



Physics of the Brout–Englert–Higgs boson: Theory

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Abstract

With the discovery of the Higgs boson a new era started with direct experimental information on the physics behind the breaking of the electroweak symmetry. This breaking plays a fundamental role in our understanding of particle physics and sits at the high-energy frontier beyond which we expect new physics that supersedes the Standard Model. The Higgs (inclusive and differential) production and decay rates offer a new way to probe this frontier.

Keywords:

Electroweak symmetry breaking, Higgs, Beyond the Standard Model

1. Introduction

The Standard Model (SM) is a triumph of the combination of the two pillars of twentieth-century physics, namely quantum mechanics and special relativity. Particles are defined as representations of the Poincaré group. Mathematically, these representations are labelled by two quantities: the spin that is quantized and takes only discrete values, and the mass, which a priori is a continuous parameter. However, the transformation laws for the various elementary particles under the gauge symmetries associated to the fundamental interactions force the masses of these particles to vanish. This would be in flagrant contradiction with the experimental measurements.

The Brout–Englert–Higgs mechanism (BEH) [1, 2, 3, 4] provides the solution to this mass conundrum. The discovery of a Higgs boson in July 2012 and the experimental confirmation of the BEH mechanism by the ATLAS and CMS collaborations [5, 6] has been a historical step in our understanding of nature: the masses of the elementary particles are not fundamental parameters defined at very high energy but rather emergent quantities appearing at low energy as a result of the particular structure of the vacuum.

This breakthrough discovery at the LHC surely stands out among the results of the first running period. For the first time a scalar particle has been found that assumes a vital role in the Standard Model and couples to all massive particles. In fact the Higgs particle dominates the structure of the vacuum.

2. The HEP landscape after the Higgs discovery

During its first run, the LHC certainly fulfilled its commitments: The machine and its detectors were mostly designed to find the Higgs boson and “[they] got it!” according to the words of R. Heuer, director general of CERN, on 4 July 2012. It was an important step in the understanding of the mechanism of electroweak symmetry breaking. But the journey is not over.

One can ask how the Higgs discovery reshaped the High Energy Physics (HEP) landscape. As it was emphasized in a recent conference by M. Mangano [7], the days of theoretically guaranteed discoveries imposed on us by some no-lose theorems are over: indeed, with the addition of a light Higgs boson with a mass around 125 GeV, the Standard Model is theoretically consistent and can be extrapolated up to very high energy, maybe as high as $10^{14\div 16}$ GeV or even the Planck scale. But

at the same time, the big questions of our field, or the ones that we have considered so far as the big questions, remain wildly open: the hierarchy problem, the origin of flavor, the issue of the neutrino masses and mixings, the question of the identity of Dark Matter, the source of dynamical preponderance of matter over antimatter during the cosmological evolution of our Universe. . . are left unanswered. In the next decades, future progress in HEP is in the hands of experimentalists whose discoveries will reveal the way Nature has solved these big questions, forcing the theorists to renounce/review/question deeply rooted bias/prejudice. The Higgs discovery sets a large part of the agenda for the theoretical and experimental HEP programs over the next couple of decades.

3. Open questions about the Higgs

Run 1 accumulated striking evidence that the Higgs field is the cause of the screening of the weak interaction at long distances. Indeed, the observation and measurement of the $H \rightarrow ZZ^* \rightarrow 4\ell$ channel indicate that the Higgs field develops a vacuum expectation value (vev) that is not invariant under the $SU(2)_L \times U(1)_Y$ gauge symmetry of the SM. Furthermore, this vev seems to be the common source of the Z-boson mass and the coupling between the Higgs boson and the Z boson.

However, this evidence only addresses the question of *how* the symmetry of the weak interaction is broken. It does not address the question of *why* the symmetry is broken or why the Higgs field acquires an expectation value. The situation is simply summarized in the following tautology

Why is electroweak symmetry broken?

Because the Higgs potential is unstable at the origin.

Why is the Higgs potential unstable at the origin?

Because otherwise EW symmetry would not be broken.

The discovery of a Higgs boson allowed first glimpses into a new sector of the microscopic world. Now comes the time of the detailed exploration of this new Higgs sector. And some key questions about the Higgs boson emerge:

1. Is it the SM Higgs?
2. Is it an elementary or a composite particle?
3. Is it unique and solitary? Or are there additional states populating the Higgs sector?
4. Is it eternal or only temporarily living in a metastable vacuum?
5. Is its mass natural following the criteria of Dirac, Wilson or 't Hooft?

