



Latest Developments in Differential Distributions at NNLO

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Abstract. Differential distributions for on-shell top-quark pair production at next-to-next-to-leading order in perturbative QCD for the Large Hadron Collider at $\sqrt{s} = 8$ TeV are presented and discussed. The theoretical framework of the calculation is summarized. By comparing the prediction to measurements, it turns out that next-to-next-to-leading order corrections are necessary in order to resolve a discrepancy between data and next-to-leading order predictions in the transverse momentum distribution of the top-quark.

INTRODUCTION

Top-quarks provide an excellent testing ground for the Standard Model and in particular Quantum Chromodynamics (QCD) at the Large Hadron Collider (LHC). The large number of top-quark pairs that have been produced at $\sqrt{s} = 7$ TeV and $\sqrt{s} = 8$ TeV allows for detailed studies of top-quark properties in different kinematical regions. Additionally, they are an important background for the main physics program measurements at the LHC at $\sqrt{s} = 13$ TeV, namely precision measurements of the Higgs-boson properties and searches for physics beyond the Standard Model. Precision predictions at a differential level for top-quark pairs are demanded in order to give a reliable interpretation of the measurements within and beyond the Standard Model.

In order to test the Standard Model down to the precision that can be achieved at the LHC, theoretical predictions at the same accuracy need to be provided. The total inclusive cross-section at next-to-next-to-leading order QCD [1, 2, 3, 4, 5, 6] shows a very good agreement with measurements at 7, 8 and 13 TeV [7, 8, 9, 10, 11]. Furthermore, only at this order of perturbation theory the theoretical uncertainty is at the same level as the precision of the measurement.

At a differential level, most measurements of the decay products of the top-quark, namely leptons and jets, are currently in good agreement with next-to-leading order predictions obtained using available Monte Carlo event generators. In contrast, the p_T -distribution of the top-quark itself reveals a discrepancy between measurements and next-to-leading order predictions at $\sqrt{s} = 7, 8$ and 13 TeV [12, 13, 14, 15, 16, 17, 18]. Since the top-quark is not measured directly, but reconstructed from its decay products using Monte Carlo generators, inaccuracies in the Standard Model description of the top-quark level process could cause a tensions. Higher-order QCD corrections will help to identify the origin of those discrepancies.

Several approximations of the next-to-next-to-leading order QCD corrections to differential top-quark pair production have been presented [19, 20, 21, 22, 23, 24]. These approximations are next-to-leading order accurate and, in addition, they capture contributions of the next-to-next-to-leading order result in specific kinematical regions, e.g. the partonic threshold of the top-quark pair.

In this write-up, full next-to-next-to-leading order differential distributions for top-quark pair production at the LHC at 8 TeV are discussed. First, an overview of next-to-next-to-leading order computations for top-quark pair production is presented. This is followed by results for the p_T -, $m_{t\bar{t}}$ -, $y_{t\bar{t}}$ - and y_t -distribution. Finally, the results are summarized and an outlook for further studies is given. For additional details please refer to the main publication [25].

Next-to-next-to-leading order computations for top-quark pairs

In order to perform numerical computations beyond leading order, a subtraction framework is needed to consistently cancel soft and collinear singularities between virtual and real contributions to the cross section. At next-to-leading

order general algorithms are established and allow an automated simulation of a large class of processes, that are relevant for the phenomenology at the LHC [26, 27]. At next-to-next-to-leading order different subtraction schemes have been proposed and applied to single processes.

In the context of top-quark pair production three approaches can be distinguished: The *antenna-subtraction* scheme has been used to calculate differential distributions in the leading colour approximation for the partonic contribution $q\bar{q} \rightarrow t\bar{t}$ [28, 29]. The *q_T-subtraction* scheme has been applied to top-quark pair production in [30] and results for the $q\bar{q}'$ -channel to the total inclusive cross section have been obtained.

The sector improved residue subtraction scheme, STRIPPER, has been proposed in [31] and subsequently successfully applied to the calculation of the total inclusive cross section for top-quark pair production [1, 2, 3, 4, 5, 6, 32]. The result includes all partonic processes without further approximations except the truncation of the perturbation series beyond next-to-next-to-leading order. Afterwards, the same framework has been used to predict the forward-backward asymmetry at the Tevatron [33], which has been the first differential prediction for on-shell top-quark pair production at next-to-next-to-leading order.

