

## Measurement of the neutron capture cross-section of $^{238}\text{U}$ using the neutron activation technique

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Significant effort has been aimed at generating nuclear power based on the concept of fast reactor [1-2] and advanced heavy water reactor (AHWR) [3]. In AHWR  $^{232}\text{Th}$ - $^{233}\text{U}$  in the oxide form is the primary fuel, whereas in the fast reactor  $^{238}\text{U}$ - $^{239}\text{Pu}$  in the form of carbide is used as the primary fuel. The  $^{239}\text{Pu}$  is first generated in a research reactor from  $^{238}\text{U}(\text{n}, \gamma)^{239}\text{U}$  reaction and by successive two beta decays. Then the fissile material  $^{239}\text{Pu}$  along with  $^{238}\text{U}$  is used as a fuel in fast reactor for power generation. The  $^{238}\text{U}$  is used as the breeding material to regenerate the fissile material  $^{239}\text{Pu}$ . In the fast reactor, there is a fast neutron spectrum. Thus for the production of  $^{239}\text{Pu}$ , it is necessary to have knowledge about  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section at various neutron energies. In the present work we have determined the  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section using the neutrons from  $^7\text{Li}(\text{p}, \text{n})$  reaction and by activation technique followed by off-line  $\gamma$ -ray spectrometry.

The experiment was carried out using the 14UD BARC-TIFR Pelletron facility at Mumbai, India. The neutron beam for irradiations of U was obtained from the  $^7\text{Li}(\text{p}, \text{n})$  reaction by using the 12 MeV proton beam of 400 nA at main line at 6 m above the analyzing magnet of the Pelletron facility. After irradiation and sufficient cooling, the  $\gamma$ -rays of fission/reaction products from the irradiated U sample were counted in an energy and efficiency calibrated 80 c.c. HPGe detector coupled to a PC-based 4K channel analyzer. From the observed photo-peak activity ( $A_{\text{obs}}$ ) for 743.3 keV  $\gamma$ -line of  $^{97}\text{Zr}$  the neutron flux ( $\Phi$ ) of  $1.3 \times 10^7 \text{ n cm}^{-2} \text{ s}^{-1}$  was obtained by using Eq.1, as explained in ref. [4]. Eq.1 can be

used for estimating  $\sigma$  when  $\Phi$  is known or vice versa.

$$A_{\text{obs}}(\text{CL/LT}) = N\sigma\Phi\alpha\epsilon(1-e^{-\lambda t})e^{-\lambda T}(1-e^{-\lambda CL})/\lambda \quad (1)$$

The neutrons from  $\text{Li}(\text{p}, \text{n})$  reaction are not mono-energetic, and their energy spectra were obtained from literature as shown in ref. [4]. The contribution to the neutron flux from the tail region is 49 % at the proton energy of 12.0 MeV. The average energy of neutrons under quasi mono-energetic peak is  $9.85 \pm 0.38$  MeV for 12 MeV protons after removing the tail of the spectrum.

The  $^{238}\text{U}(\text{n}, \gamma)$  cross-section ( $\sigma$ ) was calculated from the observed activity of 277.85 keV of  $^{239}\text{Np}$  in the  $\gamma$ -ray spectrum of the un-separated sample, which is  $2.333 \pm 0.123$  mb. The contribution from the tail region to  $^{238}\text{U}(\text{n}, \gamma)$  reaction has been estimated using the ENDF/B-VII [5] and JENDL-4.0 [6] by folding the evaluated cross-sections with neutron flux distributions. The contribution from the above evaluation at  $E_p = 12$  MeV are 1.023 and 0.614 mb from ENDF/B-VII and JENDL-4.0 respectively. The true value of  $^{238}\text{U}(\text{n}, \gamma)$  reaction-cross section due to the neutrons from the quasi mono-energetic peak consisting of the  $n_0$  and  $n_1$  groups of the neutron spectrum is obtained by subtracting these average cross-section due to neutron tail region from the above measured total cross section data. Thus the actual experimentally obtained  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-sections at average neutron energies of  $39.85 \pm 0.38$  MeV corresponding to proton energy of 12 MeV is  $1.42 \pm 0.09$  mb.

The present experimental  $^{238}\text{U}(\text{n}, \gamma)$  cross-sections are within the range of evaluated data of ENDF/B-VII and JENDL 4.0. However, the evaluated value from JEFF-3.1 [7] and CENDL-3 [8] are not in agreement with the present experimental value. In order to examine this aspect, the  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-sections from the present work and similar data from EXFOR [9] are plotted in Fig. 1. It can be seen from Fig. 1 that the  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section from present work at  $9.85 \pm 0.38$  MeV is in agreement with the value of McDaniels et al taken from EXFOR. Further, the  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section decreases from 100 keV to 7 MeV. Thereafter it decreases from neutron energy of 7 MeV to 15 MeV. Beyond this, it increases suddenly from neutron energy of 17 MeV to 20 MeV. In order to examine this, the evaluated data from ENDF/B-VII [5], JENDL-4.0 [6], JEFF-3.1 [7], CENDL [8] and INDC (VN)-8 [10] were plotted in Fig. 1. Similar data based on activation technique from the review article of Ding et al [11] were also plotted in Fig. 1 which agrees with evaluations only at low energies. The  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section at different neutron energy beyond 1 keV was also calculated theoretically using computer code TALYS of version 1.2 [12] and shown in Fig. 1. The trend of evaluated  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section is well reproduced by TALYS. However, they are slightly higher than the experimental and evaluated data.

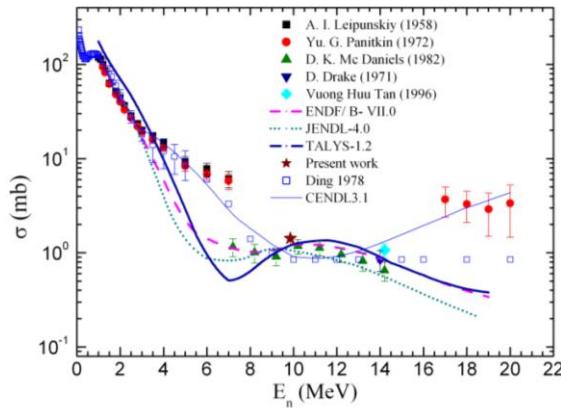


Fig. 1. Plot of experimental, evaluated and TALYS  $^{238}\text{U}(\text{n}, \gamma)$  reaction cross-section as a function of  $\text{E}_n = 1$  keV to 20 MeV.

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