6. Is it the first superparticle ever observed?
7. Is it really responsible for the masses of all the elementary particles?
8. Is it mainly produced by top quarks or by new heavy vector-like particles?
9. Is it a portal to a hidden world forming the dark matter component of the Universe?
10. Is it at the origin of the matter-antimatter asymmetry?
11. Has it driven the primordial inflationary expansion of the Universe?

The answers to these questions will have profound implications on our understanding of the fundamental laws of physics. Establishing that the Higgs boson is weakly coupled, elementary and solitary, would surely be as shocking as unexpected, but it may well indicate the existence of a multiverse ruled by anthropic selection rules. If instead deviations from the SM emerge in the dynamics of the Higgs, we will have to use them as a diagnostic tool of the underlying dynamics. The pattern of these deviations will carry indirect information about the nature of the completion of the SM at higher energies. In supersymmetric models, and more generally in models with an extended electroweak symmetry breaking sector, the largest deviations will be observed in the couplings to leptons and to the down-type quarks, as well as in the decay amplitudes to photons and gluons. In models of strong interactions, in which the Higgs boson is a bound state, the effects of compositeness uniformly suppress all the Higgs couplings while the self-interactions of the particles inside the Higgs sector, namely the Higgs particle and the longitudinal components of the W and Z bosons, will increase with the transferred energy. Moreover, the measurements of the Higgs couplings will also reveal the symmetry properties of the new state. For instance, it can be established whether the new scalar is indeed “a Higgs” fitting into a $SU(2)$ doublet together with the degrees of freedom associated with the longitudinal W and Z and not some exotic impostor, like for instance a pseudo-dilaton. If the Higgs is found to have an internal structure, a detailed study of the Higgs couplings can also establish whether it is just an ordinary composite, like a σ particle, or whether it is a pseudo-Nambu–Goldstone boson endowed with additional symmetry properties, like the π 's of QCD.

4. What is the SM Higgs the name of?

4.1. The SM Higgs boson as a UV regulator

The SM Higgs boson ensures the proper decoupling of the longitudinal polarizations of the massive EW

gauge bosons at high energy. Indeed, these longitudinal modes of $W^\pm$ and Z can be described by Nambu–Goldstone bosons associated to the coset $SU(2)_L \times SU(2)_R / SU(2)_{\text{isospin}}$. Their kinetic term corresponds to the gauge boson mass terms,

$$\frac{1}{2} m_Z^2 Z_\mu Z^\mu + m_W^2 W_\mu^+ W^{-\mu} = \frac{v^2}{4} \text{Tr}(D_\mu \Sigma^\dagger D^\mu \Sigma) \quad (1)$$

with $\Sigma = e^{i\sigma^a \pi^a / v}$, where σ^a ($a = 1, 2, 3$) are the usual Pauli matrices. Due to the Goldstone boson equivalence theorem the non-trivial scattering of the longitudinal gauge bosons V ($V = W^\pm, Z$) is controlled by the contact interactions among four pions from the expansion of the Lagrangian Eq. (1), leading to amplitudes growing with the energy,

$$\begin{aligned} \mathcal{A}(V_L^a V_L^b \rightarrow V_L^c V_L^d) &= \mathcal{A}(s) \delta^{ab} \delta^{cd} + \mathcal{A}(t) \delta^{ac} \delta^{bd} \\ &+ \mathcal{A}(u) \delta^{ad} \delta^{bc} \quad \text{with} \quad \mathcal{A}(s) \approx \frac{s}{v^2}. \end{aligned} \quad (2)$$

Here s, t, u denote the Mandelstam variables, and v represents the vacuum expectation value (VEV) with $v = 246$ GeV. The amplitude grows with the center-of-mass (c.m.) energy squared s , and therefore perturbative unitarity will be lost around $4\pi v \sim 3$ TeV, unless there is a new weakly coupled elementary degree of freedom. The simplest realization of new dynamics restoring perturbative unitarity is given by a single scalar field h , which is singlet under $SU(2)_L \times SU(2)_R / SU(2)_{\text{isospin}}$ and couples to the longitudinal gauge bosons and fermions as [8, 9, 10],

$$\mathcal{L}_{\text{EWSB}} = \frac{1}{2} (\partial_\mu h)^2 - V(h) + \quad (3)$$

$$\begin{aligned} &\frac{v^2}{4} \text{Tr}(D_\mu \Sigma^\dagger D^\mu \Sigma) \left(1 + 2a \frac{h}{v} + \sum_{n \geq 2} b_n \frac{h^n}{v^n} + \dots \right) - \\ &\frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^j) \Sigma \left(1 + c \frac{h}{v} + \sum_{n \geq 2} c_n \frac{h^n}{v^n} + \dots \right) \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.} \end{aligned}$$

with

$$V(h) = \frac{1}{2} m_h^2 h^2 + \frac{d_3}{6} \left(\frac{3m_h^2}{v} \right) h^3 + \frac{d_4}{24} \left(\frac{3m_h^2}{v^2} \right) h^4 + \dots \quad (4)$$

