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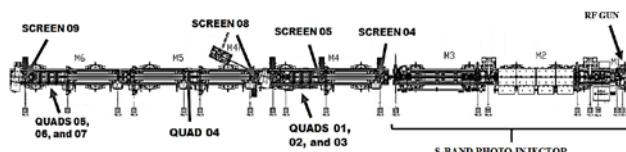


Figure 2: Location of quadrupole magnets and profile monitors in the low energy linac [2].

### Transfer Line

The transfer line starts at the dipole magnet. There are two quadrupole triplets. The dispersion of the beam after dipole magnet 01 is corrected by dipole 02. Figure 3 shows the components of the transfer line.

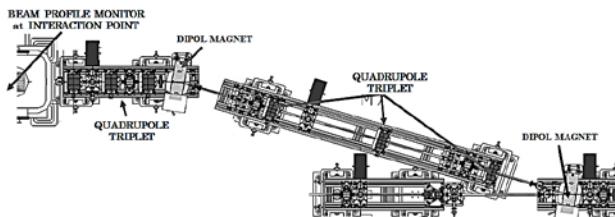


Figure 3: Location of quadrupole magnets and profile monitors in transfer line [2].

## RESULTS

### Elegant Simulations

For the simulations with the Elegant code 50 000 particles were used. The beam charge  $Q = 250\text{pC}$ , output beam energy at the injector exit  $E = 81\text{MeV}$ , bunch length  $\sigma_z \approx 280\mu\text{m}$ , the transverse normalized emittance ( $1\sigma$ ) at the injector exit (screen 04), calculated from the Elegant output data is  $\epsilon_{n,x,y} \approx 0.43 [\mu\text{m} \cdot \text{rad}]$ . Input Twiss parameters [9] are:  $\beta_x = 31.69$ ,  $\alpha_x = -1.98$ ,  $\eta_x = 0$ ,  $\beta_y = 31.71$ ,  $\alpha_y = -1.99$ ,  $\eta_y = 0$ .

Figure 4 below presents the evolution of the transverse beam profile along the low energy linac, and transfer line.

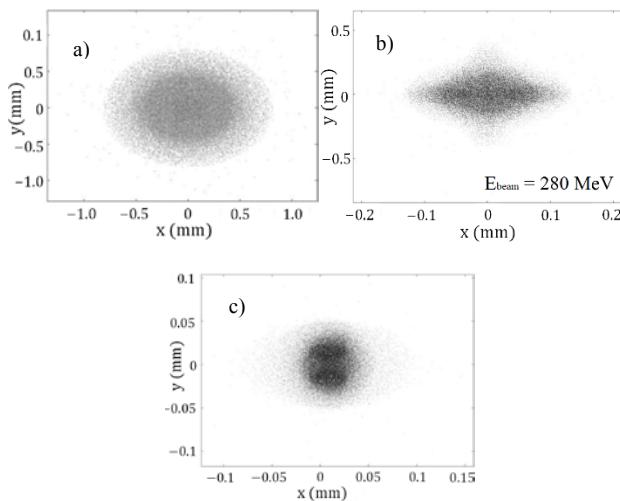


Figure 4: Transverse beam profiles at: injector exit (a), end of low energy linac (b), and interaction point IP1 (c).

The transverse (Figure 5) and longitudinal (Figure 6) beam distributions are the histograms of three coordinates  $x, y, z$ . The  $z$  coordinate is calculated from the relationship  $z = \beta c t$ , where  $\beta c$  is the velocity of the particle, and  $t$  is time of arrival at the observation point. Figures below include the  $rms$  sizes of the bunch.

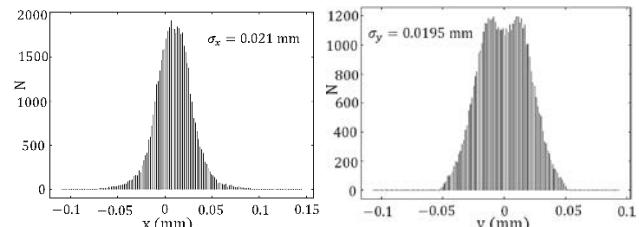


Figure 5: Transverse beam distribution at IP1.

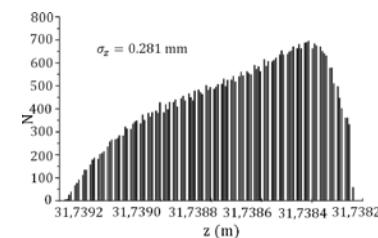


Figure 6: Longitudinal beam distribution at IP1.

