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## Measurement of $\gamma$ rays from the giant resonances in $^{12}\text{C}$ and $^{16}\text{O}$ excited by the $(p, p')$ reaction.

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## Introduction

Oxygen and carbon are the third and fourth most abundant element by mass in the solar system, respectively [1] after hydrogen and helium, and  $^{16}\text{O}$  and  $^{12}\text{C}$  are their most abundant isotopes. Thus, they have been used as the target materials in the form of water and organic liquid scintillators in many large-scale neutrino experiments designed to detect low-energy neutrinos ( $E_\nu < 100$  MeV) [2–4]. These detectors must be massive to compensate for the extremely small neutrino cross section ( $\approx 10^{-42}$  cm $^2$ ). One of the most interesting applications is the detection of neutrinos from supernova explosion in our Galaxy [5]. The main reaction for neutrino detection is the charged-current (CC) anti-neutrino reaction with a proton ( $\bar{\nu}_e + p \rightarrow e^+ + n$ ), also known as the inverse  $\beta$ -decay reaction (IBD). Of special interest is the neutral-current (NC) neutrino or anti-neutrino inelastic scattering with  $^{16}\text{O}$  and  $^{12}\text{C}$ , followed by the emission of  $\gamma$  rays that can be observed with the de-

tector [6]. This process is of a special interest because the cross section is significant enough to be detected and is independent of neutrino oscillations.

The first observation of  $^{12}\text{C}(\nu, \nu')^{12}\text{C}^*(15.11 \text{ MeV}, 1^+, T = 1)$  reaction with 15.11-MeV  $\gamma$  ray came from the KARMEN experiment [3] with a neutrino beam. The observation was based on the detection of the electromagnetic decay of  $^{12}\text{C}$  excited by neutral current interactions. The  $\gamma$ -ray emission probability ( $\Gamma_\gamma/\Gamma$ ) of excited states of  $^{12}\text{C}$  below the proton separation energy ( $S_p = 16.0$  MeV) has been well measured [7]. However, the giant resonances appear above the separation energy and they decay mainly hadronically via particle emission ( $p$ ,  $n$ ,  $d$  and  $\alpha$ ) to the daughter nuclei. Although they decay mainly to the ground state of the daughter nuclei ( $^{11}\text{B}$ ,  $^{11}\text{C}$ , etc.), some of these decays are to excited states. If these excited states are below the particle emission threshold in  $^{11}\text{B}$  ( $S_{p'} = 11.2$  MeV) or  $^{11}\text{C}$  ( $S_{p'} = 8.7$  MeV), they decay by  $\gamma$ -ray emissions. Kolbe *et al.* and Langanke *et al.* [8, 9] proposed the above decay mechanism of giant resonances and estimated the NC

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neutrino and anti-neutrino reaction cross sections for  $^{12}\text{C}$  and  $^{16}\text{O}$ .

They stressed the importance of measuring NC events, since they are more sensitive to  $\nu_\mu$  and  $\nu_\tau$  neutrinos than to  $\nu_e$  neutrinos. However, there are no experimental measurements of  $\gamma$  rays from the giant resonances of these nuclei. In this paper, we report the first measurement of  $\gamma$  rays from the excited states of  $^{12}\text{C}$  and  $^{16}\text{O}$ , including giant resonances in the energy region  $E_x = 16\text{--}32$  MeV.

## 1. Experiment and Results

The experiment (E398) to measure the  $\gamma$  rays emitted from giant resonances in  $^{12}\text{C}$  and  $^{16}\text{O}$  was carried out at the Research Center for Nuclear Physics (RCNP), Osaka University. An unpolarized proton beam at 392 MeV bombarded a natural carbon ( $^{nat}\text{C}$ ) and cellulose ( $\text{C}_6\text{H}_{10}\text{O}_5$ ) target. The scattered protons were measured around  $0^\circ$  and were analyzed by the high-resolution magnetic spectrometer Grand Raiden (GR) [10–12]. The  $\gamma$  rays were measured in coincidence with the scattered protons using a  $\gamma$ -ray detector made from an array of  $5\times 5$  NaI(Tl) counters. The details of the experimental setup are described in Ref. [13].

We measured the double differential cross section ( $d^2\sigma/dE_x d\Omega$ ) for both  $^{12}\text{C}(p,p')$  and  $^{16}\text{O}(p,p')$  inelastic reaction at 392 MeV and  $0^\circ$  for the energy range  $E_x = 7\text{--}32$  MeV.

For the measurements of  $\gamma$  rays from the giant resonances, the absolute values of the  $\gamma$ -ray emission probability  $R_\gamma(E_x)$  and the response functions were verified using in-situ  $\gamma$  rays (15.1 and 6.9 MeV) with an accuracy of  $\pm 5\%$  during the experiment [13]. This calibration made it possible to measure  $R_\gamma(E_x)$  reliably as a function of the excitation energy of  $^{12}\text{C}$  and  $^{16}\text{O}$  in the energy range  $E_x = 16\text{--}32$  MeV. We found that the measured value of  $R_\gamma(E_x)$  starts from zero at  $E_x = 16$  MeV and increases to  $53.3 \pm 0.4 \pm 3.9\%$  at  $E_x = 27$  MeV and begins to decrease with further increase in  $E_x$ . For  $^{16}\text{O}$ , the measured value of  $R_\gamma(E_x)$  starts from  $21.1 \pm 0.6 \pm 2.0\%$  at  $E_x = 16$  MeV and increases to  $59.8 \pm 0.9 \pm 5.9\%$  at  $E_x = 25$  MeV, then decreases. We compared the mea-

surements of  $\gamma$ -ray emission probability with a statistical model calculation based on Hauser-Feshbach formalism [14, 15] and observed a 30–40% lower  $\gamma$ -ray emission probability in the energy region  $E_x = 20\text{--}24$  MeV than that predicted by the calculation.

The present results are very important for understanding the  $\gamma$ -ray emission probability of the giant resonances of the typical light nuclei ( $^{12}\text{C}$  and  $^{16}\text{O}$ ) and for the neutrino detection in liquid scintillator and water based detectors through neutral-current interactions.

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