For $a = 1$ the scalar exchange cancels the piece growing with the energy in the $V_L V_L$ amplitude. If in addition $b_2 = a^2$ then also in the inelastic amplitude $V_L V_L \rightarrow hh$ unitarity is maintained, while for $ac = 1$ the $V_L V_L \rightarrow f\bar{f}'$ amplitude remains finite. The SM Higgs boson is defined by the point $a = b_2 = c = 1$ and $d_3 = d_4 = 1$, $c_{n \geq 2} = b_{n \geq 3} = 0$. The scalar resonance and the pions then combine to form a doublet which transforms linearly under $SU(2)_L \times SU(2)_R$.

The requirement that the Higgs boson fully unitarizes the scattering amplitudes of massive particles therefore implies that the Higgs couplings to the various SM particles are directly proportional to their masses. This fundamental property is in remarkable agreement with the state-of-the-art fit of the current Higgs data collected at the LHC, see Fig. 1.

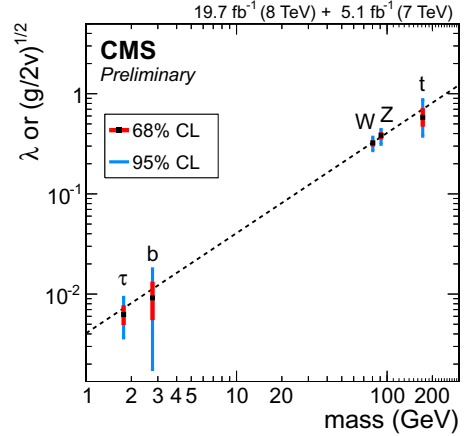


Figure 1: Comparison of the SM predictions (black dashed line) and the fit of the LHC measurements of the Higgs couplings to various SM particles. From Ref. [11].

4.2. The flavor preserving nature of the SM Higgs

In the SM, the Yukawa interactions are the only source of the fermion masses and also generate linear interaction with the physical Higgs boson

$$Y_{ij} \bar{\psi}_i H \psi_j = \frac{Y_{ij} v}{\sqrt{2}} \bar{\psi}_i \psi_j + \frac{Y_{ij}}{\sqrt{2}} h \bar{\psi}_i \psi_j. \quad (5)$$

Clearly both matrices can be diagonalized simultaneously and this ensures the absence of flavor changing neutral currents induced by the Higgs exchange.

This nice property is no longer true if the SM fermions mix with vector-like partners or in the presence of generic higher dimension Yukawa interactions (see for instance Ref. [12] for a general phenomenological discussion):

$$\begin{aligned} Y_{ij} \left(1 + \frac{c_{ij}}{f^2} |H|^2 \right) \bar{\psi}_i H \psi_j &= \frac{Y_{ij} v}{\sqrt{2}} \left(1 + \frac{c_{ij} v^2}{2f^2} \right) \bar{\psi}_i \psi_j \\ &+ \frac{Y_{ij}}{\sqrt{2}} \left(1 + \frac{3c_{ij} v^2}{2f^2} \right) h \bar{\psi}_i \psi_j. \end{aligned} \quad (6)$$

Therefore it is particularly important to probe the flavor structure of the Higgs interactions and to look for flavor-violating decays, e.g. $h \rightarrow \mu\tau$, or production,

e.g. $t \rightarrow hc$. Limits from low-energy flavor-changing interactions indirectly constrain these processes especially in the quark sector but leave the possibility of sizeable effects in the lepton sector [13, 14, 15, 16]. Interestingly enough, the CMS measurement [17] reports a 2.5σ excess in the $h \rightarrow \mu\tau$ decay and starts probing the interesting region of parameter space where the off-diagonal Yukawa couplings are set by the mass of the leptons, $|Y_{\mu\tau}Y_{\tau\mu}| \sim m_\tau m_\mu/v^2$, one order of magnitude better than the indirect bounds set by $\tau \rightarrow \mu\gamma$ and $\tau \rightarrow 3\mu$. For some recent interpretations of this excess, see Refs. [18, 19, 20, 21, 22, 23].

5. The Higgs mass as a model-discriminator

The value of the Higgs boson mass opens many decay modes at a rate accessible experimentally. Two channels are particularly accurate in accessing the Higgs mass itself: $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 2\ell^+2\ell^-$. Table 1 reports the ATLAS and CMS measurements, reaching a 0.2%-accuracy (to be compared to the 0.5%-accuracy of the top mass measurement).