However, the first formulation of the subtraction scheme and its implementation for the Tevatron setup were not suited to provide robust differential results for the LHC on an adequate time-scale. A substantial diminution of the convergence of the Monte Carlo integration has been observed as the collider energy has been increased. This behaviour is due to large logarithms of ratios between the top-quark mass and the partonic center-of-mass energy in phase space integrals. Moreover, current and future demands of the LHC require a fast Monte Carlo generator, that provides reliable predictions at next-to-next-to-leading order, which can be flexibly adjusted to experimental setups. In view of those requirements, a complete new implementation of STRIPPER has been developed, which is based on an improved four-dimensional formulation of the subtraction scheme [34]. The whole framework is general, since it exploits the process independent soft and collinear factorization of tree-level and one-loop matrix elements as well as the universal singular structure of virtual one-loop and two-loop amplitudes [32, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48]. The remaining information about a specific process under consideration can be provided by interfacing the corresponding tree-level matrix elements as well as the finite parts of the one-loop and two-loop amplitudes. The implemented event generator includes concepts that have been proven and tested at next-to-leading order to increase speed and efficiency of the numerical calculation, e.g. Monte Carlo summation over partonic subprocesses and Monte Carlo summation over external polarizations. The software allows for a simultaneous calculation of different parton distribution functions, different renormalization and factorization scales and different observables.

For the specific case of top-quark pair production the tree-level matrix elements are obtained from Ref. [49]. The four-point one-loop amplitudes have been recomputed independently, but can be found in Refs. [50, 51, 52]. The five-point one-loop amplitudes are taken from the code used in Refs. [53, 54]. The two-loop amplitudes are numerically given in form of a dense grid [55, 56]. Partial analytic results for the two-loop amplitudes can be found in Refs. [57, 58, 59, 60].

RESULTS

The differential results presented in Figure 1 and Figure 2 are calculated using the following setup. The pole mass of the top-quark is set to $m_t = 173.3$ GeV. The distributions have been calculated using the MSTW2008 parton distribution function (PDF) set [61], where PDF uncertainties are not displayed. Each order includes the PDF set of the corresponding order. The theoretical uncertainties are obtained by varying the renormalization scale μ_R and factorization scale μ_F independently around the central scale $\mu_R = \mu_F = m_t$. The additional restriction, $0.5 < \mu_R/\mu_F < 2$, is imposed [62]. All results have been cross checked with available results: Integrals over the distributions reproduce the total cross section of Refs. [1, 2, 3, 4] to better than permil level. The next-to-leading order results are cross checked with the Monte Carlo event generator MCFM [63, 64].

Figure 1 shows the $p_{T,t}$, $m_{t\bar{t}}$ and $y_{t\bar{t}}$ differential distributions for the absolute normalization as well as the corresponding next-to-leading order and next-to-next-to-leading order K-factors. It should be noticed that the widths of bins of the calculation is much smaller than the widths of the bins that are currently used by the ATLAS and CMS collaborations, which allows to use these results for a variety of phenomenological analyses.

Each of the three distributions shows a good perturbative convergence. This behaviour, which has already been known for the total inclusive cross section, can be observed for each histogram bin separately. The central value lies within the error band of the previous order. This suggests that the theoretical predictions are robust within the stated error bands, while keeping in mind that the uncertainty of the PDF set is not included yet.

