. We expect that our findings will motivate further investigation on prospects for the multiband observation of GW190521-like BBH systems, which are much alike the prototypical GW150914-like BBH systems [22].

Nevertheless, we emphasize that our study was performed with (very) simple methods. Therefore, our results should be only *indicative* and *tentative*. In the remainder of this section, we discuss what work remains to be carried out to refine our analysis in the future, so that we can eventually make a strong scientific case for multiband GW astronomy/physics.

Assessment of Prospects

First, our methodology in Section 2 should be replaced with more modern, sophisticated approaches to the GW data analysis; it will include (but not be limited to), for example, the use of complete IMR waveforms such as EOB approach as well as IMRPhenom models (see, for example, Ref. [110] and the references therein), full-fledged Bayesian posterior-based techniques (see, for example, Refs. [105,111,112]), and more ‘realistic’ noise and waveform models that account for the sky-location of sources as well as orbital configurations of B-DECIGO [11] and LISA [8].

Second, our target GW190521-like BBH system was restricted to the nonprecessing, spinning, quasicircular configuration; although there is large uncertainty, the spin precession of GW190521 is estimated as nonzero [3,4]. Future studies are therefore needed to concern both the spin precession and orbital eccentricity of the BBH system: we will briefly discuss the eccentric, non-spinning BBH system in Appendix A. In general, the BH’s spins in the precessing binary have not only their magnitudes but also orientations, which can be described by, for example, the effective inspiral spin parameter χ_{eff} (related to the components aligned with the orbital angular momentum) and the precession spin parameter χ_p (related to the components in the orbital plane for the inspiral GW waveform). Adding the orbital eccentricity to the (precessing) BBH system will be fully generic, and the waveform modeling becomes (much) more complicated. Despite that challenge, there is a considerable development of the analysis on fully generic binary systems [113].

Related to the point mentioned above, it should be noted that we have ignored subdominant ($\ell \neq 2, |m| \neq 2$) harmonics in our GW waveform model. They are more notable in the observed signal when system’s mass ratio becomes smaller (such as the analysis of GW190814 [114] with the mass ratio $q = 0.112^{+0.008}_{-0.009}$). With the subdominant harmonics, one can access the source orientation to reduce uncertainty in the distance estimation [115].

Third, we should note that there are two main hurdles to analyze the ringdown GW signals; the low SNR with aLIGO and the starting time of the ringdown phase that is *a priori* unknown to the whole observed GW signal (see, for example, Refs. [116,117] for discussions on the starting time). Our simple analysis is carried out with the single-mode waveform model (11) as the template to analyze the ringdown phase, assuming the starting time of the ringdown phase to be $t = t_0$, which corresponds to $f = f_R$ —i.e., just after the end of the merger phase. While this choice does not affect our results of parameter estimation errors, in practice the best fit values (that should be obtained rather than assumed in the context of the full-parameter estimation against the raw GW data) can be biased if one assumes the earlier starting time in the analysis (see, for example, Figure 5 in Ref. [88]). Although one may delay the starting time to avoid the bias in the parameter estimation, the SNR becomes much lower than the expected SNR with the assumption of $t = t_0$.

These obstacles will be overcome if one includes higher overtones ($n > 0$) into the ringdown GW analysis [85], for which a much larger SNR than the single-mode analysis will be expected. Indeed, a superposition of overtones ($n > 0$) in a single harmonic mode will be observed [85,118–120] in the high SNR events (with ET). Another step-functional improvement of the ringdown GW analysis will be offered by using completely different signal analysis methods than the traditional matched filtering analysis, which may not be always optimal for the ringdown GW signal. There is ongoing work to assess the improvement due to new techniques for the ringdown GW signal analysis (such as Hilbert–Huang transformation, autoregressive modeling, and neural networks) [121,122].