Under the assumption that the SM laws govern Nature up to very high energy, the precise value of the Higgs mass has thrilling implications on the stability of the EW vacuum and hence the fate of our Universe (see for instance Ref. [25] for a recent update and for further references).

The value of the Higgs mass also gives clues about the details of possible Ultra-Violet (UV) completions of the SM itself. This can be exemplified in the leading scenarios, namely the Minimal Supersymmetric Model (MSSM) and the Minimal Composite Higgs model (MCHM). In short, the Higgs mass is larger than what is typically expected in the MSSM and smaller than what is expected in the MCHM. At the classical/Born level, the mass of the lightest (SM-like) Higgs boson is bounded to be lower than the Z-boson mass since supersymmetry dictates the Higgs quartic to be fixed in terms of the gauge couplings. Some significant amount of radiative corrections, mostly from the top and stop sectors, are therefore called to raise the value of the Higgs mass. At one-loop, the Higgs mass can be approximated by

$$M_h^2 \simeq M_Z^2 \cos^2 2\beta + \frac{3\sqrt{2}G_F M_t^4}{16\pi^2} \left[\log \frac{M_t^2}{M_i^2} + \frac{X_t^2}{M_i^2} \left(1 - \frac{X_t^2}{12M_i^2} \right) \right], \quad (7)$$

where $M_i^2 = M_{Q_3} M_{U_3}$ is the geometric mean of the stop masses and X_t is the mixing between the two stops. Clearly, a Higgs boson as heavy as 125 GeV requires

either heavy stops ($M_t > 800$ GeV) and/or maximally mixed stops ($X_t \simeq \sqrt{6}M_t$), which brings back some amount of irreducible fine-tuning or call for non-trivial boundary conditions like non-universal gaugino masses at high-energy. Going beyond the minimal model, for instance by adding an extra gauge singlet, can easily help increasing the Higgs mass with significantly less amount of tuning, see for instance Ref. [26] for a recent discussion.

In the Minimal Composite Higgs models, the Higgs boson emerges from the strong sector as a pseudo-Nambu–Goldstone boson. Therefore, the strong interactions themselves are not responsible for generating a potential for the Higgs field, that is generated only at the one-loop level from the interactions between the strong sector and the SM. Computing the details of the potential from first principles remains out of reach but it is possible [27, 28] to estimate the Higgs mass using general properties about the asymptotic behavior of correlators, i.e. imposing the saturation of the Weinberg sum rules with the first few light resonances, to obtain

$$M_h^2 \simeq \frac{3M_t^2 M_Q^2}{\pi^2 f^2}, \quad (8)$$

where f is the scale of the strong interactions (the decay constant of the Higgs boson, the equivalent of f_π for the QCD pions) and M_Q is the typical mass scale of the fermion resonances (aka the top partners). This estimate can read as

$$M_Q \simeq 700 \text{ GeV} \left(\frac{M_h}{125 \text{ GeV}} \right) \left(\frac{160 \text{ GeV}}{M_t} \right) \left(\frac{f}{500 \text{ GeV}} \right). \quad (9)$$

For a natural set-up ($v^2/f^2 \leq 0.2$), we therefore expect some light top partners below one TeV. The discovery of such fermionic top-partners would be a first evidence of a strong dynamics at the origin of the breaking of the electroweak symmetry.

5.1. The Higgs profile as a probe of new physics

A dedicated study of the Higgs boson properties and couplings offers a way to infer what the structure of physics beyond the Standard Model can be. Natural models trying to give a rationale for why/how the Higgs mass is screened from high energy corrections at the quantum level generically predict some deviations in the Higgs couplings compared to the SM predictions of the order 1% to 100%. The current Higgs data accumulated at the LHC by the ATLAS and CMS collaborations already constrain the Higgs couplings to massive gauge bosons and to fermions not to deviate by more than 20–30% from the SM predictions [24, 29].

Experiment	$H \rightarrow \gamma\gamma$	$H \rightarrow ZZ^* \rightarrow 4\ell$	combined
ATLAS	$125.98 \pm 0.42(\text{stat.}) \pm 0.28(\text{syst.})$	$124.51 \pm 0.52(\text{stat.}) \pm 0.06(\text{syst.})$	$125.36 \pm 0.37(\text{stat.}) \pm 0.18(\text{syst.})$
CMS	$124.70 \pm 0.31(\text{stat.}) \pm 0.15(\text{syst.})$	$125.59 \pm 0.42(\text{stat.}) \pm 0.17(\text{syst.})$	$125.02 \pm 0.27(\text{stat.}) \pm 0.15(\text{syst.})$

Table 1: ATLAS and CMS fits of the Higgs mass (in GeV) in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 2\ell^+2\ell^-$ channels and their combination. From Refs. [24].