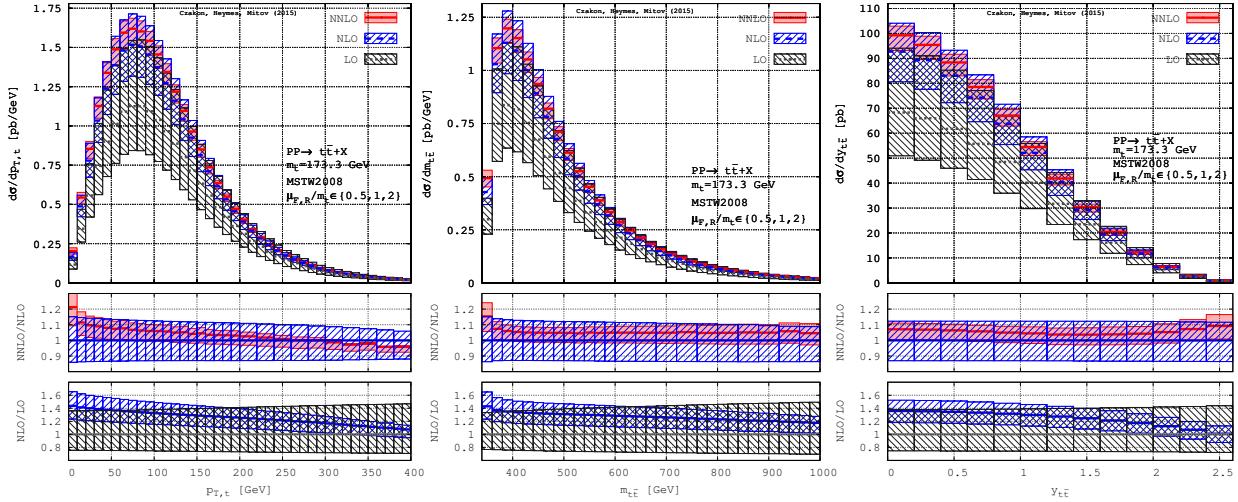


FIGURE 1. Top/antitop p_T -distribution (left), $m_{t\bar{t}}$ -distribution (center) and $y_{t\bar{t}}$ -distribution (right) in leading order (LO), next-to-leading order (NLO) and next-to-next-to-leading order (NNLO) QCD. Error bands are from scale variation [25].

A significant rise of the K-factors in the first bin of the $p_{T,t}$ -distribution and to a lesser degree in the first bin of the $m_{t\bar{t}}$ -distribution can be observed. However, this slight enhancement at low p_T and low $m_{t\bar{t}}$ is due to soft-gluon and Coulomb effects at the top-quark pair threshold, which are not captured within a fixed order calculation. Related work, in order to capture these effects to all orders can be found in Refs. [19, 20, 21, 22, 65, 66, 67, 68, 69]. A further investigation of these effects and matching resummed results to the fixed order prediction would be interesting. The next-to-next-to-leading order K-factor shows a significant slope in the p_T -distribution, which, however, is within the next-to-leading error band. The $m_{t\bar{t}}$ -distribution at next-to-next-to-leading order is remarkably stable with respect to the next-to-leading order prediction over the whole range of displayed values. It is therefore a well suited observable to look for new resonances beyond the Standard Model as suggested for example in Ref. [70]. The shape of the rapidity distribution of the top-quark pair at next-to-next-to-leading order changes slightly with respect to the next-to-leading order distribution, which is however not significant.

A first comparison of the p_T -distribution and the y_t -distribution of the top-quark with the CMS measurement in the lepton and jets channels from Ref. [15] is shown in Figure 2. Further comparisons with CMS data can be found in Ref. [71], while a comparison with the ATLAS measurements at $\sqrt{s} = 8$ TeV has been included in Ref. [72].

The distributions are normalized, such that integrating the displayed bins yields unity. Uncertainties due to scale variations for next-to-next-to-leading order distributions are shown exclusively. The lower panel displays the ratio of the Data, the leading order and the next-to-leading order distribution with respect to the next-to-next-to-leading order distribution. There is a slight mismatch between the normalization of the data and the prediction. The transverse momentum distribution of the data includes an additional bin, $400\text{GeV} < p_T < 500\text{GeV}$, which contributes 4 permil to the normalization. In the rapidity distribution, the last bin of the theory prediction extends up to $|y_t| < 2.6$, while the last bin of the measurement extends only up to $|y_t| < 2.5$. This can be seen explicitly in Figure 2.

The first observation is that the tension between data and the next-to-leading order result for the p_T -distribution is resolved at next-to-next-to-leading order. In each bin the Standard Model prediction gets closer to the CMS data and an agreement within the displayed uncertainties can be observed. The normalized rapidity distribution of the top-quark is stable with respect to next-to-next-to-leading order corrections. Within the experimental uncertainty the data seems to be well described by the next-to-next-to-leading order prediction.