Fourth, there are other GR tests that were not covered in Section 4 but that can be greatly improved using the multiband observation of GW195021-like BBH systems. Reference [91] performs various tests of GR with the BBH events in GWTC-2 (see also the reviews by Carson and Yagi [123] about the current and future test with GWs), including (i) the IMR consistency test between the inspiral and postinspiral phases divided at some cutoff frequency; (ii) constraining deviations from GR with parametric deformations to a predicted GR waveform model, (iii) “BH spectroscopy” [124,125] with ringdown GWs which contains two (or possibly more) QNMs [126,127], and so on.

The test (i) is similar to our inspiral–ringdown consistency test directly using the information of the merger phase, and makes the most of the GW waveform. It has been pointed out that the values of this test with GW150914-like BBH systems will be maximized in the multiband observation [28,29,128]. Our result suggests that the same will be true for the GW190521-like “intermediate-mass” BBH systems, too. Similarly, Gupta et al. [31,32] have showed that the multiband observations of stellar- and intermediate-mass BBHs with LISA and 3G detectors will be only workable way to carry out the most general version of test (ii). Adding B-DECIGO (or any other planned GW detectors) in the deci-Hz band to the multiband analysis, we expect the precision of this test will be unprecedented.

Based on the generalized likelihood ratio test [129], the test (iii) is performed (see, for example, Section 9.5 in Ref. [130] and references therein, and also Ref. [131] for the future O5 era). This test with aLIGO and Advanced Virgo alone is quite challenging to have any conclusive result, simply because of the too low SNRs and larger parameter estimation errors in the ringdown phase. In the 3G era, two (or possibly more) QNMs will be measurable [126,127]. It is also helpful to use the multiband observation in order to optimize ground-based detectors via the forewarnings from the low-frequency, LISA band [132]. In either cases, the multiband observation with B-DECIGO and ET will give the additional advantage of being able to perform the best test of Kerr hypothesis of remnant BHs via BH spectroscopy.

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Appendix A. Signal-to-Noise Ratio of GW190521-Like Eccentric BBH Systems

Throughout the bulk of this paper, we have looked at the GW190521-like BBH system under the assumption of a quasicircular BBH merger. While the quasicircular evolution of GW190521 is totally consistent with the LIGO/Virgo observation [3,4], due to the lack of the inspiral GW signal long enough, it appears that alternative scenarios, for example, GW190521 as an eccentric BBH merger also become relevant. Indeed, Gayathri et al. [44] demonstrated that observed GW190521 data could be explained as an equal-mass, highly eccentric ($e = 0.7$) BBH system (the orbital eccentricity ($e = 0.7$) is provided as initial data for NR simulations at a frequency of 10 Hz for a system with the total mass of $50 M_{\odot}$ that is the orbital separation $\sim 24.5 m_t$ for BBHs with the total mass m_t [133] and note that the NR waveforms are scaled by the total mass m_t), and the estimated source parameters are quite different from those derived from the quasicircular BBH scenario (recall Table 1): the primary mass $m_1^r = 102_{-11}^{+7} M_{\odot}$, the secondary mass $m_2^r = 102_{-11}^{+7} M_{\odot}$, and the total mass $m_t^r = 204_{-33}^{+14} M_{\odot}$ in the source's rest frame, the mass ratio $q = 1$, the luminosity distance $D_L = 1.84_{-0.054}^{+1.07}$ Gpc and the redshift $z = 0.35_{-0.09}^{+0.16}$ (see Ref. [44] for the spin parameters). The possibility of GW190521 with nonvanishing eccentricity is also pointed out by Refs. [134,135].