In general, new physics can deform the SM in many ways but most of these deformations are already severely constrained by electroweak precision measurements or flavor data. Assuming flavor universality among the couplings between the Higgs boson and the SM fermions, it was shown [30, 31] that eight directions among the leading CP-conserving deformations of the SM can be probed, at tree-level, only in processes with a physical Higgs boson. These deformation induce deviations in the Higgs couplings that respect the Lorentz structure of the SM interactions, or generate simple new interactions of the Higgs boson to the W and Z field strengths, or induce some contact interactions of the Higgs boson to photons (and to a photon and a Z boson) and gluons that take the form of the ones that are generated by integrating out the top quark. In other words, the Higgs couplings are described, in the unitary gauge, by the following effective Lagrangian [32, 33]

$$\begin{aligned}
\mathcal{L} = & \kappa_3 \frac{m_H^2}{2v} H^3 + \kappa_Z \frac{m_Z^2}{v} Z_\mu Z^\mu H + \kappa_W \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} H \\
& + \kappa_g \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} H + \kappa_\gamma \frac{\alpha}{2\pi v} A_{\mu\nu} A^{\mu\nu} H \\
& + \kappa_{Z\gamma} \frac{\alpha}{\pi v} A_{\mu\nu} Z^{\mu\nu} H \\
& + \kappa_{VV} \frac{\alpha}{2\pi v} \left(\cos^2 \theta_W Z_{\mu\nu} Z^{\mu\nu} + 2 W_{\mu\nu}^+ W^{-\mu\nu} \right) H \\
& - \left(\kappa_t \sum_{f=u,c,t} \frac{m_f}{v} \bar{f}_L f_R + \kappa_b \sum_{f=d,s,b} \frac{m_f}{v} \bar{f}_L f_R \right. \\
& \left. + \kappa_\tau \sum_{f=e,\mu,\tau} \frac{m_f}{v} \bar{f}_L f_R + h.c. \right) H. \tag{10}
\end{aligned}$$

In the SM, the Higgs boson does not couple to massless gauge bosons at tree level, hence $\kappa_g = \kappa_\gamma = \kappa_{Z\gamma} = 0$. Nonetheless, the contact operators are generated radiatively by loops of SM particles. In particular, the top quark gives a contribution to the 3 coefficients $\kappa_g, \kappa_\gamma, \kappa_{Z\gamma}$ that does not decouple in the infinite top mass limit. For instance, in that limit $\kappa_\gamma = \kappa_g = 1$ [34, 35].

The coefficient for the contact interactions of the Higgs boson to the W and Z field strengths is not in-

dependent but obeys the relation

$$(1 - \cos^4 \theta_W) \kappa_{VV} = \sin 2\theta_W \kappa_{Z\gamma} + \sin^2 \theta_W \kappa_\gamma. \tag{11}$$

This relation is a general consequence of the so-called custodial symmetry [36]. When the Higgs boson is part of an $SU(2)_L$ doublet, the custodial symmetry could only be broken by a single operator at the level of dimension-6 operators and it is accidentally realized among the interactions with four derivatives, like the contact interactions considered. Custodial symmetry also implies

$$\kappa_Z = \kappa_W, \tag{12}$$

leaving exactly 8 free couplings [30, 31]. Out of these 8 coefficients, only κ_V can be indirectly constrained by EW precision data at a level comparable from the direct constraints from LHC Higgs data [37, 38, 39].

The effective Lagrangian of Eq. (10) can be amended by 6 extra Higgs couplings that break the CP symmetry

$$\begin{aligned}
\mathcal{L} = & \tilde{\kappa}_g \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} H + \tilde{\kappa}_\gamma \frac{\alpha}{2\pi v} A_{\mu\nu} \tilde{A}^{\mu\nu} H \\
& + \tilde{\kappa}_{Z\gamma} \frac{\alpha}{\pi v} A_{\mu\nu} \tilde{Z}^{\mu\nu} H \\
& - i \left(\tilde{\kappa}_t \sum_{f=u,c,t} \frac{m_f}{v} \bar{f}_L f_R + \tilde{\kappa}_b \sum_{f=d,s,b} \frac{m_f}{v} \bar{f}_L f_R \right. \\
& \left. + \tilde{\kappa}_\tau \sum_{f=e,\mu,\tau} \frac{m_f}{v} \bar{f}_L f_R + h.c. \right) H, \tag{13}
\end{aligned}$$