SUMMARY AND OUTLOOK

In this write-up, differential distributions for top-quark pair production at next-to-next-to-leading order QCD for the LHC at $\sqrt{s} = 8$ TeV have been presented. These results have been obtained using a complete independent implementation of the subtraction scheme STRIPPER and include all partonic channels, where no approximations have been

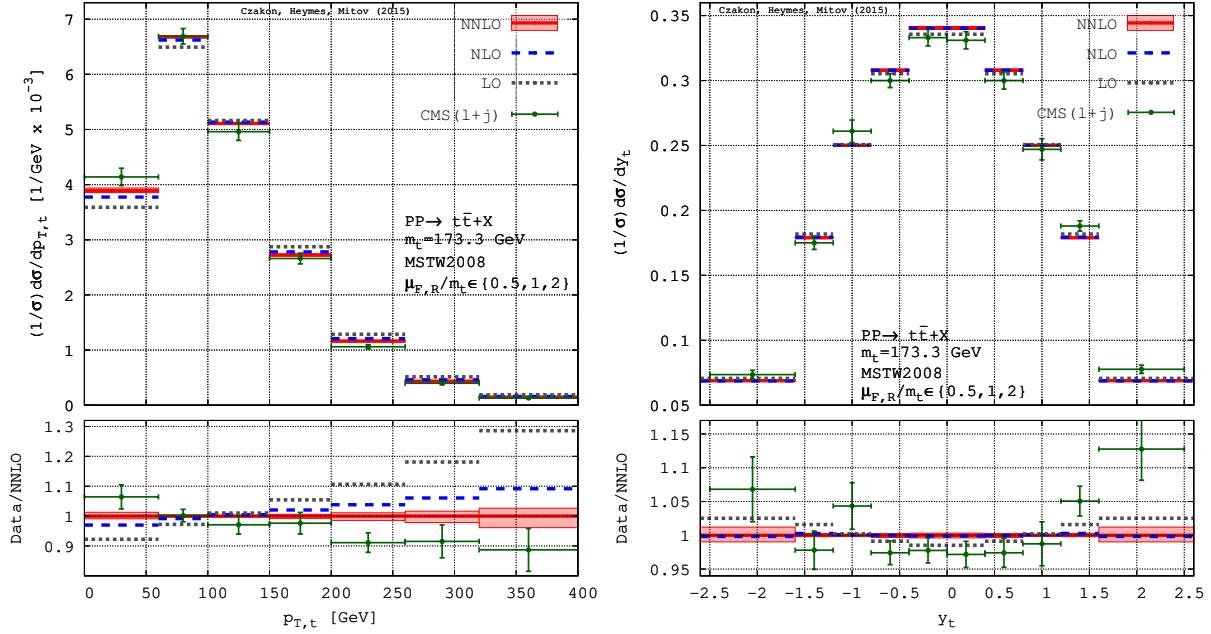


FIGURE 2. Normalised distribution for the p_T of the top/antitop-quark (left) and the rapidity y of the top/antitop-quark (right) in comparison with CMS measurements [15, 25]

made.

These results can be used for further phenomenological studies and tuning of Monte Carlo event generators in order to validate the estimated Standard Model background at the LHC. Moreover, they are valuable to constrain parton distribution functions using LHC data and to measure the strong coupling α_s . High precision Standard Model prediction allow to further explore limitations of the Standard Model and allow to constrain different scenarios beyond the Standard Model.

The presented results have been obtained for fixed scale settings $\mu_{F,R} \in \{m_t, m_t/2, 2m_t\}$, which allowed for several cross checks of the final result with previous calculations, e.g. the total inclusive cross section. However, for differential predictions this setting may not be the most appropriate one. In the near future results will become available using several dynamical scales and will include different PDF sets and PDF error estimates.

The calculation will be extended to a center-of-mass energy of $\sqrt{s} = 13$ TeV. In this context, it would be interesting to extend the kinematic regime of the $m_{t\bar{t}}$ - and p_T - distribution to higher values, since a significant number of top-quark pairs are produced in this so called *boosted* regime, $\sqrt{s} \gg m_t$. In this regime large logarithm can be resummed to all order to improve a fixed order calculation [73, 74]. A first comparison of the resummed results matched to next-to-leading order QCD and the full next-to-next-to-leading order at result $\sqrt{s} = 8$ TeV has been presented in Ref. [75]. An investigation of how the resummed result could be matched to the full next-to-next-to-leading order prediction would be interesting.