Like the quasicircular case, the multiband observation of eccentric GW190521-like BBH systems will once again help in distinguishing these two scenarios. Assuming a quadrupole GW generation from a Newtonian Kepler orbit [136,137], the typical coalescing time t_c^{ecc} and the characteristic strain amplitude $h_{c,n}^{\text{ecc}}$ of the n -th harmonic are (see, for example, Section 4.1 of Maggiore's text [138] as well as Ref. [139]):

$$t_c^{\text{ecc}} \sim t_c^{\text{circ}} (1 - e_0^2)^{7/2}, \quad h_c^{\text{ecc}} \sim h_c^{\text{circ}} g(n, e), \quad (\text{A1})$$

where t_c^{circ} and h_c^{circ} are corresponding circular-inspiral results given in Equations (1) and (3), respectively, and the function $g(n, e)$ will be defined momentarily. We see that the early inspiral phase of the eccentric binary is well within the B-DECIGO and LISA bands, too. Specifically, Holgado et al. [140] (see also the works by Amaro-Saoane [42,139,141]) pointed out that having deci-Hz GW observatories such as B-DECIGO, MAGIS [14] and TianGO [15] will be a key element to observe the eccentric inspiral GW signals in multiband networks because the GW signals may skip the LISA band entirely; $h_{c,n}^{\text{ecc}}$ can be suppressed by a function of $g(n, e)$ significantly, depending on its eccentricity in the LISA band.

To better understand the visibility of eccentric BBH systems with aLIGO, ET, B-DECIGO, and LISA, let us estimate the SNR of non-spinning, eccentric BBH inspirals accumulated in each band; see also Ref. [142] for a recent review about waveform families for the eccentric binary systems. We also see various active works on eccentric waveform approximants [143–146]. The squared SNR averaged over the all-sky positions and binary orientations may be written as [23,53,81,147]:

$$\rho_{\text{ave}}^2 \approx \frac{1}{5} \sum_n \int_{f_{\text{in}}}^{f_{\text{end}}} \frac{h_{c,n}^2}{f_n S(f)} d(\ln f_n), \quad (\text{A2})$$

where $S(f)$ is the noise PSD for a given GW detector, and

$$f_n = n f_{\text{orb}} = n f_{\text{orb}}^r (1 + z)^{-1}, \quad (\text{A3})$$

is the frequency of the harmonic in the observer frame, defined by the source's rest-frame frequency f_{orb}^r of the Kepler orbit with the redshift z . The expression for the average SNR (A2) recovers Equation (17) (of the inspiral part) in the circular orbit limit, $e \rightarrow 0$ (note that the harmonics are restricted to only $n = 2$ in the circular limit because one has $\lim_{e \rightarrow 0} g(n, e) = \delta_{2,n}$). The dimensionless characteristic strain $h_{c,n}$ of the n -th harmonic is [147]:

$$h_{c,n} \equiv \frac{1+z}{\pi D_L} \sqrt{2 \frac{dE_n^r}{df_n^r}}, \quad (\text{A4})$$

where D_L is the luminosity distance. Assuming the quadrupole formula applied to the Kepler orbit (with the chirp mass \mathcal{M} in the observer frame) [136,137],

$$\frac{dE_n^r}{df_n^r} \equiv \frac{\pi}{3} \frac{\mathcal{M}^{5/3}}{(1+z)^2} \frac{g(n, e)}{F(e)} \left(\frac{2}{n}\right)^{2/3} (\pi f_n)^{-1/3}, \quad (\text{A5})$$

is the emitted GW energy per unit frequency f_n^r at the n -th harmonic measured in the source's rest frame, and we define:

$$g(n, e) \equiv \frac{n^4}{32} \left\{ \left[J_{n-2}(ne) - 2e J_{n-1}(ne) + \frac{2}{n} J_n(ne) + 2e J_{n+1}(ne) - J_{n+2}(ne) \right]^2 + (1-e^2) [J_{n-2}(ne) - 2J_n(ne) + J_{n+2}(ne)]^2 + \frac{4}{3n^2} J_n^2(ne) \right\}, \quad (\text{A6})$$

$$F(e) \equiv \frac{1 + \frac{73}{24}e^2 + \frac{37}{96}e^4}{(1-e^2)^{7/2}} \left(= \sum_{n=1}^{\infty} g(n, e) \right), \quad (\text{A7})$$

with the Bessel functions of the first kind $J_n(x)$ (n : integer); see, for example, Ref. [148] for the derivation of Equation (A5).