where $\tilde{F}_{\mu\nu} = \epsilon_{\mu\nu\rho\sigma} F^{\rho\sigma}$ is the dual field-strength of $F_{\mu\nu}$. It is certainly tempting to consider new sources of CP violation in the Higgs sector, potentially bringing in one of the necessary ingredients for a successful baryogenesis scenario. The prospects for measuring at the LHC these CP violating sources in the $h\bar{t}\bar{t}$, $h\tau\bar{\tau}$ and $h\gamma\gamma$ couplings have been studied in Refs. [40], [41] and [42] respectively. It should however be noted [43] that these CP violating couplings would induce quark and electron electric dipole moments at one- (for $\tilde{\kappa}_\gamma$ and $\tilde{\kappa}_{Z\gamma}$) or two-loops (for $\tilde{\kappa}_f$). Unless the Yukawa couplings of the Higgs to the electron and light quarks are significantly

reduced compared to their SM values, these constraints severely limit the possibility to observe any CP violating signal in the Higgs sector at the LHC.

The coefficient κ_3 can be accessed only through double Higgs production processes, hence it will remain largely unconstrained at the LHC and a future machine like an ILC or a future circular collider might be needed to pin down this coupling [44, 45, 46]. The LHC will also have a limited sensitivity on the coefficient κ_τ since the lepton contribution to the Higgs production cross section remains subdominant and the only way to access the Higgs coupling is via the $H \rightarrow \tau^+ \tau^-$ and possibly $H \rightarrow \mu^+ \mu^-$ channels. Until the associated production of a Higgs with a pair of top quarks is observed, the Higgs coupling to the top quark is only probed indirectly via the one-loop gluon fusion production or the radiative decay into two photons. However, these two processes are only sensitive to the two combinations $(\kappa_t + \kappa_g)$ and $(\kappa_t + \kappa_\gamma)$ and a deviation in the Higgs coupling to the top quark can in principle always be masked by new contact interactions to photons and gluons [47, 48, 49, 50]. In the next section, we shall discuss how individual information on $\kappa_{\gamma,g,t}$ can be obtained by studying either the hard recoil of the Higgs boson against an extra jet or the off-shell Higgs production in $gg \rightarrow h^* \rightarrow ZZ \rightarrow 4\ell$.

The operators already bounded by EW precision data modify in general the Lorentz structure of the Higgs couplings and hence induce some modifications of the kinematical differential distributions [51, 52]. A promising way to have a direct access to the Wilson coefficients of these operators in Higgs physics is to study the VH associated production with a W or a Z at large invariant mass [51, 53]. It has not been estimated yet whether the sensitivity on the determination of the Wilson coefficients in these measurements can compete with the one derived for the study of anomalous gauge couplings. In any case, these differential distributions could also be a way to directly test the hypothesis that the Higgs boson belongs to an $SU(2)_L$ doublet together with the longitudinal components of the massive electroweak gauge bosons. A sensitive issue when using the information on the VH production at large $m(HV)$ -invariant mass is the validity/reliability of the effective field theory treatment [54, 55, 56, 57].

Various dynamics produce different patterns among the Higgs coupling deviations. Table 2 summarizes the largest effects expected in popular classes of models of new physics addressing the hierarchy problem. The correlations among these deviations can thus reveal the nature of the dynamics above the weak scale while their magnitude will indicate the scale of this new dynamics.

m_H [GeV]	$\frac{\sigma_{\text{NLO}}(m_t)}{\sigma_{\text{NLO}}(m_t \rightarrow \infty)}$	$\frac{\sigma_{\text{NLO}}(m_t, m_b)}{\sigma_{\text{NLO}}(m_t \rightarrow \infty)}$
125	1.061	0.988

Table 3: Finite top mass effects in the inclusive gluon fusion production. From Ref. [61].

6. A sketch of the Higgs sector at energy scales above the Higgs mass

So far the LHC has mostly produced Higgs bosons on-shell in processes with a characteristic scale around the Higgs mass. This gives a rather good portrait of the Higgs couplings around the weak scale itself. However to fully accomplish its role as a UV regulator of the scattering amplitudes, what matter are the couplings of the Higgs at asymptotically large energy. To probe the Higgs couplings at large energy, one can rely to the associated production with additional boosted particle(s) but the price to pay is to deal with significantly lower production rates.