Finally, a particle level comparison with data would be possible, once decays of the top-quark are included into the calculation at next-to-next-to-leading order. It should be noticed that next-to-next-to-leading order corrections for on-shell top-quark decays have already been presented in Refs. [76, 77].

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REFERENCES

- [1] P. Bärnreuther, M. Czakon and A. Mitov, Phys. Rev. Lett. **109**, 132001 (2012) [arXiv:1204.5201 [hep-ph]].
- [2] M. Czakon and A. Mitov, JHEP **1212**, 054 (2012) [arXiv:1207.0236 [hep-ph]].
- [3] M. Czakon and A. Mitov, JHEP **1301**, 080 (2013) [arXiv:1210.6832 [hep-ph]].
- [4] M. Czakon, P. Fiedler and A. Mitov, Phys. Rev. Lett. **110**, no. 25, 252004 (2013) [arXiv:1303.6254 [hep-ph]].
- [5] M. Cacciari, M. Czakon, M. Mangano, A. Mitov and P. Nason, Phys. Lett. B **710**, 612 (2014) [arXiv:1111.5869 [hep-ph]].
- [6] M. Czakon and A. Mitov, Comput. Phys. Commun. **185**, 2930 (2014) [arXiv:1112.5675 [hep-ph]].
- [7] S. Chatrchyan *et al.* [CMS Collaboration], JHEP **1211**, 067 (2012) [arXiv:1208.2671 [hep-ex]].
- [8] S. Chatrchyan *et al.* [CMS Collaboration], JHEP **1402**, 024 (2014) [JHEP **1402**, 102 (2014)] [arXiv:1312.7582 [hep-ex], arXiv:1312.7582].
- [9] G. Aad *et al.* [ATLAS Collaboration], Eur. Phys. J. C **74**, no. 10, 3109 (2014) [arXiv:1406.5375 [hep-ex]].
- [10] V. Khachatryan *et al.* [CMS Collaboration], arXiv:1510.05302 [hep-ex].
- [11] [ATLAS Collaboration] ATLAS-CONF-2015-049.
- [12] S. Chatrchyan *et al.* [CMS Collaboration], Eur. Phys. J. C **73**, no. 3, 2339 (2013) [arXiv:1211.2220 [hep-ex]].
- [13] G. Aad *et al.* [ATLAS Collaboration], Phys. Rev. D **90**, no. 7, 072004 (2014) [arXiv:1407.0371 [hep-ex]].
- [14] G. Aad *et al.* [ATLAS Collaboration], JHEP **1506**, 100 (2015) [arXiv:1502.05923 [hep-ex]].
- [15] V. Khachatryan *et al.* [CMS Collaboration], arXiv:1505.04480 [hep-ex].
- [16] G. Aad *et al.* [ATLAS Collaboration], arXiv:1510.03818 [hep-ex].
- [17] [CMS Collaboration] CMS-PAS-TOP-15-005.
- [18] [CMS Collaboration] CMS-PAS-TOP-15-010.
- [19] V. Ahrens, A. Ferroglia, M. Neubert, B. D. Pecjak and L. L. Yang, JHEP **1009**, 097 (2010) [arXiv:1003.5827 [hep-ph]].
- [20] V. Ahrens, A. Ferroglia, M. Neubert, B. D. Pecjak and L. L. Yang, JHEP **1109**, 070 (2011) [arXiv:1103.0550 [hep-ph]].
- [21] N. Kidonakis, Phys. Part. Nucl. **45**, no. 4, 714 (2014) [arXiv:1210.7813 [hep-ph]].
- [22] N. Kidonakis, Phys. Rev. D **91**, no. 3, 031501 (2015) [arXiv:1411.2633 [hep-ph]].