For an inspiraling eccentric BBH, the evaluation of SNR through Equation (A2) requires the knowledge of the slowly evolving orbital eccentricity e and the frequency f_n in time, under the gravitational radiation losses. Again, in the quadrupole formalism, this is given by [136,137]:

$$\frac{f_{\text{orb}}}{f_{\text{orb},0}} = \left[\frac{1-e_0^2}{1-e^2} \left(\frac{e}{e_0} \right)^{12/19} \left(\frac{1 + \frac{121}{304}e^2}{1 + \frac{121}{304}e_0^2} \right)^{870/2299} \right]^{-3/2}, \quad (\text{A8})$$

with reference eccentricity e_0 and orbital frequency $f_{\text{orb},0}$ (in the observer frame). One can set these constants by the values at the last stable orbit (LSO) of the eccentric geodesic (of a test particle) in the Schwarzschild geometry [149]: namely, $e_0 = e_{\text{LSO}}$ and:

$$f_{\text{orb},0} = f_{\text{LSO}} = \frac{1}{2\pi(1+z)m_t^r} \left(\frac{1-e_{\text{LSO}}^2}{6+2e_{\text{LSO}}} \right)^{3/2} \sim 8.0 (1+z)^{-1} \left(\frac{m_t^r}{204M_{\odot}} \right)^{-1} \left(\frac{1-e_{\text{LSO}}^2}{6+2e_{\text{LSO}}} \right)^{3/2} \text{Hz}, \quad (\text{A9})$$

with the total mass m_t^r in the source's rest frame. The frequency evolution (A8) is therefore completely determined with a given single parameter e_{LSO} .

We compute the averaged SNR ρ_{ave} given in Equation (A2) for the GW190521-like, non-spinning, eccentric BBHs with the source-frame masses $(m_1^r, m_2^r) = (102 M_{\odot}, 102 M_{\odot})$ and the luminosity distance $D_L = 1.9 \text{ Gpc}$ (i.e., the redshift $z \sim 0.35$), which mimics the eccentric BBH merger obtained by the NR simulations in Ref. [44]. We assume the five year observation prior to the final merger determined by Equations (A1) and (A9), and apply the same setup described in Section 3.1 to each n -th harmonic of GW strains. Because the frequency evolution (A8) is expressed in term of eccentricity, in practice, we change the integration variable of Equation (A2) from f_n to e for the computational efficiency, making use of $df_n = n|df_{\text{orb}}/de| de$ with (see, for example, Ref. [150]):

$$\frac{df_{\text{orb}}}{de} = -\frac{18}{19} \frac{f_{\text{orb}}}{e} \frac{F(e)}{(1-e^2)^{9/2} \left(1 + \frac{121}{304}e^2 \right)}. \quad (\text{A10})$$

Additionally, the infinite summation over the harmonics n in Equation (A2) is truncated at some finite value of n_{max} ; we chose $n_{\text{max}} = O(10^4)$ so that the resultant SNRs are

guaranteed to have at least 3 significant digits (a good analytical estimator of n_{\max} can be found in, for examples, Refs. [151,152]).