A good approximation for the beam shape in phase space is elliptic. An ellipse can be defined by specifying: area, shape, and orientation. The *beta* function  $\beta(s)$  is related to the beam shape and size. The beam envelope is determined by the beam emittance and the function  $\beta(s)$ . The *beta* function is highly dependent on the particular arrangement of the quadrupole magnets. The *alpha* function  $\alpha(s)$  is related to the tilt (i.e. orientation) of the beam ellipse. Figures 7 and 8 depict the functions  $\beta(s)$  and  $\alpha(s)$  along the low energy linac and the transfer line up to IP1. The *gamma* function depends on  $\alpha$ - and  $\beta$ -functions in accordance to the relation  $\gamma(s) = \frac{1+\alpha^2(s)}{\beta(s)}$ .

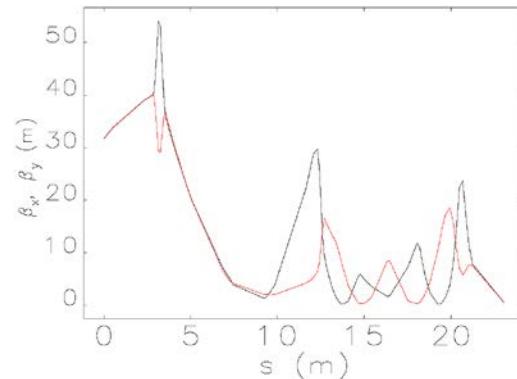


Figure 7: Function *beta* from photo-injector exit down to low energy interaction point IP1 ( $\beta_x$ -black line,  $\beta_y$ -red line).

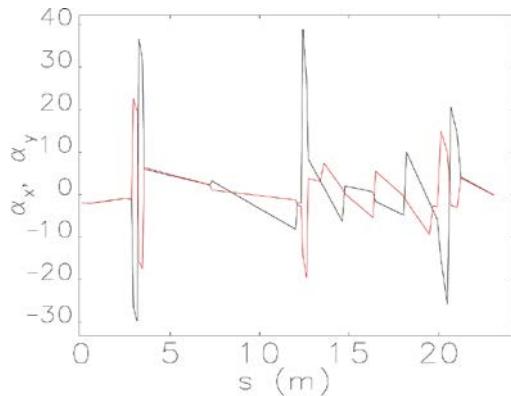


Figure 8: Function *alpha* from photo-injector exit down to low energy interaction point IP1 ( $\alpha_x$ -black line,  $\alpha_y$ -red line).

The normalized emittance ( $1\sigma$ ) at the IP1 calculated from the Elegant output data is  $\epsilon_{nxy} = 0.43\text{--}0.45 \text{ } [\mu\text{m} \cdot \text{rad}]$ .

In Fig. 9 the transverse phase space at IP1 is given.

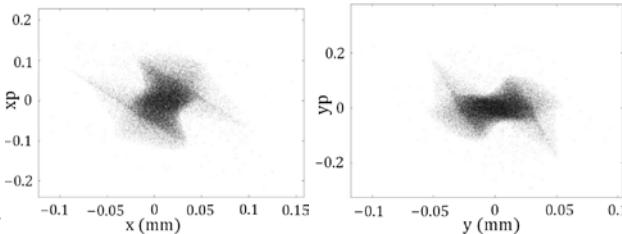


Figure 9: Transverse (horizontal – left, vertical - right) phase space at IP1.

Figure 10 presents the beam energy spread and beam energy at IP1, and Fig. 11 the energy distribution along the bunch.

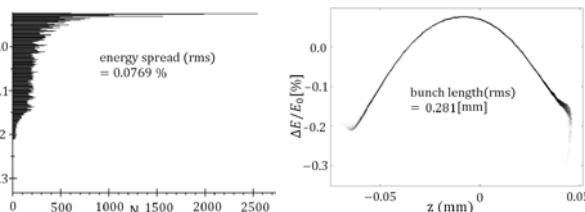


Figure 10: Beam energy spread (left) and energy distribution (right) at IP1.

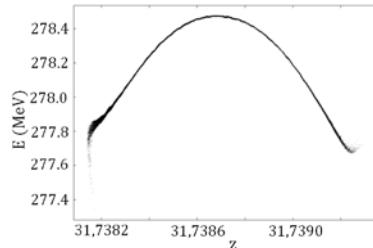


Figure 11: Energy distribution along bunch at IP1.

## CONCLUSIONS

We presented the results obtained in simulating the beam transport in the low energy linac and the transfer line for the ELI-NP Gamma Beam System. We used the Elegant code. Our study will be extended for the high energy linac. Calculations of the injector will also be done with use of the Astra code.

## ACKNOWLEDGEMENTS

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