6.1. Boosted Higgs

The dominant production mode of the Higgs at the LHC is the gluon fusion channel. This is a purely radiative process. The lightness of the Higgs boson plays a malicious role and makes it impossible to disentangle short- and long-distance contribution to the total rate. This limitation is embodied in the Higgs low energy theorem [34, 35] that prevents to resolve the loop contribution itself (the NLO gluon fusion inclusive cross section for a finite and infinite top mass differ only by 1%, see Table 3). New Physics could modify the top Yukawa and also generate a contact interaction to the gluons without leaving any impact on the total rate, provided that $\kappa_t + \kappa_g = 1$. Concrete examples are top partners in composite Higgs models or mixed stops in the MSSM. Still, extra radiation in the $gg \rightarrow h$ process will allow us to explore the structure of the top loop. When the extra radiation carries away a large amount of energy and boosts the Higgs boson, the process effectively probes the ultraviolet structure of the top loop. Notice that the extra radiation cannot be in the form of a photon, as the amplitude for $gg \rightarrow h + \gamma$ vanishes due to Furry's theorem. One is therefore led to consider the production of h in association with a jet.¹

Table 4 reports the evolution of the cross section as a function of the cut on the p_T of the Higgs [62]. Figure 2

¹The process $gg \rightarrow hZ$ is also interesting, despite being subleading with respect to the tree-level contribution from $q\bar{q} \rightarrow hZ$. See Refs. [59, 60] for very recent studies.

Models	κ_f	κ_V	κ_g	κ_γ	$\kappa_{Z\gamma}$	κ_3
MSSM	✓					✓
NMSSM	✓	✓	✓	✓	✓	✓
MCHM aka PNCB Higgs	✓	✓			✓	✓
SUSY Composite Higgs	✓	✓				✓
Higgs as a Dilaton			✓	✓	✓	✓
Partly-Composite Higgs			✓	✓	✓	✓
Bosonic Technicolor						✓

Table 2: Largest deviations in the Higgs couplings expected in a variety of new physics models. From Ref. [58].

$\sqrt{s}$ [TeV]	$p_T^{\min}$ [GeV]	$\sigma_{p_T^{\min}}^{\text{SM}}$ [fb]	gg, qg [%]
14	100	2180	67, 31
	200	351	65, 34
	300	74.9	61, 38
	400	19.9	56, 43
	500	6.24	52, 47
	600	2.22	48, 51
	800	0.368	42, 57

Table 4: Summary table of the cross sections for $pp \rightarrow h + \text{jet}$ at proton-proton colliders with $\sqrt{s} = 14$ TeV as function of the cut on the Higgs transverse momentum, $p_T > p_T^{\min}$. The fourth column shows the fraction of the SM cross section coming from the partonic subprocesses gg and qg (the contribution of the $q\bar{q}$ channel is always smaller than 2%). From Ref. [62].

gives an idea of the sensitivity on the boosted analysis to resolve the κ_t/κ_g degeneracy plaguing the inclusive rate measurement [62]. Similar results have been obtained in Refs. [63, 64] and a more realistic analysis of $h \rightarrow 2\ell + \text{jet}$ via $h \rightarrow \tau\tau$ and $h \rightarrow WW^*$ has been performed in Ref. [65].

It should be noted that the $gg \rightarrow h + \text{jet}$ process has been computed only at leading order with the full mass dependence on the fermion running in the loops. The theoretical uncertainty can be estimated by relying on the NNLO K -factors computed in the $m_t \rightarrow \infty$ limit. It is however clear that an exact NLO computation of the SM Higgs p_T spectrum would be very welcome, and we hope that the QCD community will fill this gap in the near future.²

6.2. Off-shell Higgs

As for any other quantum particle, the influence of the Higgs boson is not limited to its mass shell. Recently,

²A first step in this direction has been made in Ref. [66]. For recent progress in the predictions for $h + \text{jet}$, see also Refs. [67, 68].

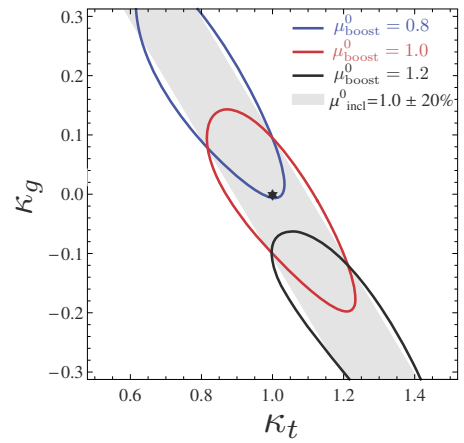


Figure 2: 95% CL contours obtained from a χ^2 fit of the inclusive and boosted (with $150 \text{ GeV} < p_T < 650 \text{ GeV}$) cross sections. The inclusive measurement, assumed to agree with the SM within 20%, constrains the parameter space to the light gray band, while the addition of the boosted cross section information shrinks the band to an ellipse. The colors blue, red and black correspond to a measured boosted cross section equal to 0.8, 1 and 1.2 times the SM prediction. The SM point is indicated by the black star. An integrated luminosity of $\int \mathcal{L} dt = 3 \text{ ab}^{-1}$ and $\sqrt{s} = 14 \text{ TeV}$ is assumed. From Ref. [62].