In Table A1, we summarize the averaged SNRs ρ_{ave} of the GW190521-like, non-spinning, eccentric BBH accumulated in each GW band, for sample values of final eccentricities at the last stable orbit $e_{\text{LSO}} = \{10^{-6}, 10^{-3}, 0.1, 0.4, 0.6\}$. For references, the ‘initial’ eccentricities at f_{in} for each band are also listed. It is approximated by solving Equation (A8) for e at the detector’s initial GW frequency of the second harmonics $f_{\text{orb}} = f_{\text{in}}/2$. We found three main results: (i) B-DECIGO and ET always accumulate a total SNR greater than at least 10, independent of the value of the final eccentricity e_{LSO} ; (ii) LISA has the detectable SNR only when $e_{\text{LSO}} < 10^{-3}$. That is, the inspiral GW signal from the GW190521-like BBH that has the high-eccentricity in the aLIGO band would entirely skip the LISA band; (iii) the SNR with B-DECIGO becomes bigger when e_{LSO} becomes smaller, while the SNR with ET and aLIGO shows the opposite behavior. Therefore, the multiband observation could provide much louder SNR than the single-band SNR across the full range of e_{LSO} . Using Equation (21) with the result in Table A1, we find that the multiband SNR with B-DECIGO and ET are ~ 180 for $e_{\text{LSO}} = \{10^{-6}, 10^{-3}\}$, ~ 140 for $e_{\text{LSO}} = \{0.4, 0.6\}$ and ~ 100 for $e_{\text{LSO}} = 0.1$.

Table A1. Averaged SNRs ρ_{ave} of the GW190521-like, non-spinning, eccentric BBHs accumulated in each band five years prior to coalescence, for given values of final eccentricity $e = e_{\text{LSO}}$. We assume the source-frame component masses $(m_1^r, m_2^r) = (102 M_{\odot}, 102 M_{\odot})$, and the luminosity distance $D_L = 1.9$ Gpc (i.e., the redshift $z \sim 0.35$) [44]. The values in parentheses indicate the ‘initial’ eccentricity estimated from the initial GW frequency (f_{in} defined in Section 3) of the second harmonic mode in each detectors. Note that aLIGO’s ‘initial’ eccentricity when $e_{\text{LSO}} = 0.6$ is not displayed because the second harmonic mode is not detectable in this case.

SNR and Eccentricity at f_{in}					
	$e_{\text{LSO}} = 10^{-6}$	$e_{\text{LSO}} = 10^{-3}$	$e_{\text{LSO}} = 0.1$	$e_{\text{LSO}} = 0.4$	
aLIGO	4.95 (1.64×10^{-6})	4.95 (1.64×10^{-3})	5.55 (0.150)	1.07×10^1 (0.405)	1.11×10^1 (\dots)
ET	7.64×10^1 (8.98×10^{-6})	7.63×10^1 (8.97×10^{-3})	8.73×10^1 (0.505)	1.42×10^2 (0.746)	1.42×10^2 (0.793)
B-DECIGO	1.67×10^2 (2.41×10^{-3})	1.70×10^2 (0.709)	5.08×10^1 (0.982)	1.85×10^1 (0.992)	1.20×10^1 (0.993)
LISA	7.60 (3.86×10^{-3})	2.05 (0.985)	<1.00 (>0.999)	<1.00 (>0.999)	<1.00 (>0.999)

Finally, although the parameter estimation is not the main focus here, we briefly discuss the potential accuracy of the eccentricity measurement. In the small-eccentricity and high-SNR limit, the orbital eccentricity may be measured within the fractional error [153] (see also Ref. [154]):

$$\delta\hat{e}_0 \sim 5 \times 10^{-5} \frac{(1+z)^{5/3}}{\sqrt{2+3\alpha}} \left(\frac{\mathcal{M}^r}{65.1 M_{\odot}} \right)^{5/3} \left(\frac{f_0}{0.1 \text{Hz}} \right)^{5/3} \left(\frac{e_0}{0.1} \right)^{-1}, \quad (\text{A11})$$

from Equation (2) in Ref. [153] by using a rough approximation with the quasicircular amplitude and the eccentric phase, where $\delta\hat{e}_0$ denotes the parameter estimation error of e_0 (at the GW frequency f_0 for $n = 2$) normalized by the SNR. Here, the power α is due to the approximation of the noise PSD by the power law, $S_n \sim f^{2\alpha}$ (assuming $\alpha > -2/3$), and B-DECIGO may have $\alpha = 1$, for example. This estimator implies that B-DECIGO would be able to precisely measure the eccentricity of GW190521-like BBH systems. Because we find that the ‘B-DECIGO + ET’ combination always provides the multiband SNRs larger than ~ 100 independent of the value of e_{LSO} , one might expect that this combination best

observes GW190521-like BBHs over the full range of eccentricity, helping to understand the population properties of BBH mergers [155]. We will explore this possibility in future work.