- [23] M. Guzzi, K. Lipka and S. O. Moch, JHEP **1501**, 082 (2015) [arXiv:1406.0386 [hep-ph]].
- [24] A. Broggio, A. S. Papanastasiou and A. Signer, JHEP **1410**, 98 (2014) [arXiv:1407.2532 [hep-ph]].
- [25] M. Czakon, D. Heymes and A. Mitov, arXiv:1511.00549 [hep-ph].
- [26] S. Catani and M. H. Seymour, Nucl. Phys. B **485**, 291 (1997) [Nucl. Phys. B **510**, 503 (1998)] doi:10.1016/S0550-3213(96)00589-5 [hep-ph/9605323].
- [27] S. Frixione, Z. Kunszt and A. Signer, Nucl. Phys. B **467**, 399 (1996) doi:10.1016/0550-3213(96)00110-1 [hep-ph/9512328].
- [28] G. Abelof, A. Gehrmann-De Ridder and I. Majer, arXiv:1506.04037 [hep-ph].
- [29] G. Abelof and A. Gehrmann-De Ridder, JHEP **1412**, 076 (2014) [arXiv:1409.3148 [hep-ph]].
- [30] R. Bonciani, S. Catani, M. Grazzini, H. Sargsyan and A. Torre, arXiv:1508.03585 [hep-ph].
- [31] M. Czakon, Phys. Lett. B **693**, 259 (2010) [arXiv:1005.0274 [hep-ph]].
- [32] M. Czakon, Nucl. Phys. B **849** (2011) 250 [arXiv:1101.0642 [hep-ph]].
- [33] M. Czakon, P. Fiedler and A. Mitov, Phys. Rev. Lett. **115**, no. 5, 052001 (2015) [arXiv:1411.3007 [hep-ph]].
- [34] M. Czakon and D. Heymes, Nucl. Phys. B **890**, 152 (2014) [arXiv:1408.2500 [hep-ph]].
- [35] S. Catani and M. Grazzini, Nucl. Phys. B **570** (2000) 287 [hep-ph/9908523].
- [36] J. M. Campbell and E. W. N. Glover, Nucl. Phys. B **527**, 264 (1998) [hep-ph/9710255].
- [37] S. Catani and M. Grazzini, Phys. Lett. B **446**, 143 (1999) [hep-ph/9810389].
- [38] V. Del Duca, A. Frizzo and F. Maltoni, Nucl. Phys. B **568**, 211 (2000) [hep-ph/9909464].
- [39] Z. Bern, L. J. Dixon, D. C. Dunbar and D. A. Kosower, Nucl. Phys. B **425**, 217 (1994) [hep-ph/9403226].
- [40] Z. Bern, V. Del Duca and C. R. Schmidt, Phys. Lett. B **445**, 168 (1998) [hep-ph/9810409].
- [41] D. A. Kosower, Nucl. Phys. B **552**, 319 (1999) [hep-ph/9901201].
- [42] D. A. Kosower and P. Uwer, Nucl. Phys. B **563**, 477 (1999) [hep-ph/9903515].
- [43] Z. Bern, V. Del Duca, W. B. Kilgore and C. R. Schmidt, Phys. Rev. D **60**, 116001 (1999) [hep-ph/9903516].
- [44] G. Somogyi and Z. Trocsanyi, JHEP **0701**, 052 (2007) [hep-ph/0609043].
- [45] S. Catani and M. Grazzini, Nucl. Phys. B **591**, 435 (2000) [hep-ph/0007142].
- [46] S. Catani, S. Dittmaier and Z. Trocsanyi, Phys. Lett. B **500**, 149 (2001) [hep-ph/0011222].
- [47] I. Bierenbaum, M. Czakon and A. Mitov, Nucl. Phys. B **856**, 228 (2012) [arXiv:1107.4384 [hep-ph]].