the CMS and ATLAS collaborations reported the differential cross-section measurement of $pp \rightarrow Z^{(*)}Z^{(*)} \rightarrow 4\ell, 2\ell 2\nu$ ($\ell = e, \mu$) at high invariant-mass of the ZZ system [69, 70]. This process receives a sizable contribution from a Higgs produced off-shell by gluon fusion [71, 72]. As such, this process potentially carries information relevant for probing the EFT at large momenta and could thus reveal the energy-dependence of the Higgs couplings controlled by higher-dimensional operators with extra derivatives. It has been proposed [73] to use the off-shell Higgs data to bound, in a model-independent way, the Higgs width. However, as it was emphasized in Ref. [74], this bound actually holds under the assumption that the Higgs couplings re-

main the same over a large range of energy scales. Instead the off-shell measurement offers a rather unique access to the structure of the Higgs couplings at high energy. Again this channel reveals to be particularly efficient to disentangle the long and short distance contribution to the Higgs production by gluon fusion.

Figure 3 gives an idea of the sensitivity on the off-shell analysis to resolve the κ_t/κ_g degeneracy plaguing the inclusive rate measurement [75]. Similar results have been obtained in Ref. [76, 77, 78, 79]. For a recent discussion of off-shell Higgs production within the SM and beyond and an extensive list of references, see Ref. [80].

An important issue that must be taken into account when simulating $gg \rightarrow ZZ$ is that the Higgs contribution is known to NLO, $O(\alpha_s^3)$, [81, 82, 83] in QCD with exact top mass dependence and to NNLO, $O(\alpha_s^4)$, [84, 85, 86] in the infinite top mass limit, whereas the interfering background is known only at LO, $O(\alpha_s^2)$. As a consequence, assessing the higher order corrections to the full process is problematic. Some estimates have been done and used but a full computation is still in demand.

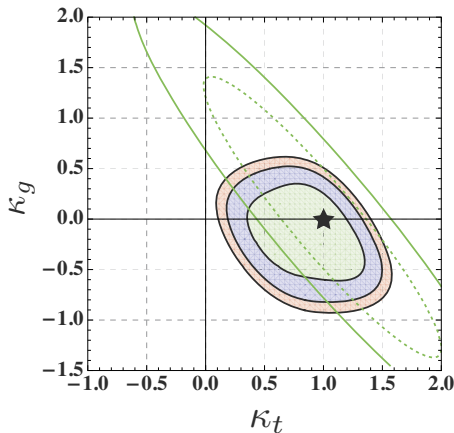


Figure 3: Prospects for a 14 TeV off-shell analysis with an integrated luminosity of 3 ab^{-1} and for the injected SM signal: 68%, 95% and 99% expected probability regions in the (κ_t, κ_g) plane. The dashed and solid green lines indicate the 68% and 95% contours for the linear analysis when only the interference between the higher-dimensional operator and the SM is kept for the computation of the signal yields. No theoretical uncertainty is included. The SM point is indicated by the black star. From Ref. [75].

Multi-Higgs productions, either in vector boson fusion, in gluon fusion or in association with a pair of top quarks, could provide further interesting information to complete the Higgs portrait [9, 87, 88, 89, 46]. But they remain challenging channels at the LHC, even in its high-luminosity run.

7. Conclusions

The first run of the LHC operations crowned the Standard Model as the successful description of the fundamental constituents of matter and their interactions to the tiniest details, from the QCD jet production over many orders of magnitude, to the multiple productions of electroweak gauge bosons as well as the production of top quarks. Undeniably, the Higgs boson discovery will remain the acme of the LHC run 1 and it has profound theoretical and phenomenological implications. In a few months from now, the second run of the LHC will start, with an increased center-of-mass energy, $\sqrt{s} = 13 \text{ TeV}$. This run will provide us in the coming years with crucial experimental information on the physics behind the breaking of the electroweak symmetry and it carries the hopes to finally reveal the first cracks in the SM grounds. If naturalness turned out to be a good guide, the LHC should soon find new states and revolutionize the field. If we are not so lucky and such new states are too heavy for the LHC reach, we might still detect indirectly their presence through the deviations they can induce on the Higgs properties. Precise measurements of such properties are therefore crucial and could be extremely useful to guide future direct searches at higher energies, either at the LHC itself or at other future facilities.

The Higgs boson might also be a portal to a hidden sector whose existence is anticipated to account for the total matter and energy budget of the Universe. The Higgs boson could also be one key agent in driving the early exponentially growing phase of our Universe and thus allowing large scale structures to emerge from original quantum fluctuating seeds.

Whatever the outcome of the next LHC run, we are exploring new territory and living in exciting times! The future editions of ICHEP hold great promises.

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