Appendix B. Some More Noise Power Spectral Densities of (Next-Generation) GW Detectors

In this appendix, we summarize some (fitting) curves of the noise PSD for both ground and space-based, current and future GW detectors. These sensitivity curves are not used in the bulk of this paper, but it will serve as a convenient all-in-one-place summary with our notation; these curves with f in unit of Hz are shown in Figure 4 of Ref. [156] (except DECIGO and TianQin).

- “LIGO O3a-Livingston” rough fitting curve (during the first half of LIGO/Virgo third observing run by using Ref. [157]):

$$S_n^{\text{O3a-L}} = \left(2.13068 \times 10^{-12} f^{-7.938724592} + 4.0 \times 10^{-22} f^{-1.0} \right. \\ \left. + 3.0 \times 10^{-24} + 1.74546 \times 10^{-27} f^{1.178746922} \right)^2 \text{Hz}^{-1}. \quad (\text{A12})$$

- “LIGO O5” rough fitting curve (will be in the fifth observing run by using Ref. [157]):

$$S_n^{\text{O5}} = \left(480985000.0 f^{-30.28419138} + 6.63263 \times 10^{-20} f^{-3.122716032} \right. \\ \left. + 6.15101 \times 10^{-21} f^{-2.089976737} + 1.32853 \times 10^{-27} f^{1.059219544} \right)^2 \text{Hz}^{-1}. \quad (\text{A13})$$

- “ET-B” (another sensitivity curve for ET in Ref. [158] other than Equation (5); see also, for examples, Ref. [159] and ET sensitivity page [160]):

$$S_n^{\text{ET-B}} = 1.0 \times 10^{-50} \left(45540.5 f^{-15.64} + 6804.96 f^{-2.145} \right. \\ \left. + 3.05853 f^{-0.12} + 0.00258062 f^{1.1} \right)^2 \text{Hz}^{-1}. \quad (\text{A14})$$

- “CE2” rough fitting curve (for Cosmic Explorer presented in Ref. [21]):

$$S_n^{\text{CE2}} = \left(1.74408 \times 10^{-16} f^{-8.908164528} + 2.0 \times 10^{-25} + 8.23008 \times 10^{-32} f^{2.095903274} \right)^2 \text{Hz}^{-1}. \quad (\text{A15})$$

- “DECIGO” (the noise PSD of the L-shaped configuration [161]):

$$S_n^{\text{DECIGO}} = \left\{ 7.05 \times 10^{-48} \left[1 + \left(\frac{f}{f_p} \right)^2 \right] + 4.8 \times 10^{-51} f^{-4} \left[1 + \left(\frac{f}{f_p} \right)^2 \right]^{-1} \right. \\ \left. + 5.53 \times 10^{-52} f^{-4} \right\} \text{Hz}^{-1}. \quad (\text{A16})$$

with $f_p \equiv 7.36$. Note that this expression accounted for the factor of $(\sqrt{3}/2)^{-2}$ due to DECIGO having arms that make an angle of 60° .

- “TianQin” [13,162] (the sky averaged noise PSD; the expression below is quoted from Equations (9) and (10) of Ref. [163]):

$$S_n^{\text{TianQin}} = 3.0 \times 10^{-51} \left(0.009505539123 f^{-5} + 95.05539123 f^{-4} + 0.07550033531 f^{-3} \right. \\ \left. + 755.0033531 f^{-2} + 3.703703703 \times 10^{10} \right. \\ \left. + 2.941767614 \times 10^{11} f^2 \right) \text{ Hz}^{-1}. \quad (\text{A17})$$

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