- [48] A. Ferroglio, M. Neubert, B. D. Pecjak and L. L. Yang, JHEP **0911**, 062 (2009) [arXiv:0908.3676 [hep-ph]].
- [49] M. Bury and A. van Hameren, Comput. Phys. Commun. **196** (2015) 592 [arXiv:1503.08612 [hep-ph]].
- [50] J. G. Körner, Z. Merebashvili and M. Rogal, Phys. Rev. D **71**, 054028 (2005) [hep-ph/0412088].
- [51] C. Anastasiou and S. M. Aybat, Phys. Rev. D **78**, 114006 (2008) [arXiv:0809.1355 [hep-ph]].
- [52] B. Kniehl, Z. Merebashvili, J. G. Körner and M. Rogal, Phys. Rev. D **78**, 094013 (2008) [arXiv:0809.3980 [hep-ph]].
- [53] S. Dittmaier, P. Uwer and S. Weinzierl, Phys. Rev. Lett. **98**, 262002 (2007) [hep-ph/0703120 [HEP-PH]].
- [54] S. Dittmaier, P. Uwer and S. Weinzierl, Eur. Phys. J. C **59**, 625 (2009) [arXiv:0810.0452 [hep-ph]].
- [55] M. Czakon, Phys. Lett. B **664** (2008) 307 [arXiv:0803.1400 [hep-ph]].
- [56] P. Bärnreuther, M. Czakon and P. Fiedler, JHEP **1402**, 078 (2014) [arXiv:1312.6279 [hep-ph]].
- [57] R. Bonciani, A. Ferroglio, T. Gehrmann, A. von Manteuffel and C. Studerus, JHEP **1312**, 038 (2013) [arXiv:1309.4450 [hep-ph]].
- [58] R. Bonciani, A. Ferroglio, T. Gehrmann, A. von Manteuffel and C. Studerus, JHEP **1101**, 102 (2011) [arXiv:1011.6661 [hep-ph]].
- [59] R. Bonciani, A. Ferroglio, T. Gehrmann and C. Studerus, JHEP **0908**, 067 (2009) [arXiv:0906.3671 [hep-ph]].
- [60] R. Bonciani, A. Ferroglio, T. Gehrmann, D. Maitre and C. Studerus, JHEP **0807**, 129 (2008) [arXiv:0806.2301 [hep-ph]].
- [61] A. D. Martin, W. J. Stirling, R. S. Thorne and G. Watt, Eur. Phys. J. C **63**, 189 (2009) [arXiv:0901.0002 [hep-ph]].
- [62] M. Cacciari, S. Frixione, M. L. Mangano, P. Nason and G. Ridolfi, JHEP **0809**, 127 (2008) [arXiv:0804.2800 [hep-ph]].
- [63] J. M. Campbell and R. K. Ellis, arXiv:1204.1513 [hep-ph].
- [64] P. Nason, S. Dawson and R. K. Ellis, Nucl. Phys. B **303**, 607 (1988); Nucl. Phys. B **327**, 49 (1989) [Erratum-ibid. B **335**, 260 (1990)].
- [65] N. Kidonakis and G. F. Sterman, Nucl. Phys. B **505**, 321 (1997) [hep-ph/9705234].
- [66] E. Laenen, G. Oderda and G. F. Sterman, Phys. Lett. B **438**, 173 (1998) [hep-ph/9806467].
- [67] N. Kidonakis, E. Laenen, S. Moch and R. Vogt, Phys. Rev. D **64**, 114001 (2001) [hep-ph/0105041].
- [68] K. Hagiwara, Y. Sumino and H. Yokoya, Phys. Lett. B **666**, 71 (2008) [arXiv:0804.1014 [hep-ph]].
- [69] Y. Kiyo, J. H. Kuhn, S. Moch, M. Steinhauser and P. Uwer, Eur. Phys. J. C **60**, 375 (2009) [arXiv:0812.0919 [hep-ph]].
- [70] R. Frederix and F. Maltoni, JHEP **0901**, 047 (2009) doi:10.1088/1126-6708/2009/01/047 [arXiv:0712.2355 [hep-ph]].
- [71] CMS Collaboration [CMS Collaboration], CMS-PAS-TOP-15-011.
- [72] G. Aad *et al.* [ATLAS Collaboration], arXiv:1511.04716 [hep-ex].
- [73] A. Ferroglio, B. D. Pecjak and L. L. Yang, Phys. Rev. D **86**, 034010 (2012) doi:10.1103/PhysRevD.86.034010 [arXiv:1205.3662 [hep-ph]].
- [74] A. Ferroglio, S. Marzani, B. D. Pecjak and L. L. Yang, JHEP **1401**, 028 (2014) doi:10.1007/JHEP01(2014)028 [arXiv:1310.3836 [hep-ph], arXiv:1310.3836].
- [75] A. Ferroglio, B. D. Pecjak, D. J. Scott and L. L. Yang, arXiv:1512.02535 [hep-ph].
- [76] M. Brucherseifer, F. Caola and K. Melnikov, JHEP **1304**, 059 (2013) doi:10.1007/JHEP04(2013)059 [arXiv:1301.7133 [hep-ph]].
- [77] J. Gao, C. S. Li and H. X. Zhu, Phys. Rev. Lett. **110**, no. 4, 042001 (2013) doi:10.1103/PhysRevLett.110.042001 [arXiv:1210.2808 [hep-ph]].